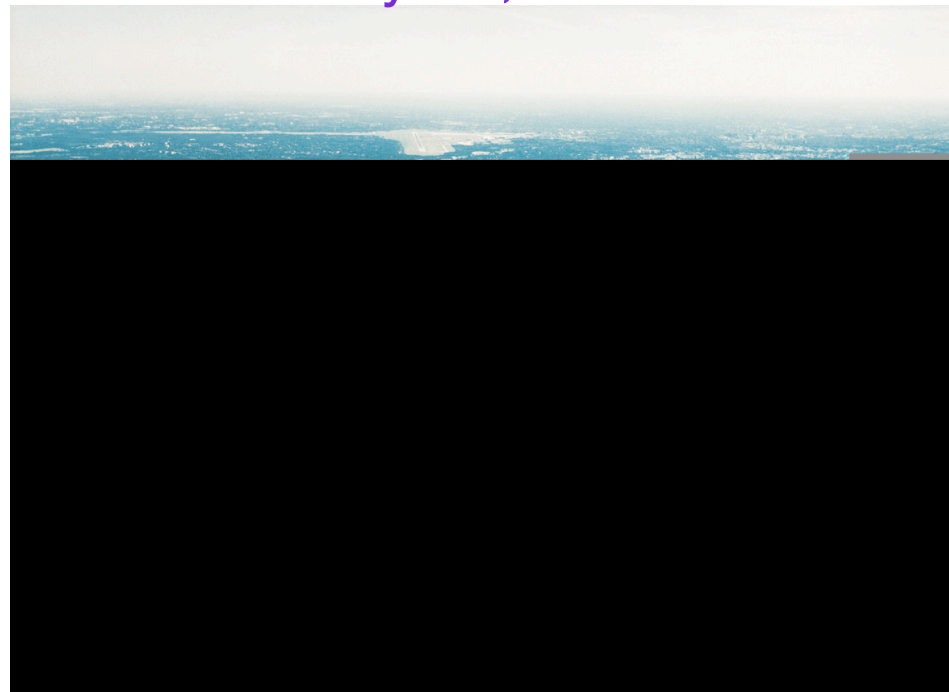


Electroweak physics at HERA

Michele Rosin
on behalf of the ZEUS and H1 collaborations
July 18, 2010



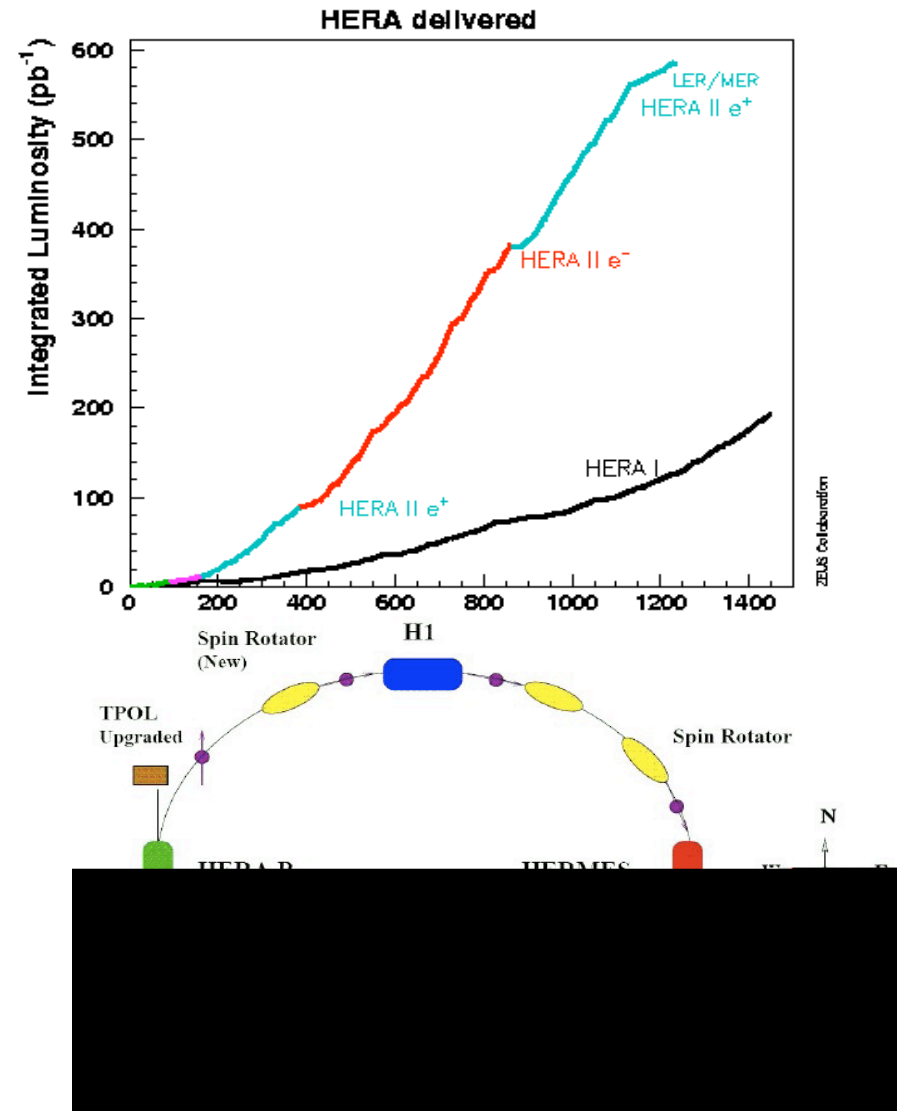
HERA ep collider (1992-2007)

