

Leading Neutrons at HERA

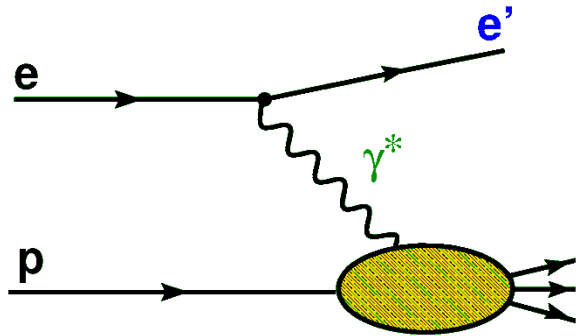
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On behalf of the H1 and ZEUS collaborations



International Workshop on Diffraction in High-Energy Physics
La Londe-les-Maures, France, September 9-14 2008

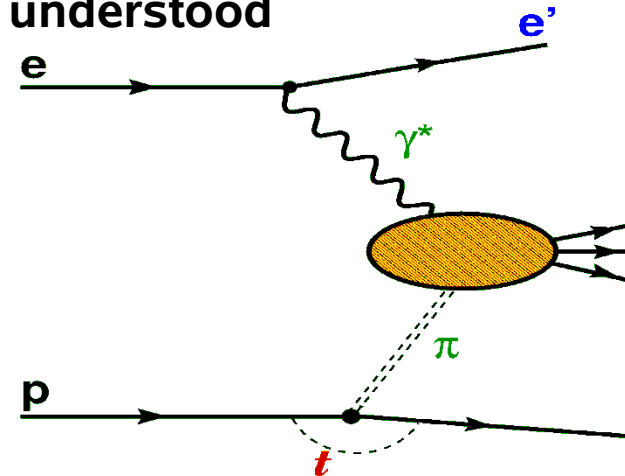
Motivation

- Significant fraction of ep scattering events contain a leading neutron in the final state carrying a substantial portion of the energy of the incoming proton:
 $e+p \rightarrow e+n+X$
- Production mechanism is not yet completely understood



Leading neutron can come from “standard fragmentation”

- From hadronization of p remnant
- Implemented in MC models (e.g. Lund)



Leading neutron can be produced via exchange of virtual particle

- charged iso-vector (π^+ , ρ^+)

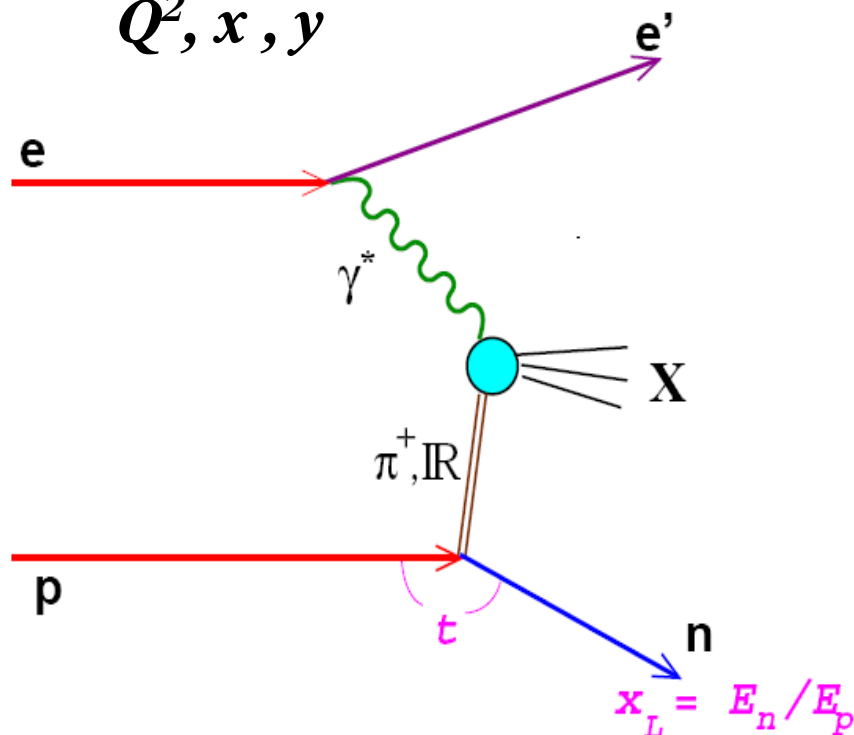
Results discussed in this talk:

- Leading Neutron spectra in DIS and photoproduction
- Leading Neutron production cross section
- Comparison with models

Photon vertex:

lepton variables:

$$Q^2, x, y$$



Proton vertex:

leading baryon variables:

$$x_L = E_n / E_p$$

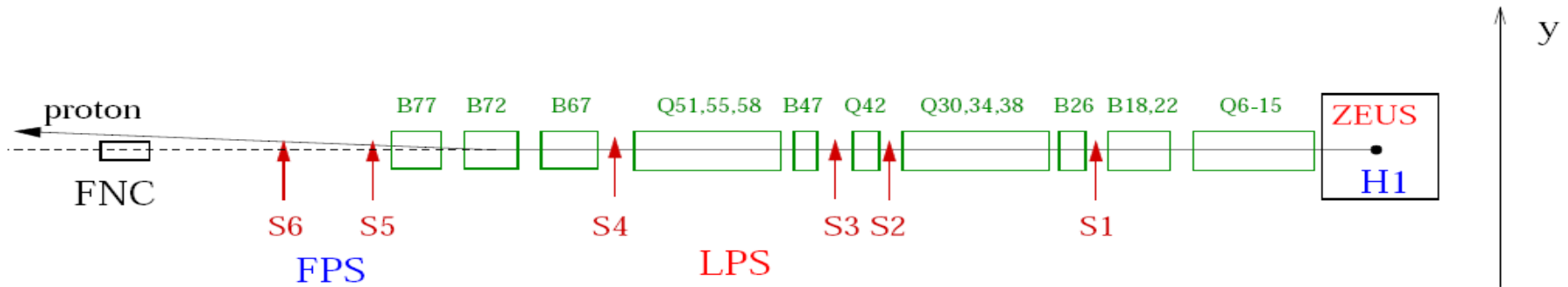
$$t = (p - p')^2 \approx -\frac{p_T^2}{x_L} - \frac{(1 - x_L)^2}{x_L} m_n^2$$

In the π -exchange model cross section dependence on leading baryon variables is independent of kinematics at photon vertex

$$\sigma_{ep \rightarrow enX}(x, Q^2, x_L, t) = f_{\pi^+/p}(x_L, p_T) \cdot \sigma_{e\pi \rightarrow eX}(x/(1 - x_L), Q^2)$$

Leading Neutron production affected by absorption and rescattering effects:
evidences of vertex factorization violation

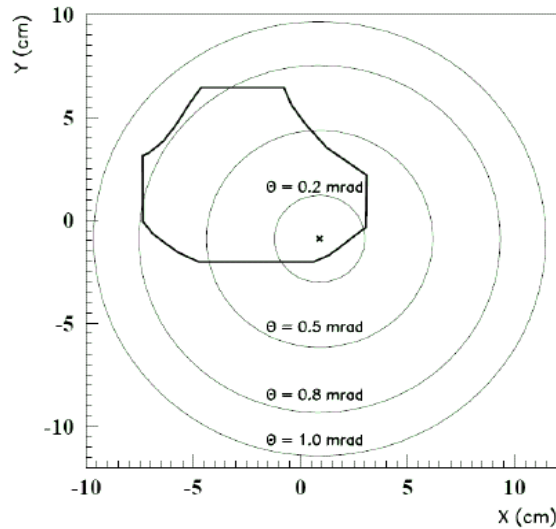
H1 and ZEUS Forward Neutron Calorimeters



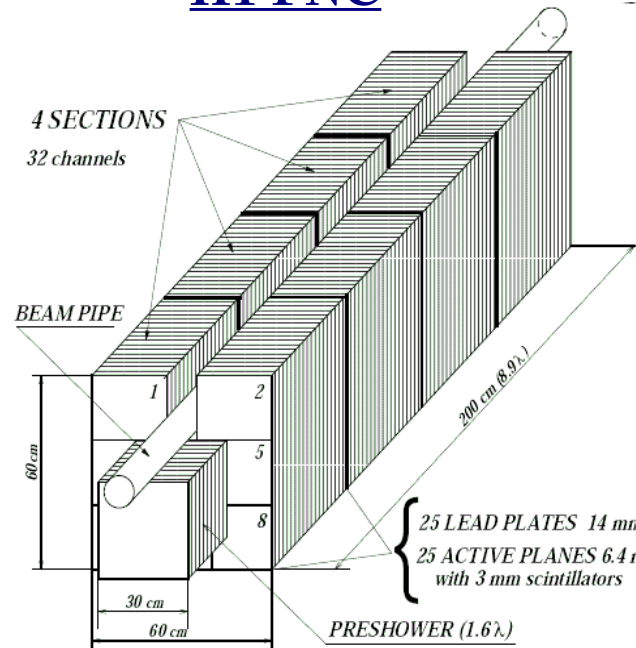
Position – 105m from the interaction point.

Geometrical acceptance is limited by beam-line elements: < 0.8 mrad.