- $E_{e\pm} = 27.5 \text{ GeV}$
- 1992-2000 (HERA I) $E_p = 820,920 \text{ GeV}$
- 2003-2007 (HERA II) $E_p = 920 \text{ GeV}$ $\sqrt{s} = 318 \text{ GeV}$

- Increased luminosity
- Longitudinal polarization of lepton beam

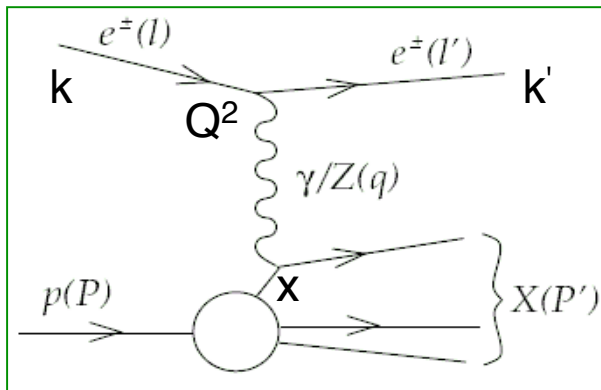
$$P_e = \frac{N_{RH} - N_{LH}}{N_{RH} + N_{LH}} \sim 30 - 40\%$$

- For colliding beam experiments:
two large multipurpose detectors H1 and ZEUS
- $\sim 0.5 \text{ fb}^{-1}$ recorded luminosity for each experiment

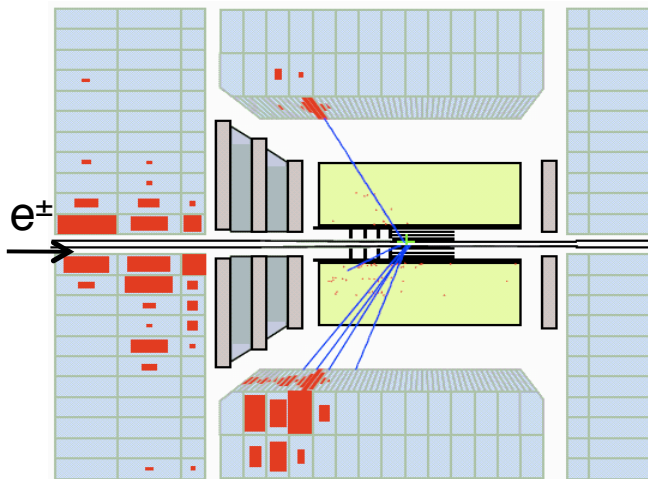


Deep Inelastic Scattering (DIS)

Neutral Current (NC)