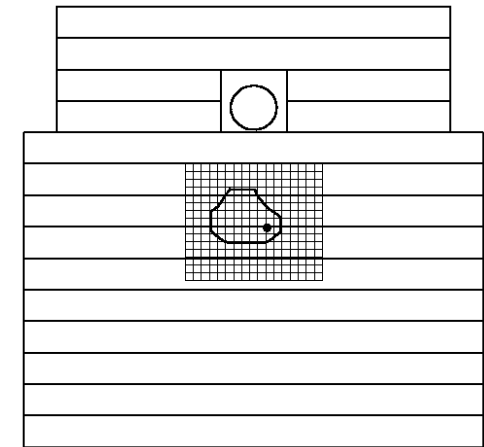
Geometrical Acceptance



H1 FNC



ZEUS FNC



14 towers, 17x15 grid of
the FNT hodoscopes,
 $\sigma_E/E \approx 0.7/\sqrt{E}$

$\sim 10 \lambda$ lead-scintillator sandwich calorimeters

H1: $F_2^{\text{LN}(3)}(Q^2, x, x_L)$ - triple diff. cross section of LN production in DIS

- HERA-II data, much higher statistics compared to old result H1prelim-08-111
- Upgraded Forward Neutron Calorimeter (better resolution and photon identification)
- Kinematic range: $6 \text{ GeV}^2 < Q^2 < 100 \text{ GeV}^2$, $1.5 \times 10^{-4} < x < 3 \times 10^{-2}$, $p_{T,n} < 0.2 \text{ GeV}$

ZEUS: relative to inclusive cross section vs x_L and p_T in DIS and γp

Leading Neutrons are selected from inclusive data sets (i.e. no LN tag): Nucl.Phys.B776 (2007) 1

- DIS: $Q^2 > 2\text{-}3 \text{ GeV}^2$, 3 subsets $Q^2 \approx 2.7, 8.9, 40 \text{ GeV}^2$
- γp : $Q^2 < 0.02 \text{ GeV}^2$, e^+ tagged $\Rightarrow 150 < W_{\gamma p} < 270 \text{ GeV}$

Leading Neutron yields:

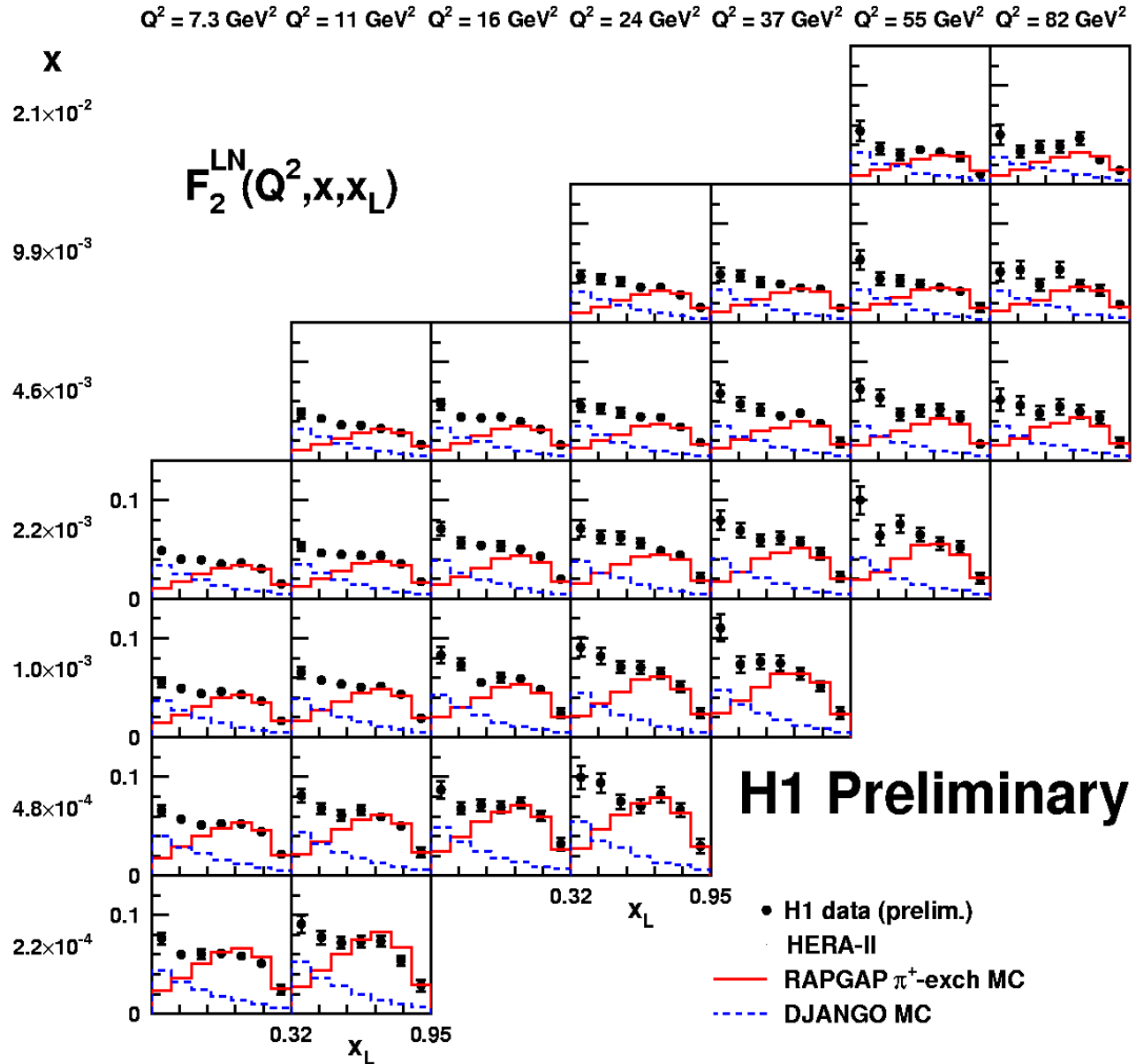
- DIS, γp have very different inclusive cross sections σ_{inc}
- For sensible comparisons look at LN yields: $r_{\text{LB}} = \sigma_{\text{LB}}/\sigma_{\text{inc}}$
- Additional benefit: systematic uncertainties of central detector cancel

Triple differential reduced cross section

$$\frac{d^3 \sigma(ep \rightarrow eNX)}{dQ^r dx dx_L} =$$

$$= \frac{4\pi\alpha^2}{x Q^\epsilon} \left[1 - y + \frac{y^2}{\Upsilon} \right] F_2^{\text{LN}}(Q^r, x, x_L)$$

- DJANGO (standard fragmentation) predicts too low cross section + x_L spectrum shape is too different
- RAPGAP π^+ -exch. describes data well for $x_L > 0.7$



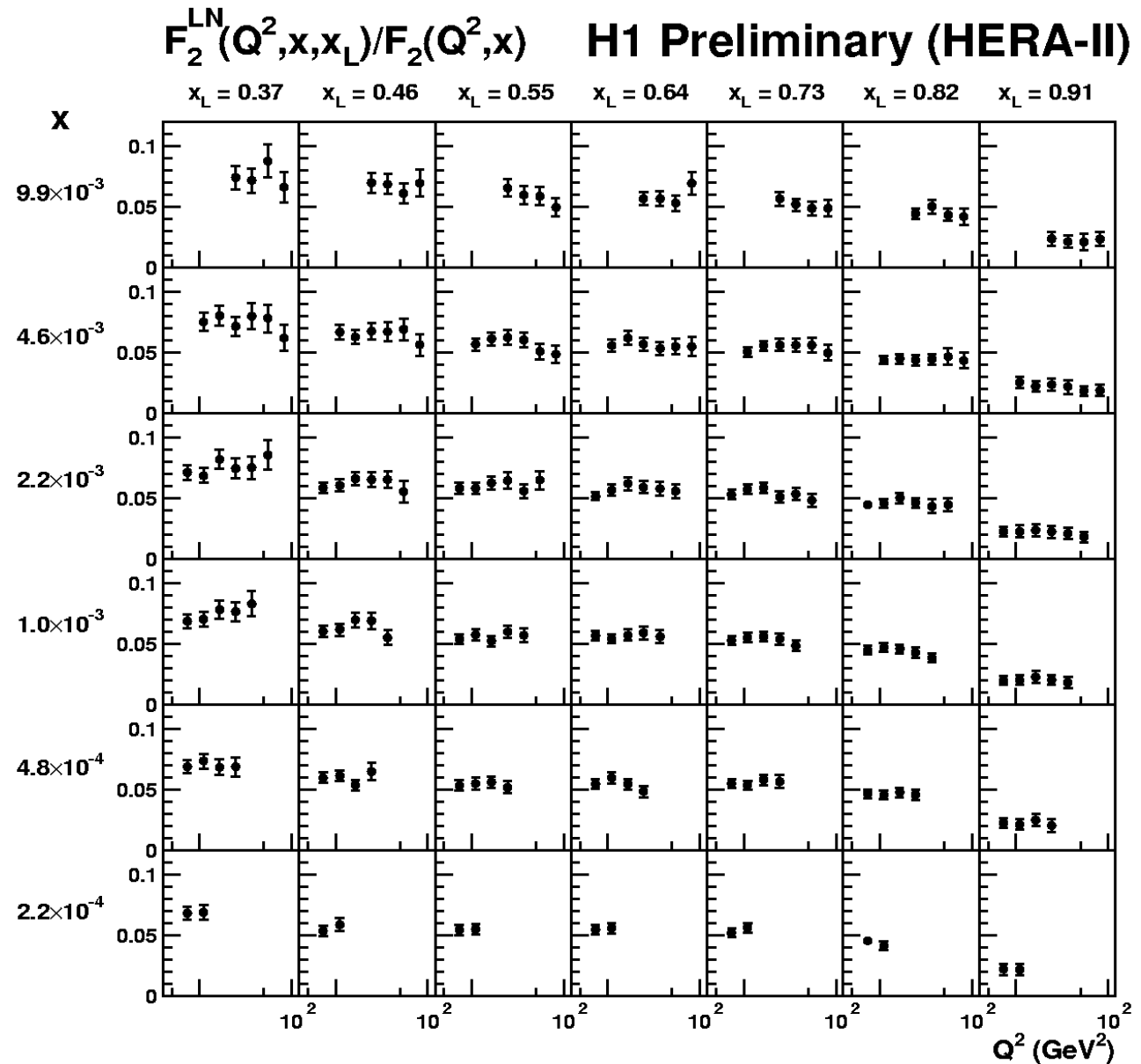
$F_2^{\text{LN}}(Q^2, x, x_L)$ to $F_2(Q^2, x)$ ratio

$F_2(Q^2, x)$ from the H1 parameterization
(Eur.Phys.J.C21 (2001) 33)

$F_2^{\text{LN}}(Q^2, x, x_L)/F_2(Q^2, x)$
is mostly flat in Q^2 and x

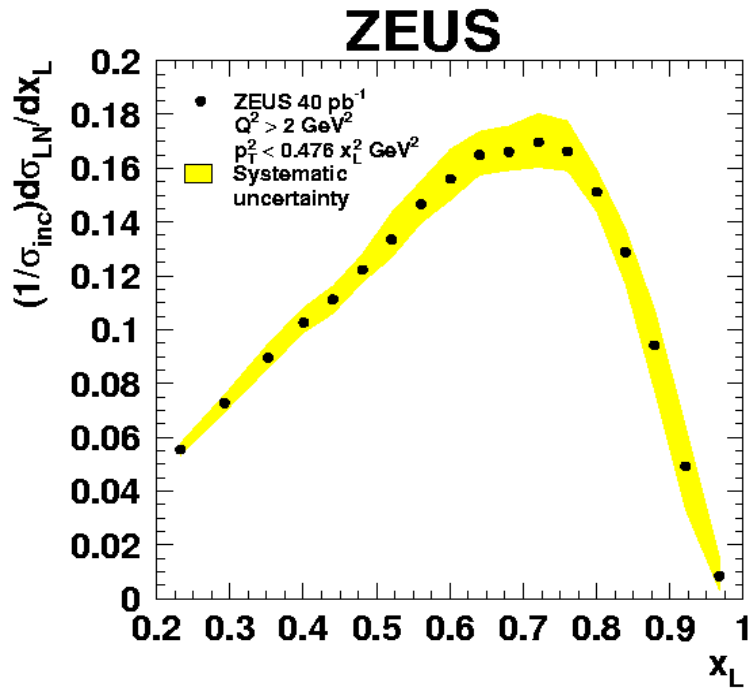
=> consistent with vertex factorization
(or common suppression of LN cross section)