ZEUS



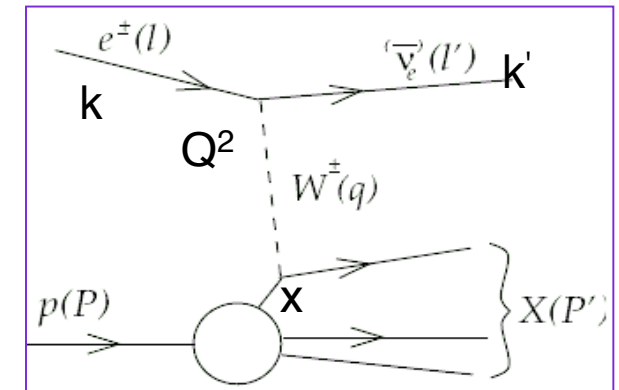
Boson virtuality

$$Q^2 = -q^2 = -(l - l')^2 = -(\vec{k} - \vec{k}')^2 + (l - l')^2$$

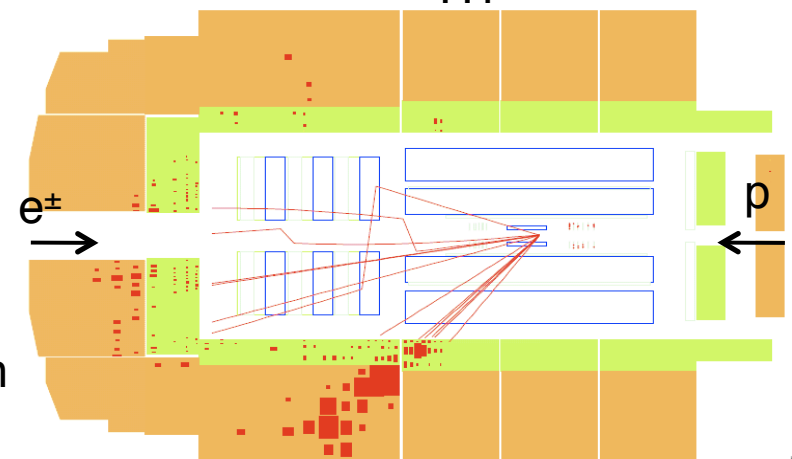
Bjorken x

inelasticity

Charged Current (CC)



H1



Explore proton structure
& quark-gluon interaction
dynamics

NC cross section

$$\frac{d^2\sigma^{NC}(e^\pm p)}{dx dQ^2} = \frac{2\pi\alpha^2}{xQ^4} \left[Y_+ \tilde{F}_2^\mp \mp Y_- x \tilde{F}_3^\pm - y^2 \tilde{F}_L^\pm \right]$$

dominant contribution
important at high Q^2
sizeable at high y in QPM $\tilde{F}_L^\pm = 0$

$$\kappa = \frac{1}{4 \sin^2 \theta_w \cos^2 \theta_w} \frac{Q^2}{Q^2 + M_Z^2}$$

$$Y_\pm = 1 \pm (1 - y)^2$$

generalized structure functions:

$$\tilde{F}_2^\pm = F_2^\gamma + \kappa(-v_e \pm P_e a_e) F_2^{\gamma Z} + \kappa^2(v_e^2 + a_e^2 \pm 2P_e v_e a_e) F_2^Z$$

$$P_e \quad x F_3^{\gamma Z} \quad P_e \quad x F_3^Z$$

Polarization dependence due to γZ interference and Z terms
with opposite signs for positron and electron

$$[F_2^\gamma, F_2^{\gamma Z}, F_2^Z] = \sum_q [e_q^2, 2e_q v_q, v_q^2 + a_q^2] x(q + \bar{q})$$

F_2 functions are related to the sum of quark and anti-quarks

$$[x F_3^{\gamma Z}, x F_3^Z] = \sum_q [e_q a_q, v_q a_q] 2x(q - \bar{q})$$

$x F_3$ functions are related to difference bet. quark and anti-quarks
→ sensitivity to valence quarks

CC cross section

$$\frac{d^2\sigma^{CC}(e^\pm p)}{dx dQ^2} = (1 \pm P_e) \frac{G_F^2}{4\pi x} \left(\frac{M_W^2}{M_W^2 + Q^2} \right)^2 \tilde{\sigma}_{CC}^{e^\pm p}$$