**LN production rate,
kinematics is approx.
independent of (Q^2, x)**



DIS: x_L and p_T^2 distributions

Leading Neutron: $p_T^2 < 0.476 x_L^2 \text{ GeV}^2$

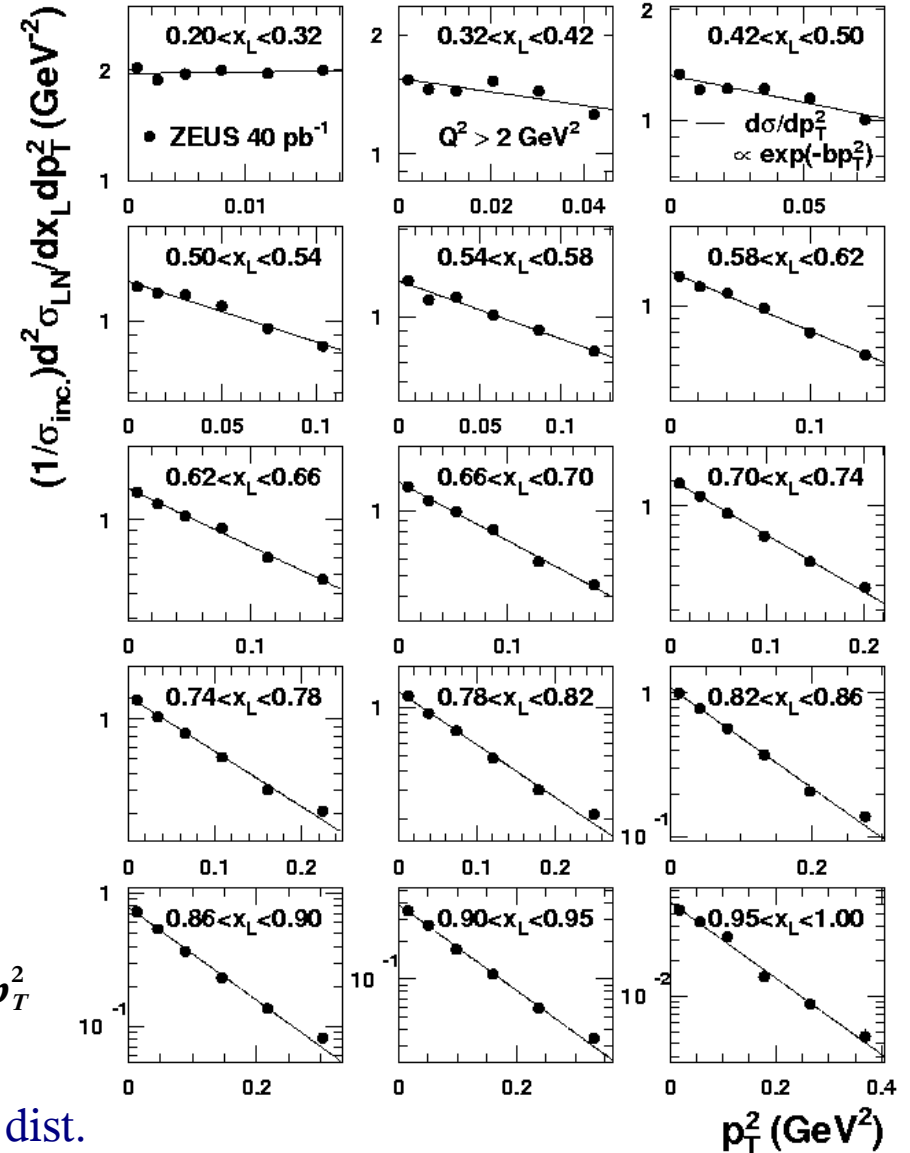


- LN yield $\rightarrow 0$ at kin. limit $x_L \rightarrow 1$
- Below $x_L \approx 0.7$ yield drops due to decreasing p_T^2 range

Fit by exp. in p_T^2 :
$$\frac{1}{\sigma_{\text{inc}}} \frac{d^2 \sigma_{\text{LN}}}{dx_L dp_T^2} = a(x_L) \cdot e^{-b(x_L) p_T^2}$$

Intercept $a(x_L)$ and slope $b(x_L)$ fully describe (x_L, p_T^2) dist.

Log Scale

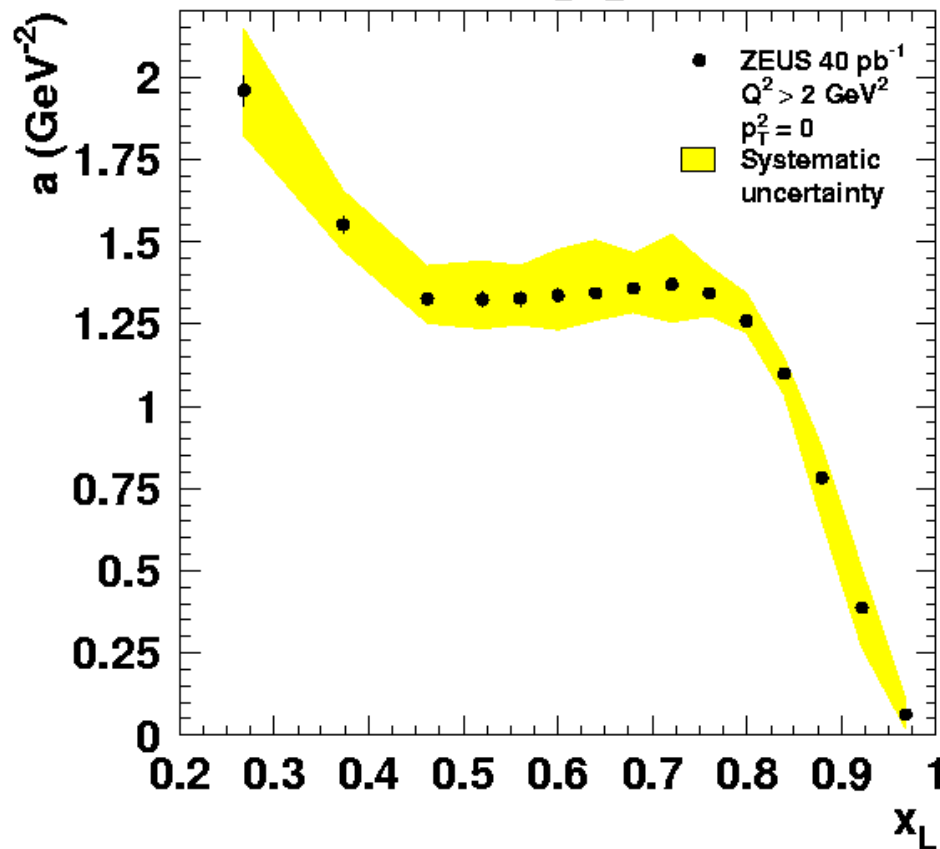


DIS p_T^2 distributions: intercepts and slopes

$$\frac{1}{\sigma_{\text{inc}}} \frac{d^2 \sigma_{\text{LN}}}{dx_L dp_T^2} = a(x_L) \cdot e^{-b(x_L) p_T^r}$$

Intercepts $a(x_L)$

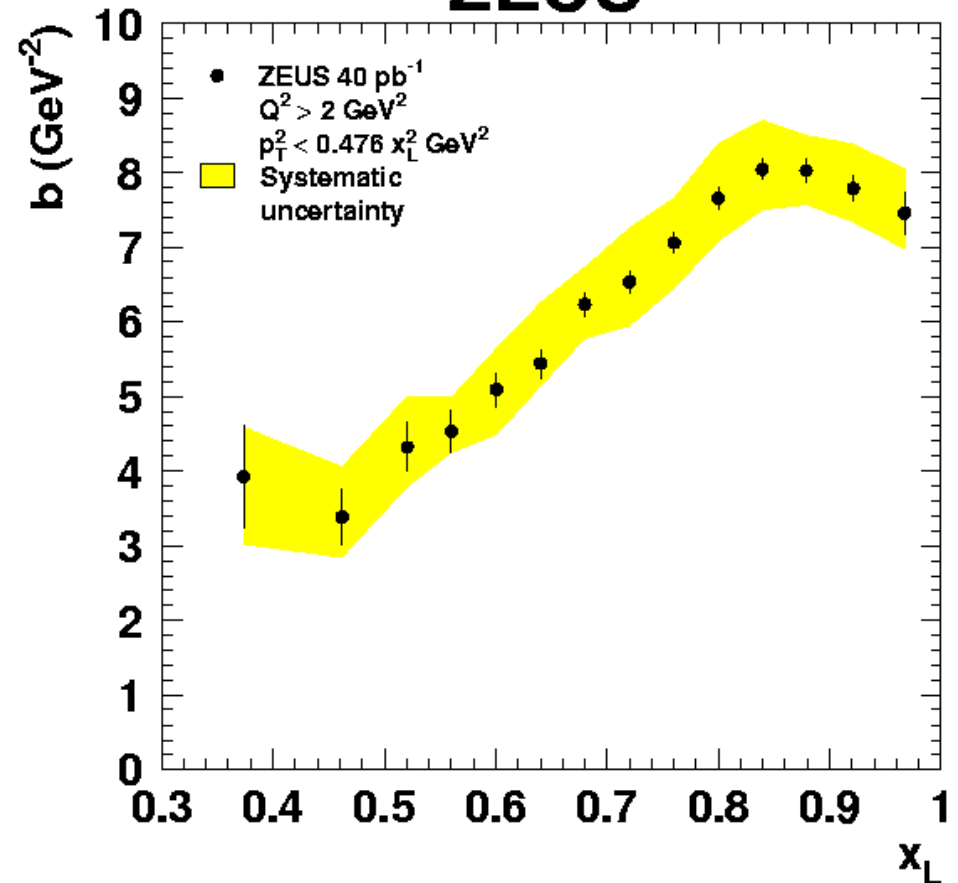
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- LN intercepts fall with x_L
- bump/plateau/shoulder $0.4 < x < 0.8$

Slopes $b(x_L)$

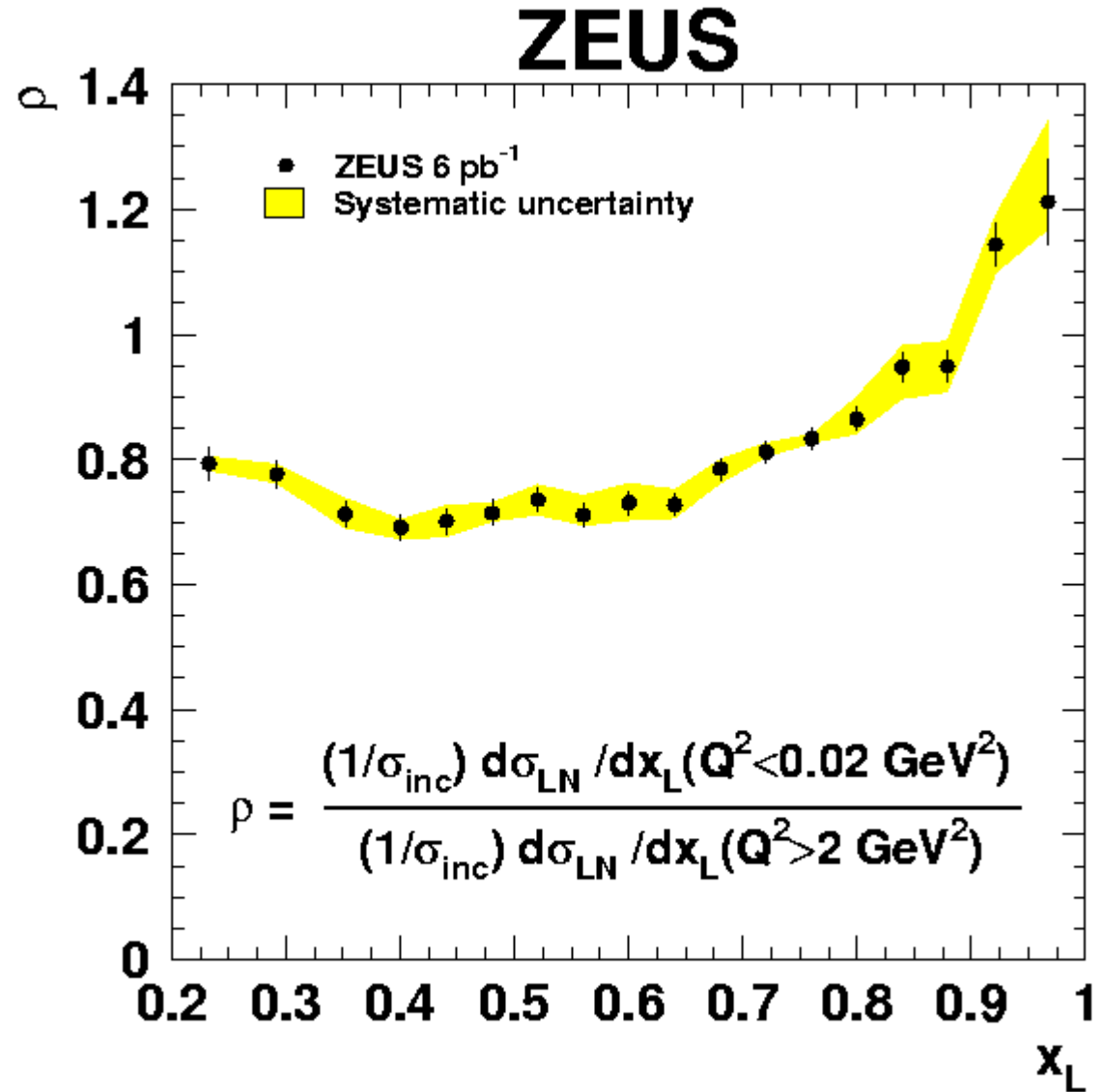
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- LN slopes sharp rise with x_L

Compare γp /DIS: x_L distributions

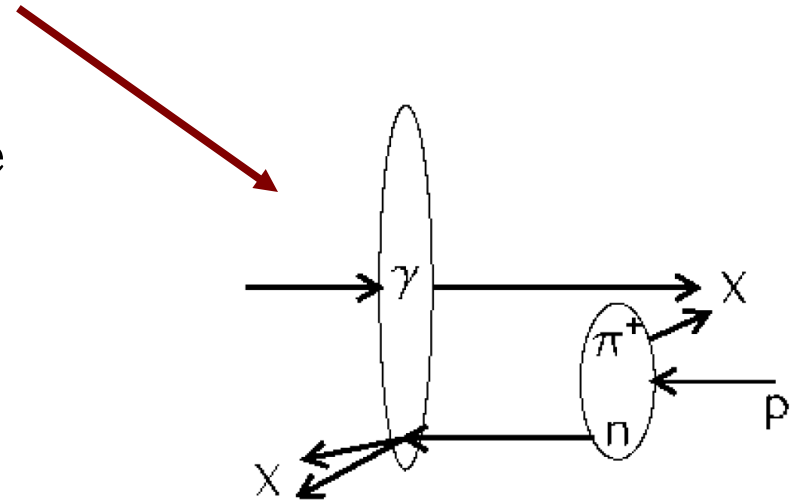
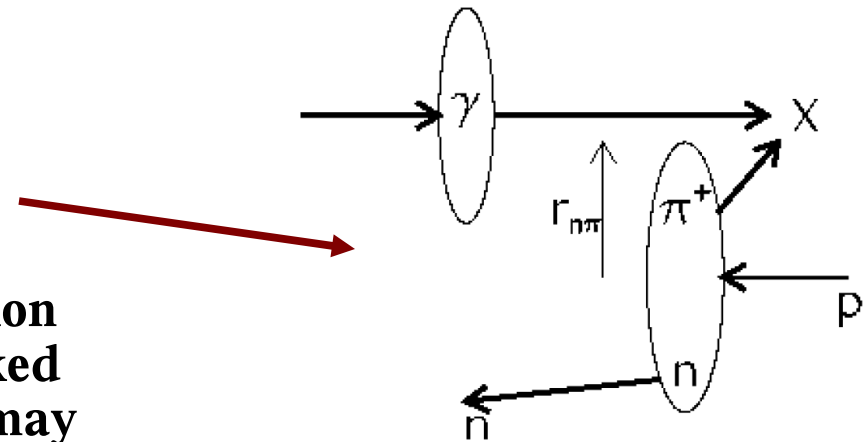
- Combine all DIS $Q^2 > 2 \text{ GeV}^2$, compare to γp x_L dist.
- Ratio $\sim 70\%$ mid- x_L , rising above 1 as $x_L \rightarrow 0.9$



Exchange model refinements: Absorptions (Rescattering)

For e.g. LN production via π -exchange:

- In DIS γ^* is small; small chance both n, π scatter on γ^* : n reaches detector
- In photoproduction γ^* large; if n - π separation smaller rescattering of n may occur: n kicked to lower x_L and higher p_T (migration) and may escape detection (absorption loss)
- In another language: multi-Pomeron exchange
- Compare photoproduction and DIS:
 - x_L, p_T distributions
 - effects of absorption?
- Effects of absorption/migration estimated:
D' Alesio, Pirner;
Nikolaev, Speth, Zakharov;
Kaidalov, Khoze, Martin, Ryskin;
Kopeliovich, Potashnikova, Schmidt, Soffer

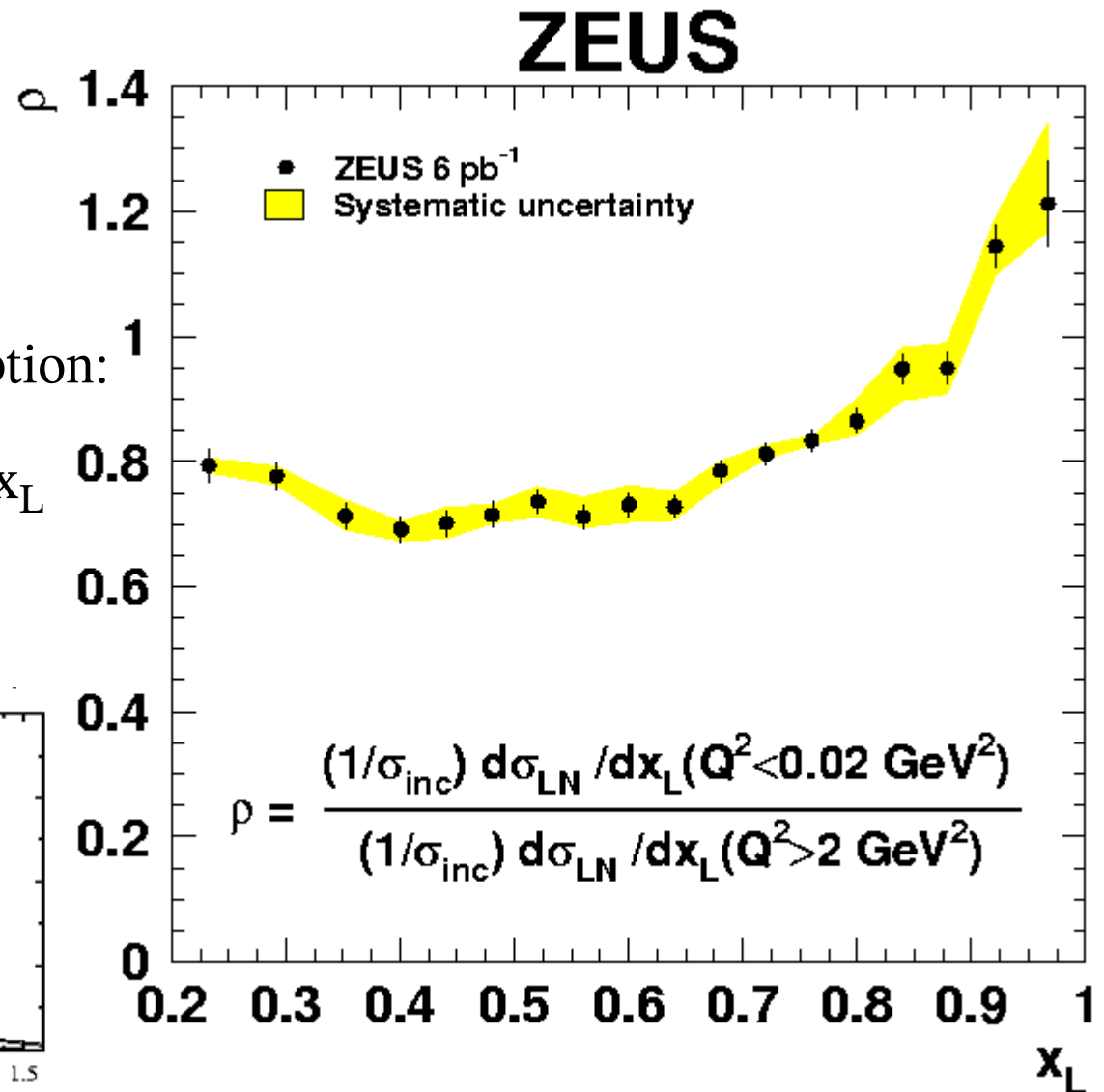
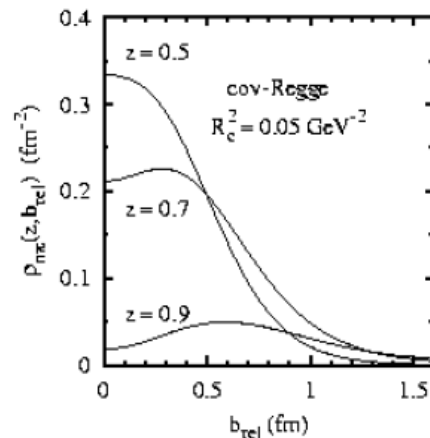
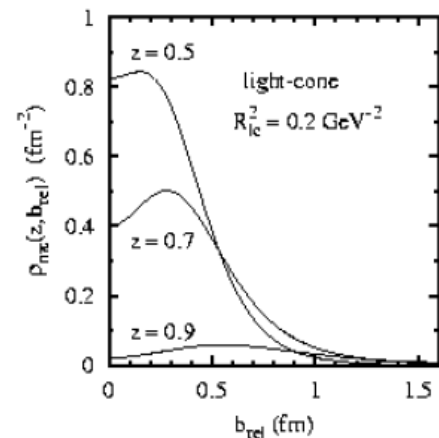