CC reduced cross section

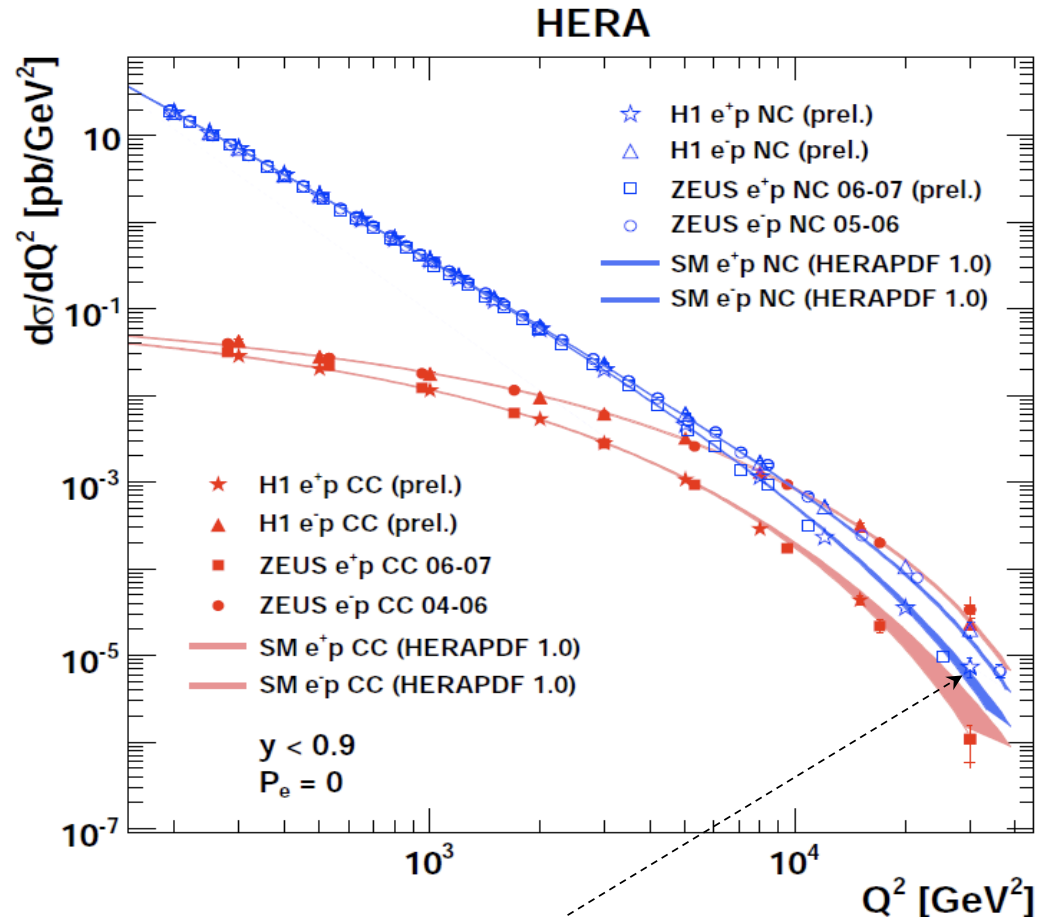
e^+/e^- sensitive to different quark densities:

$$\tilde{\sigma}_{CC}^{e^+ p} = x [\bar{u} + \bar{c}] + (1 - y)^2 x [d + s]$$

in L.O:

CC gives sensitivity to different combinations of quarks as NC.

Electroweak unification



difference in e⁺ and e⁻ for NC in high Q² region
comes from contribution of Z exchange

$$\text{NC: } \frac{d\sigma}{dQ^2} \sim \frac{1}{Q^4}$$

$$\text{CC: } \frac{d\sigma}{dQ^2} \sim \frac{1}{(Q^2 + M_W^2)^2}$$

EW component of SM:

NC and CC cross sections become similar at $Q^2 \approx M_Z^2, M_W^2$

Data compared with SM (HERAPDF 1.0) → described later

Good agreement over full range

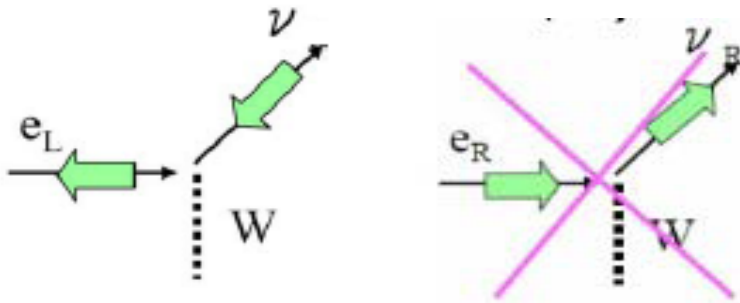
Total CC cross sections for e^+p & e^-p as function of polarization