Compare γp /DIS: x_L distributions

- Combine all DIS $Q^2 > 2 \text{ GeV}^2$, compare to γp x_L dist.
- Ratio $\sim 70\%$ mid- x_L , rising above 1 as $x_L \rightarrow 0.9$

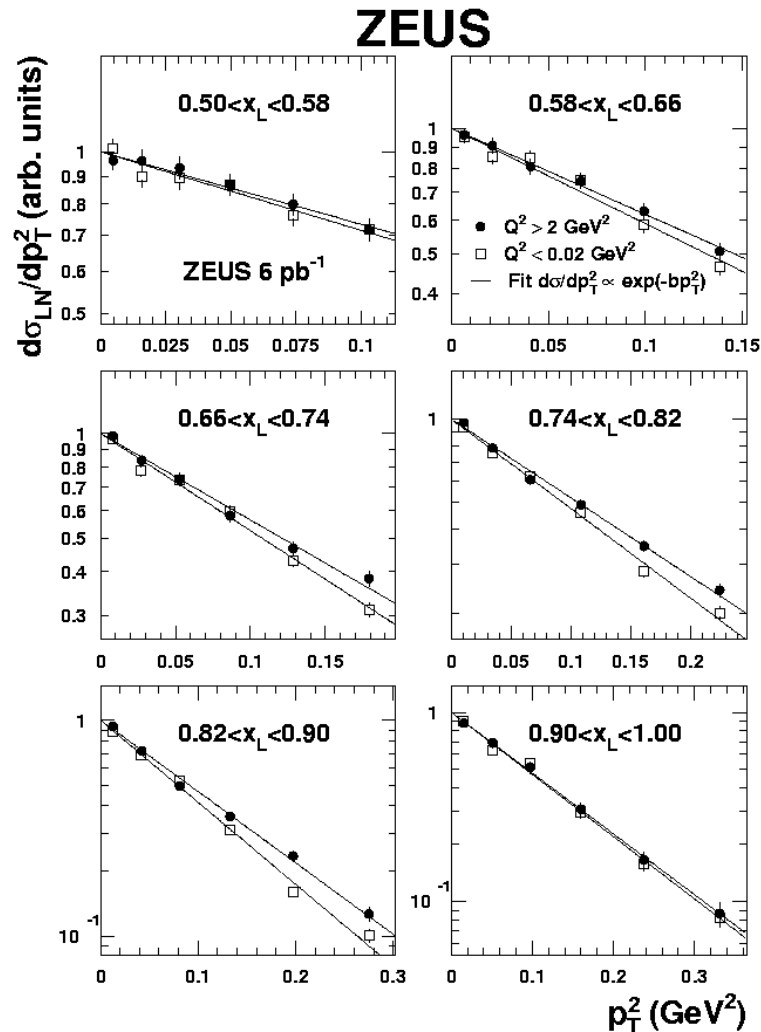
Qualitatively consistent with absorption:

- Exchange model: mean n - π separation $r_{n\pi}$ decreases at lower x_L
- Smaller $r_{n\pi} \Rightarrow$ more absorption at lower x_L

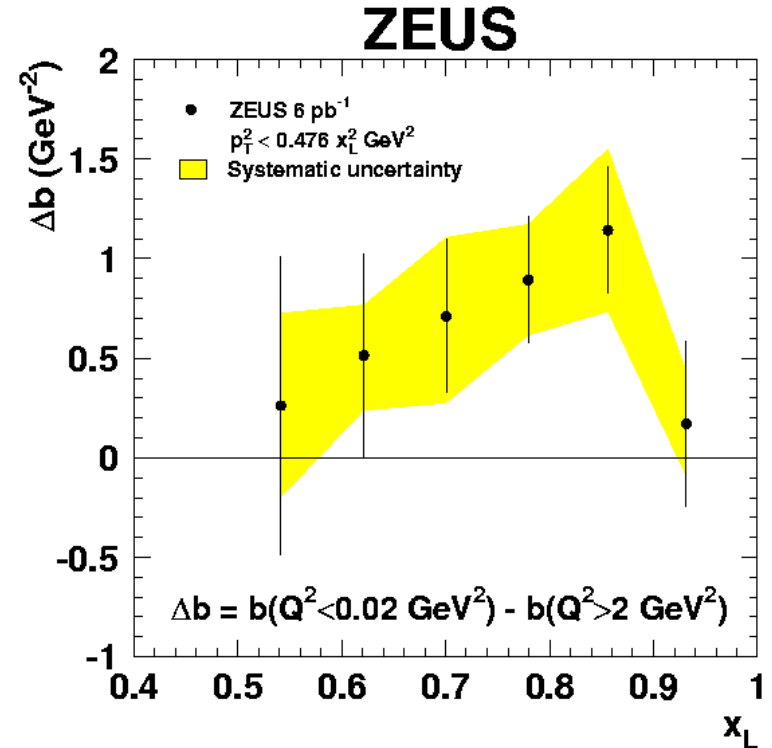


Compare γp /DIS: p_T^2 distributions

normalized at $p_T^2 = 0$

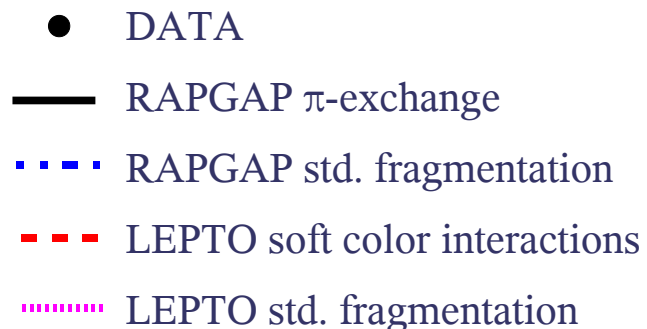


$$\Delta b = b(\gamma p) - b(\text{DIS})$$

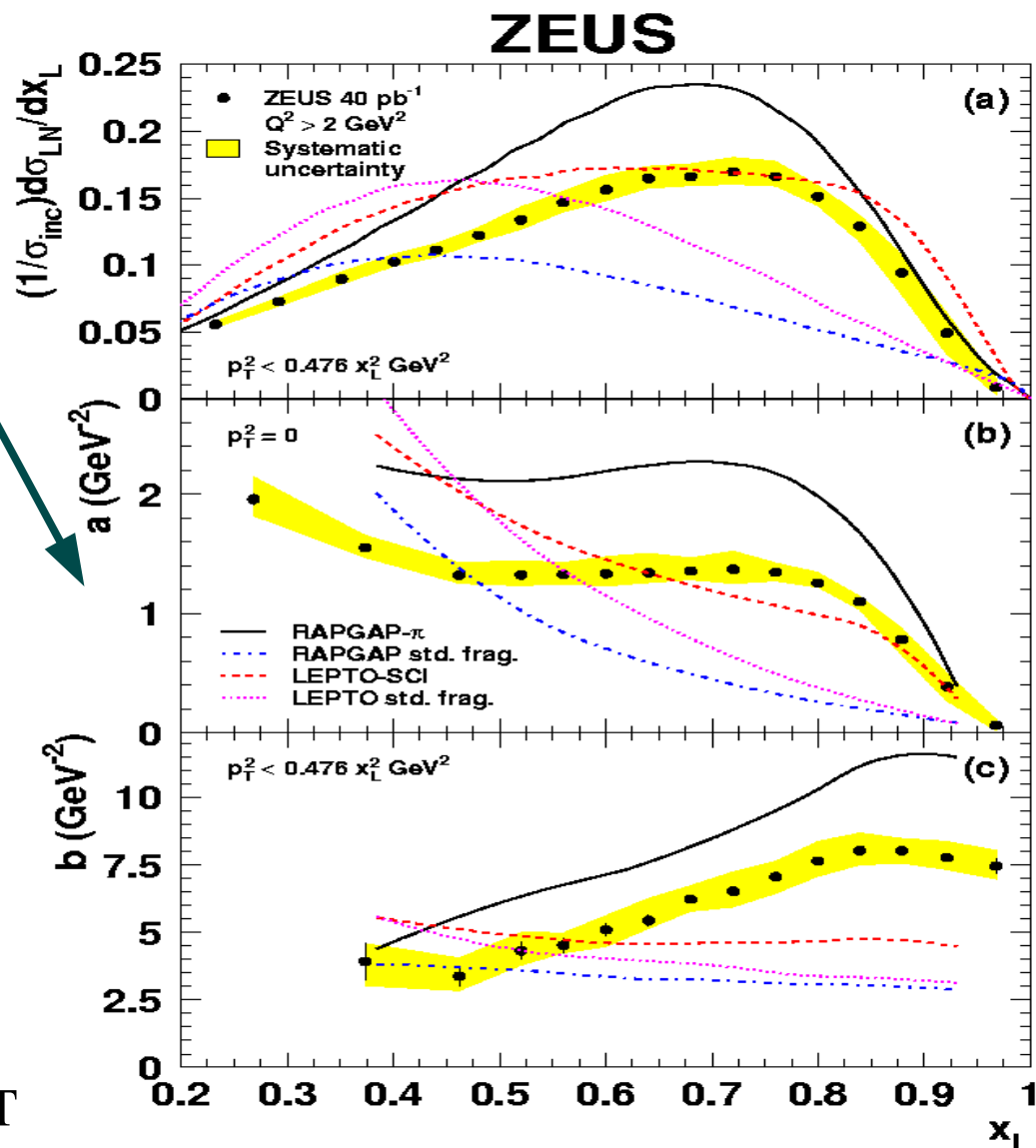


- Small but clear difference:
 $b(\gamma p) > b(\text{DIS})$ for $0.6 < x_L < 0.9$
- Qualitatively consistent with absorption:
more absorp. at small $r_{n\pi} \sim$ large p_T
fewer LN at high $p_T \Rightarrow$ larger slope

Model comparisons: DIS

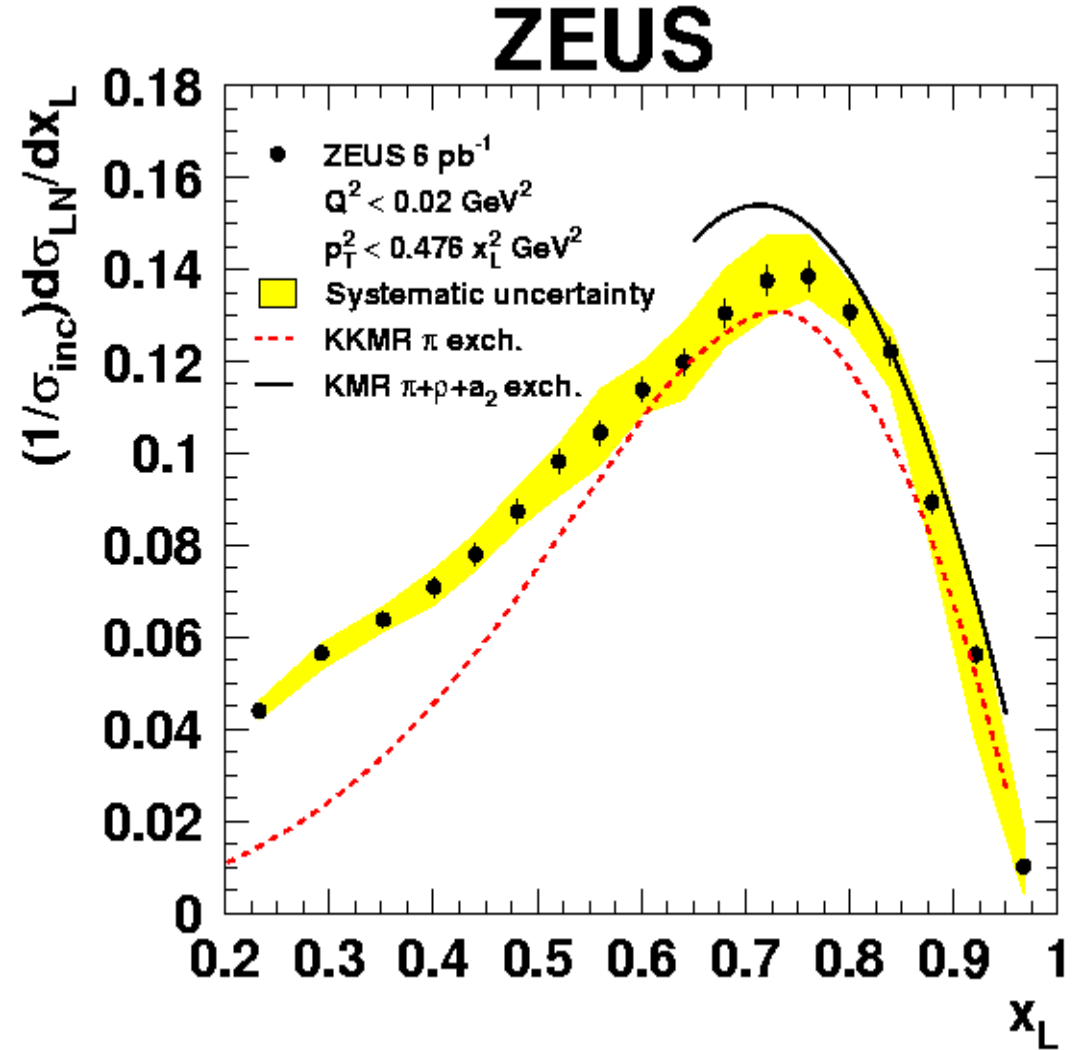


- None of models describe data well
- Both std. frag. too few n, too low x_L
- LEPTO-SCI ~ OK in shape, magnitude, but slopes too small, not x_L dependent
- RAPGAP π -exch. closest to data (but slopes too high)
- Other DIS, γp std. frag. models also fail: ARIADNE, CASCADE, PYTHIA, PHOJET

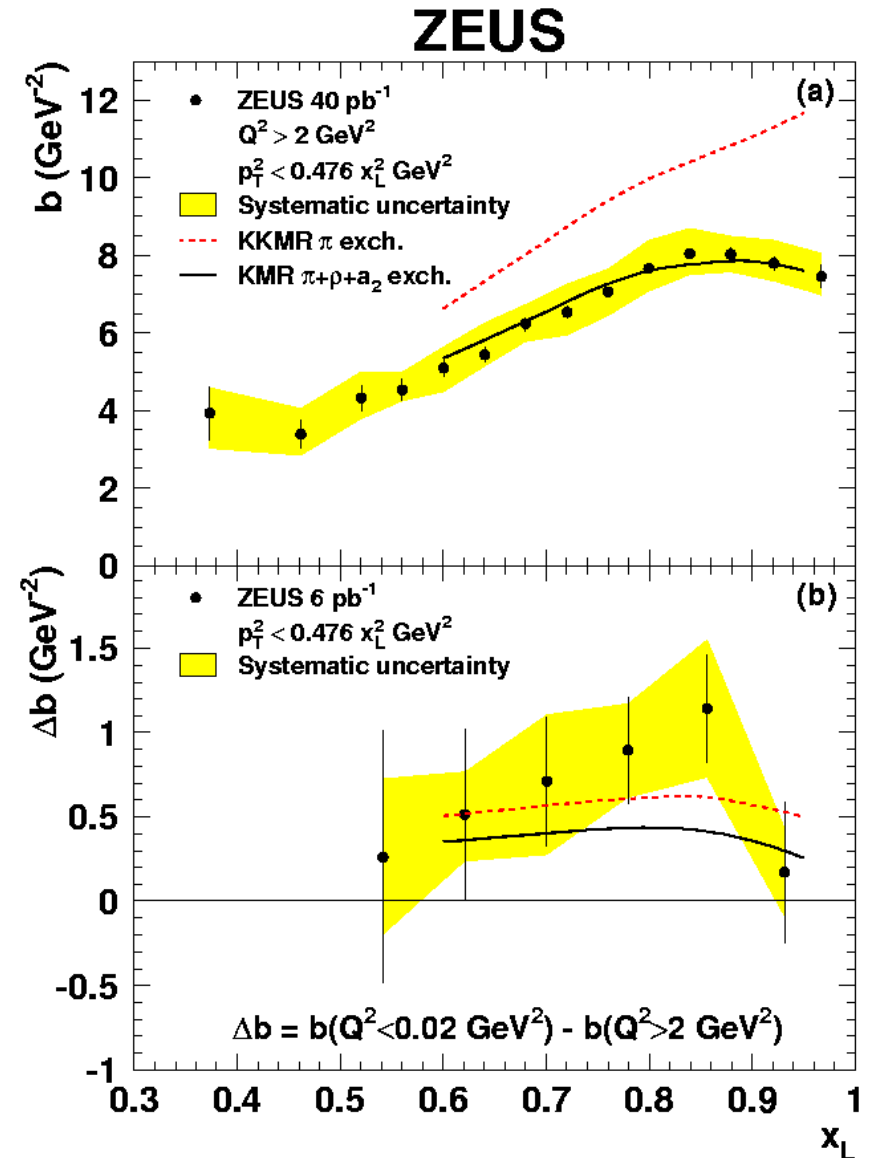


Compare: π -exch. with absorption and (ρ, a_2) exchanges

- Compare with calculations of Kaidalov, Khoze, Martin and Ryskin:
 - start with pure π -exch.
 - some n rescatter on γ
 - rescattered n migrate in (x_L, p_T)
 - add (ρ, a_2) exchanges
- Overall loss $\sim 50\%$ from pure OPE
- Reasonable agreement with LN in γp



- Absorption+migration with pion exchange alone doesn't describe slopes; too high in magnitude, no turnover at high x_L ; Δb is \sim OK
- Addition of (ρ, a_2) exchanges gives good description of both slopes magnitude and x_L dependence, Δb is still OK



Pion Structure Function from F_2^{LN}

According to π -exchange model we can extract F_2^π from measured F_2^{LN} :

$$F_2^{\text{LN}(3)}(\beta, Q^2, x_L) = \Gamma_\pi(x_L) \cdot F_2^\pi(\beta, Q^2)$$

where

$$\beta = x/(1-x_L)$$

$\Gamma_\pi(x_L)$ is integrated over t pion flux

For pion flux

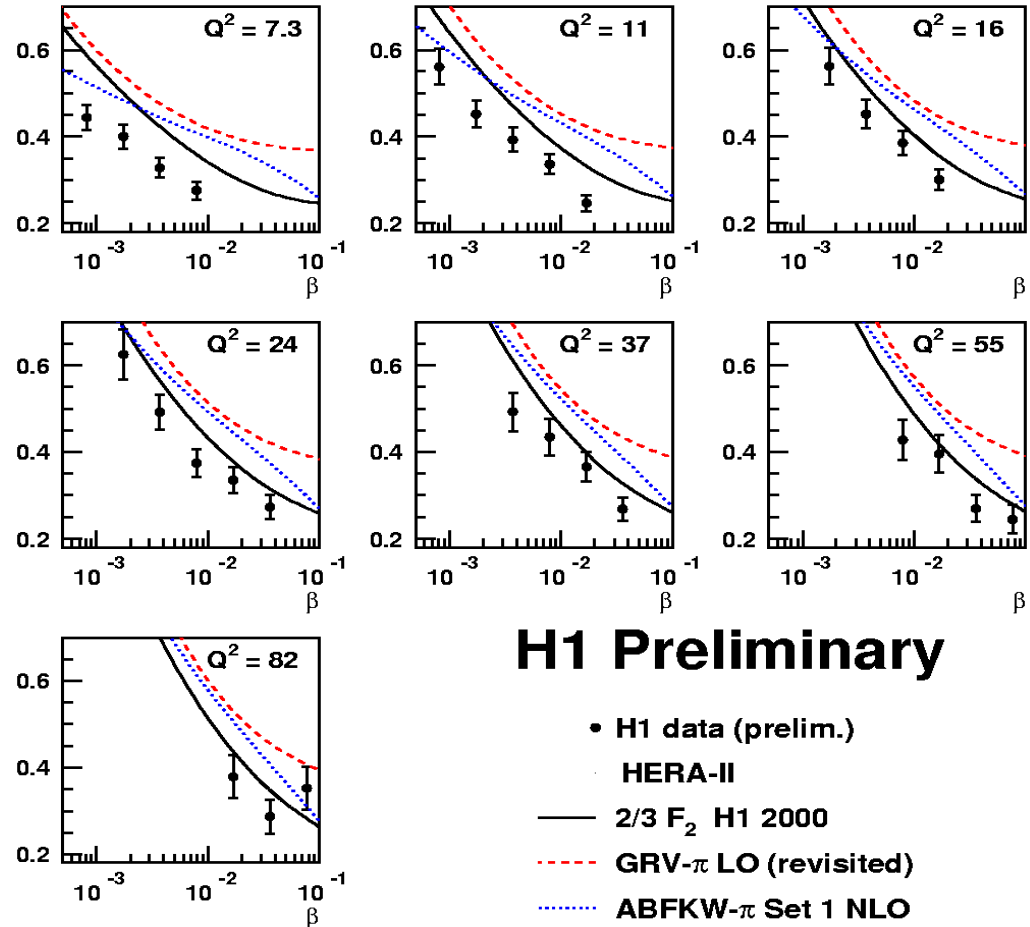
$$f_{\pi^+/p} = \frac{1}{2\pi} \frac{g_{p\pi n}^2}{4\pi} (1-x_L) \frac{-t}{(m_\pi^2 - t)^2} \exp\left(-R_{\pi n}^2 \frac{m_\pi^2 - t}{1-x_L}\right)$$

at $x_L = 0.73$ we obtained $\Gamma_\pi = 0.131$

Data compared to parameterizations:

- 2/3 of $F_2 \Rightarrow$ additive quark model
- GRV- π LO (revisited)
- ABFKW- π Set 1 NLO

$$F_2^{\text{LN}(3)}(x_L = 0.73)/\Gamma_\pi, \Gamma_\pi = 0.131$$



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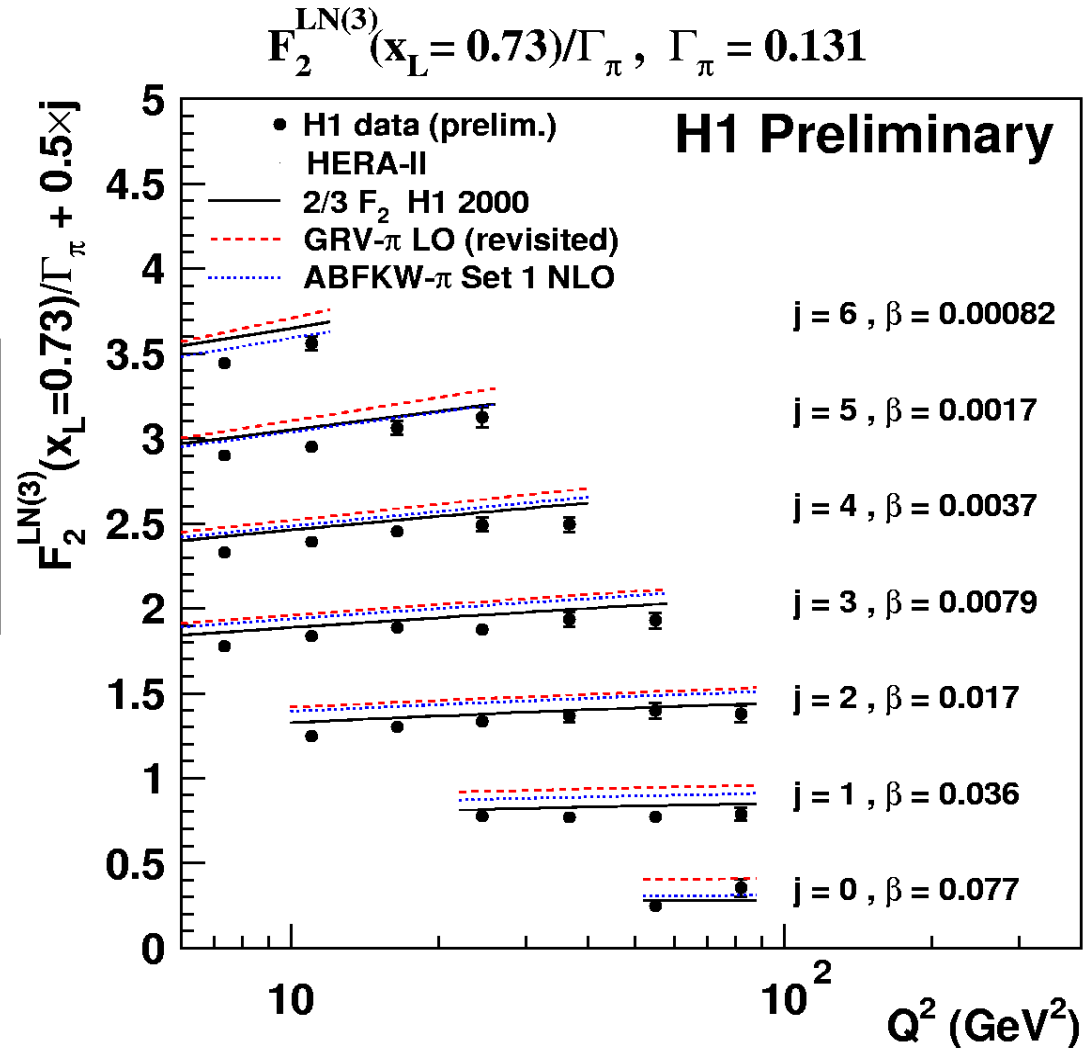
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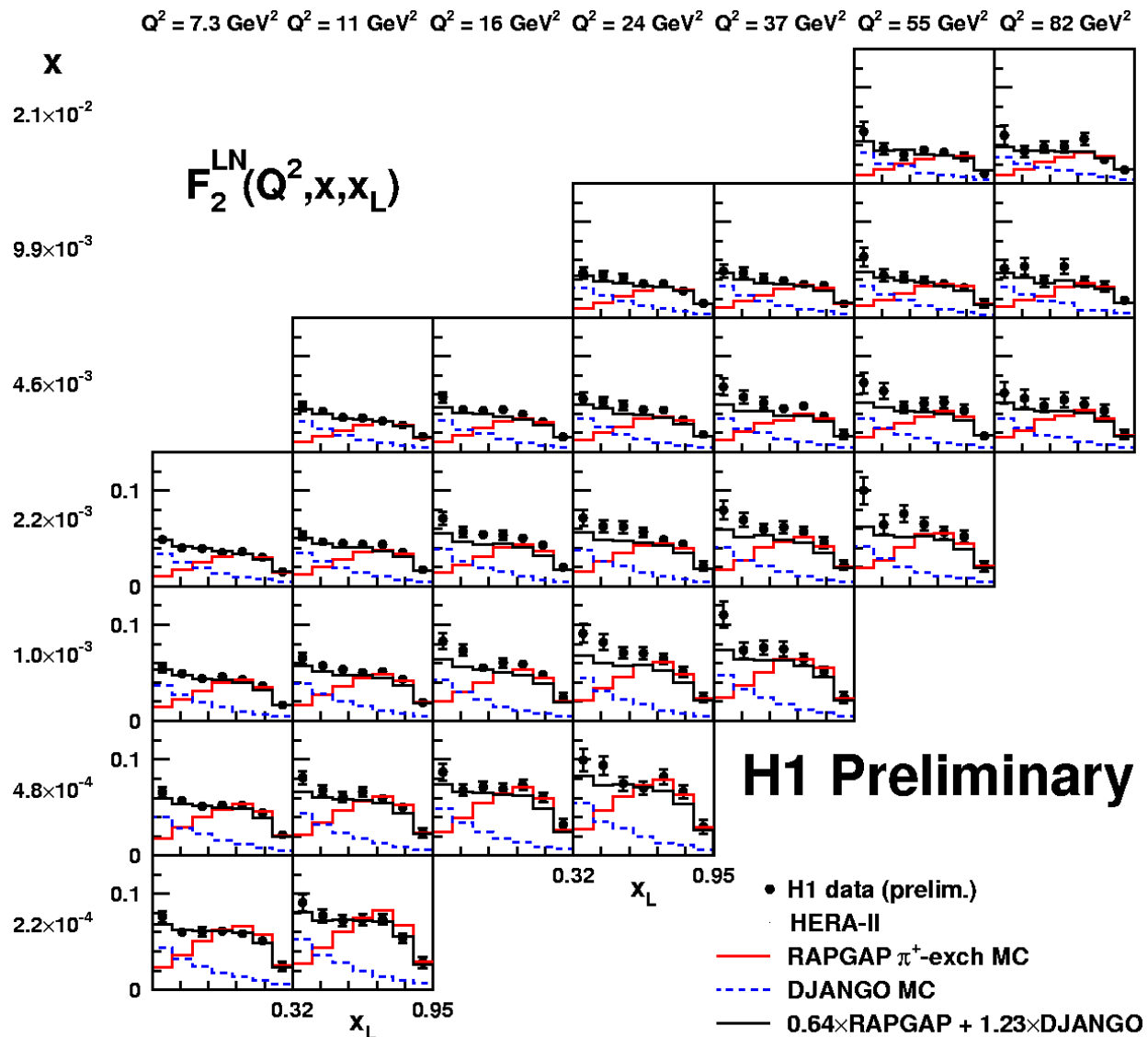
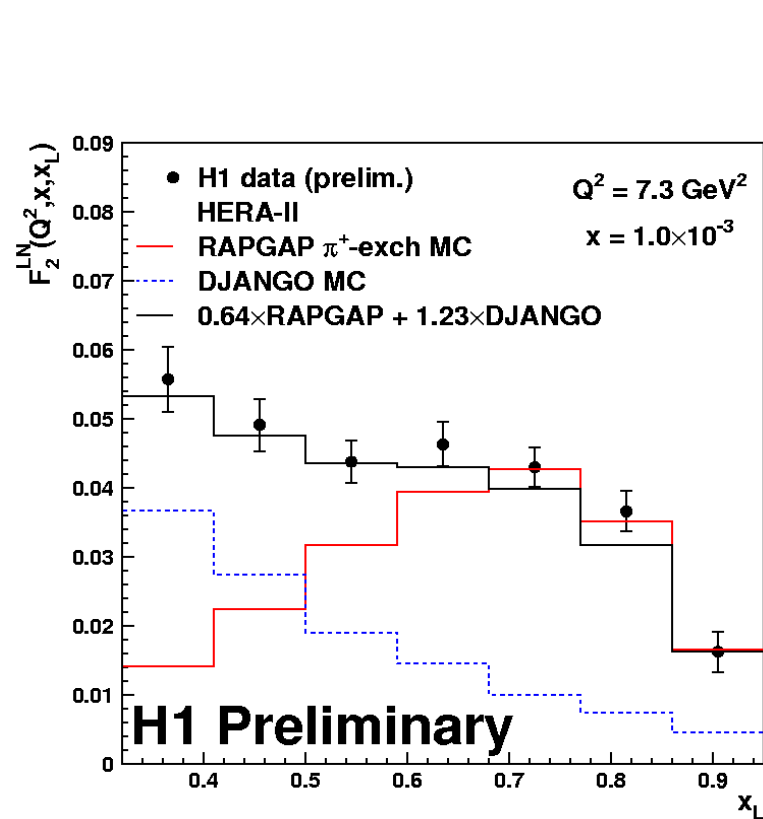
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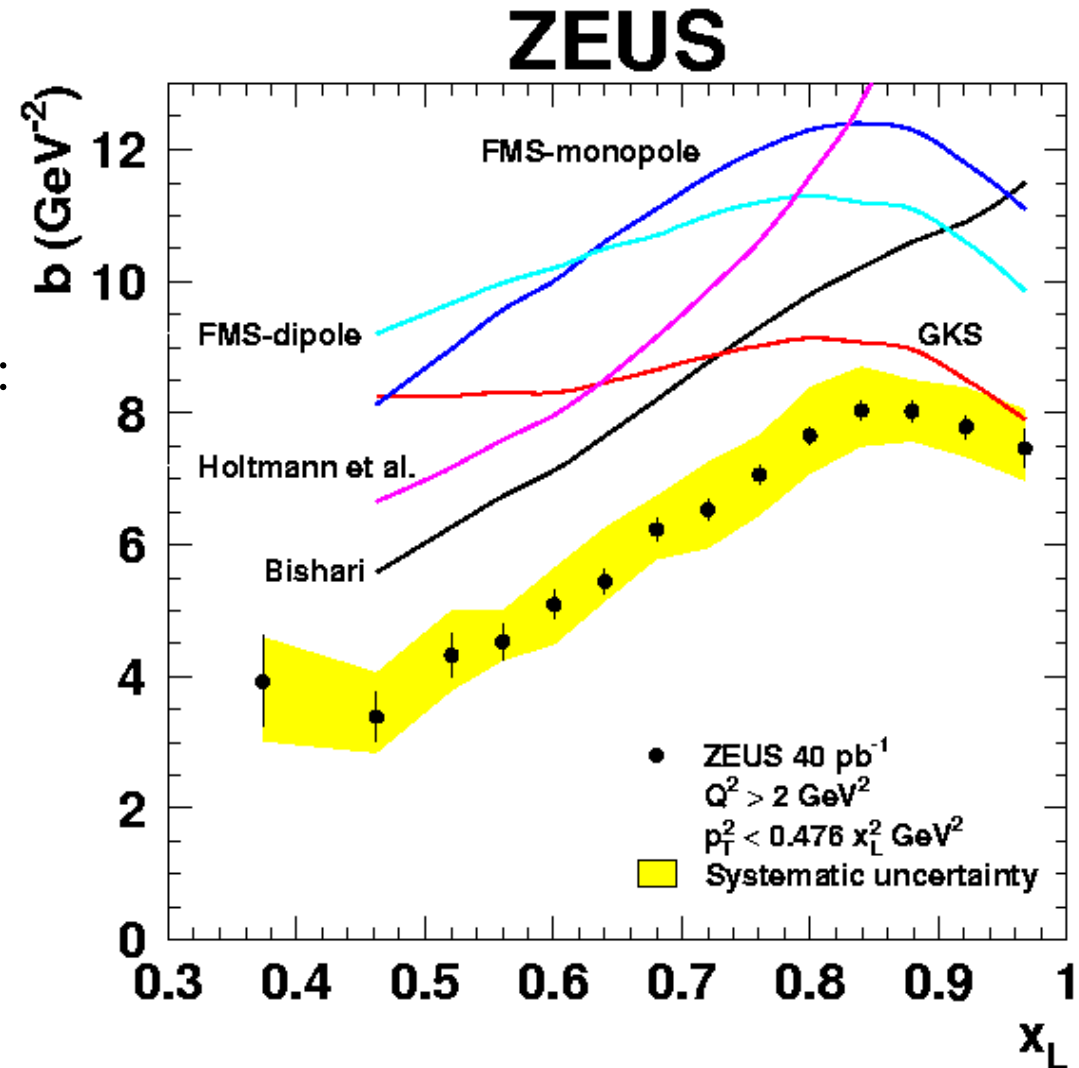


Summary

- New measurement of $F_2^{\text{LN}(3)}(Q^2, x, x_L)$ structure function from HERA-II data.
 - $\Rightarrow F_2^{\text{LN}}/F_2$ ratio is consistent with vertex factorisation
 - $\Rightarrow F_2^{\text{LN}(3)}$ is interpreted in terms of pion structure function F_2^π and compared to different parameterizations of F_2^π
- Standard fragmentation models do not describe leading neutron production
- Models with virtual particle exchange describe data much better: x_L shape and magnitude are OK but p_T slopes are still off
- Account for absorption and (ρ, a_2) exchanges \Rightarrow promising agreement with data



- Numerous parameterizations of pion flux $f_{\pi/p}(x_L, p_T)$ in literature
- Here compare to measured DIS $b(x_L)$:
- Best agreeing models shown here; others wildly off
- All give too large $b(x_L)$
- More refinement needed: absorption/migration

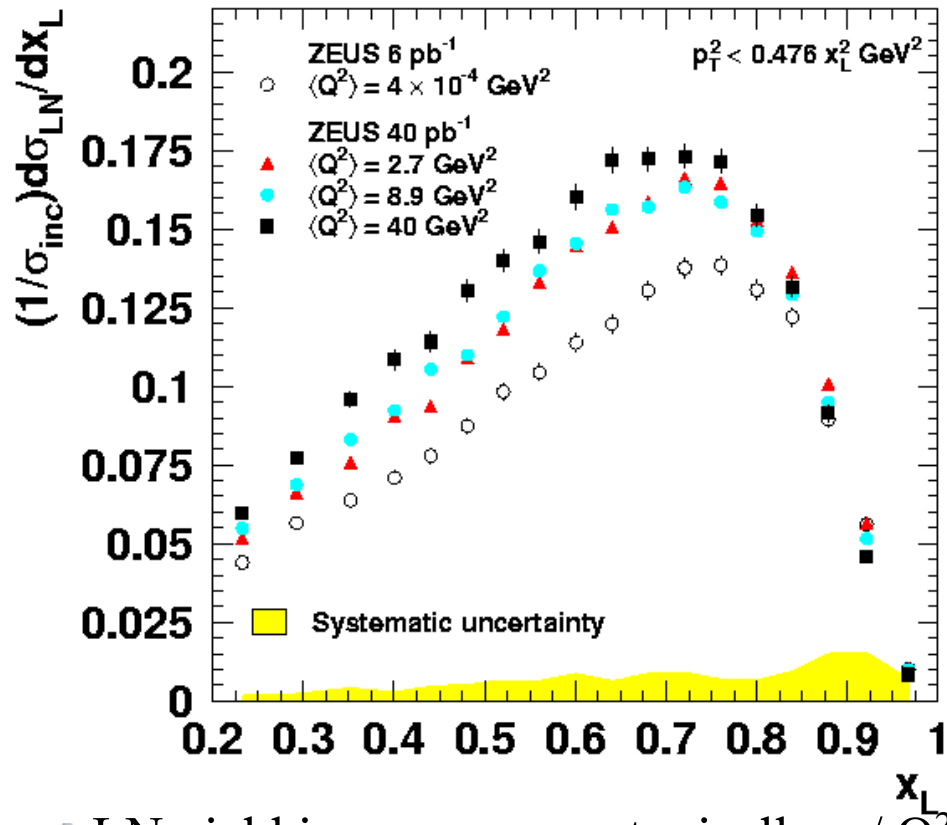


Q^2 dependence of LN production

3 Q^2 bins in DIS + 1 Q^2 bin in γp

- x_L distributions:

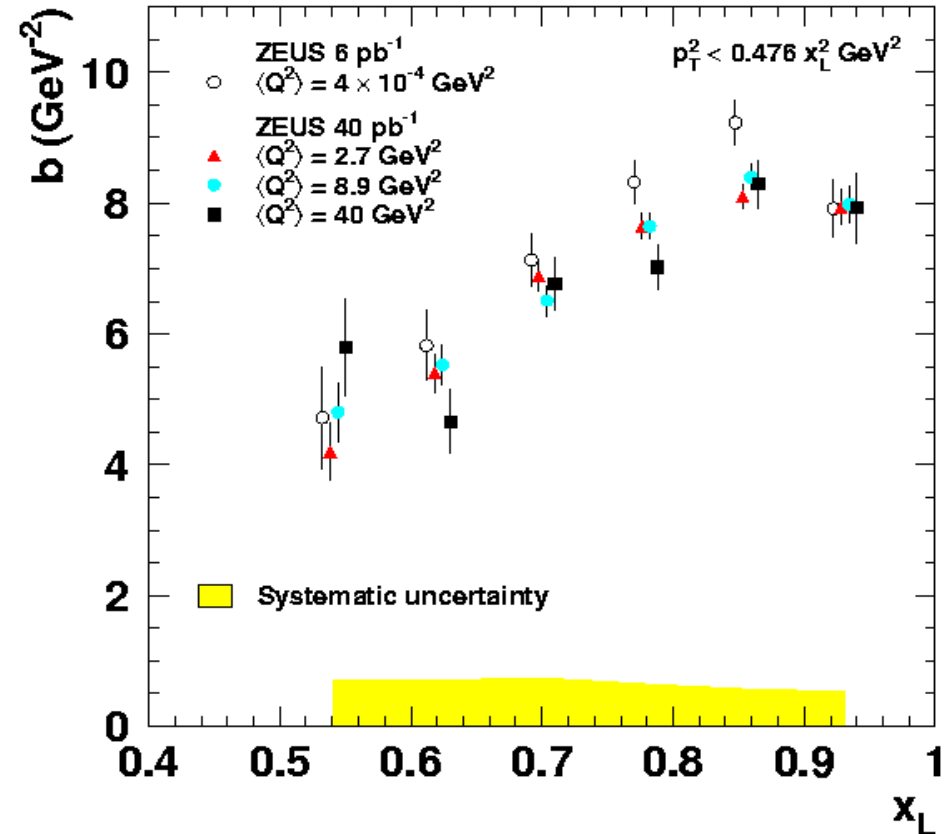
ZEUS



- LN yield increases monotonically w/ Q^2
- Consistent with absorption:
larger $Q^2 \Rightarrow$ smaller γ^* less absorption

- slopes $b(x_L)$:

ZEUS



- Slopes for 3 Q^2 approx. same
- Slope for γp significantly larger