Linear dependence of σ^{CC} on P_e

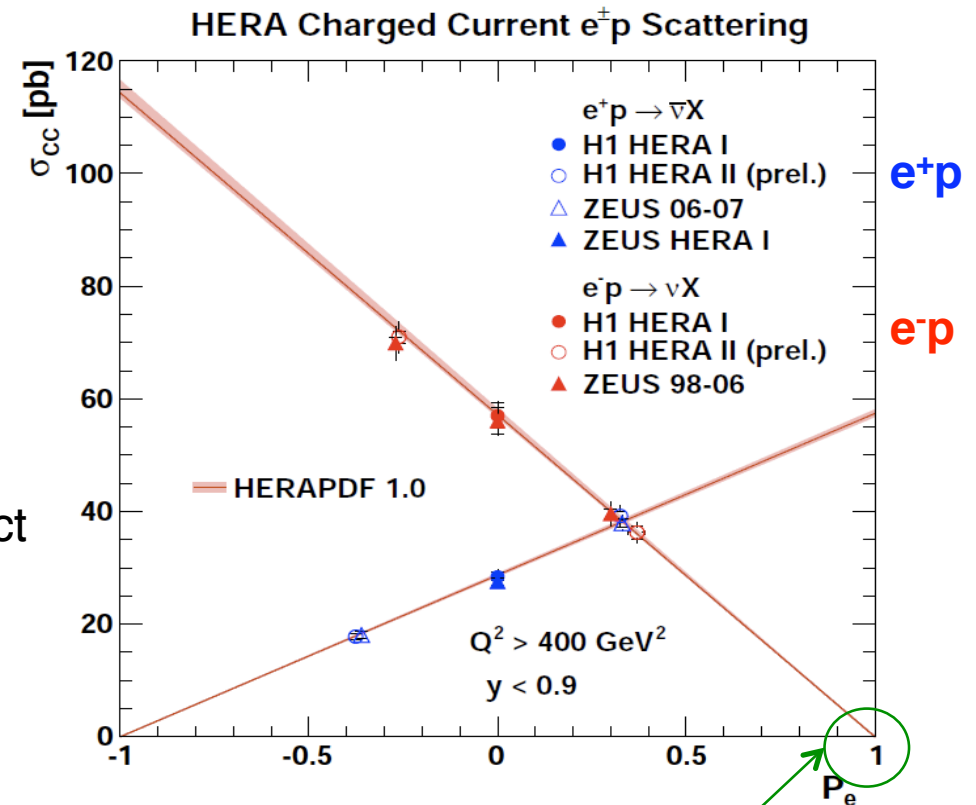
$$\sigma^{CC}(e^\pm p) = (1 \pm P_e) \sigma_{P_e=0}^{CC}(e^\pm p)$$

$$P_e = \frac{N_{RH} - N_{LH}}{N_{RH} + N_{LH}}$$

SM: weak CC interactions:
only left handed particles (anti-particles) interact
→ right (left) handed currents forbidden



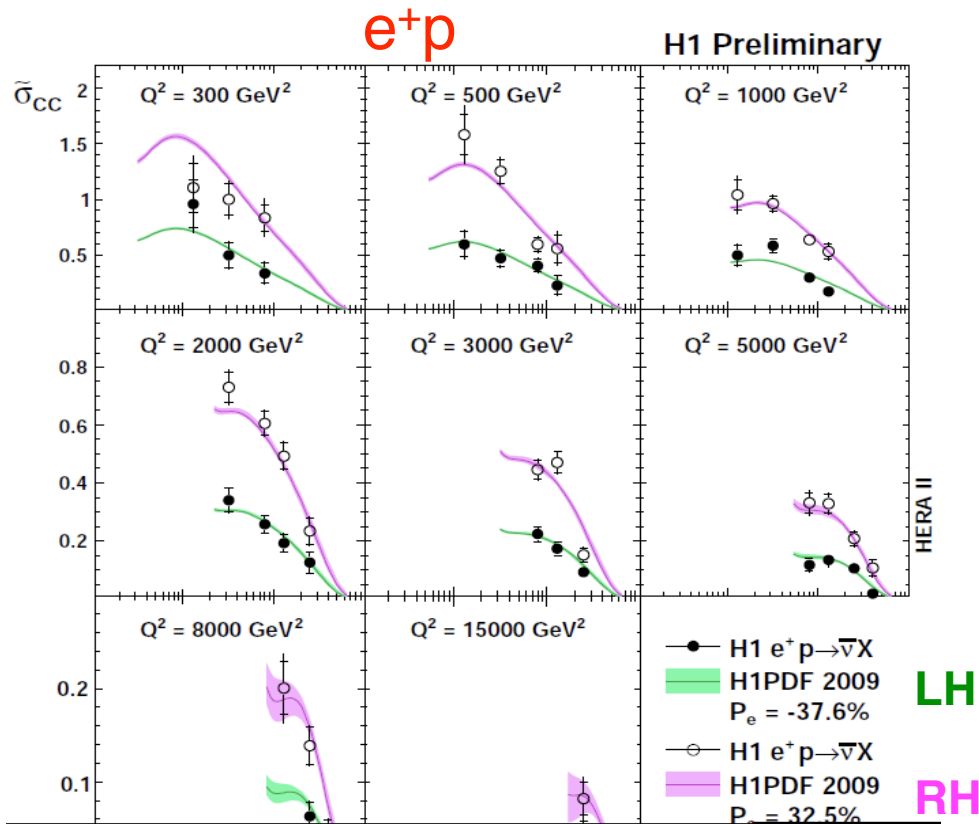
ZEUS and H1 in agreement with SM



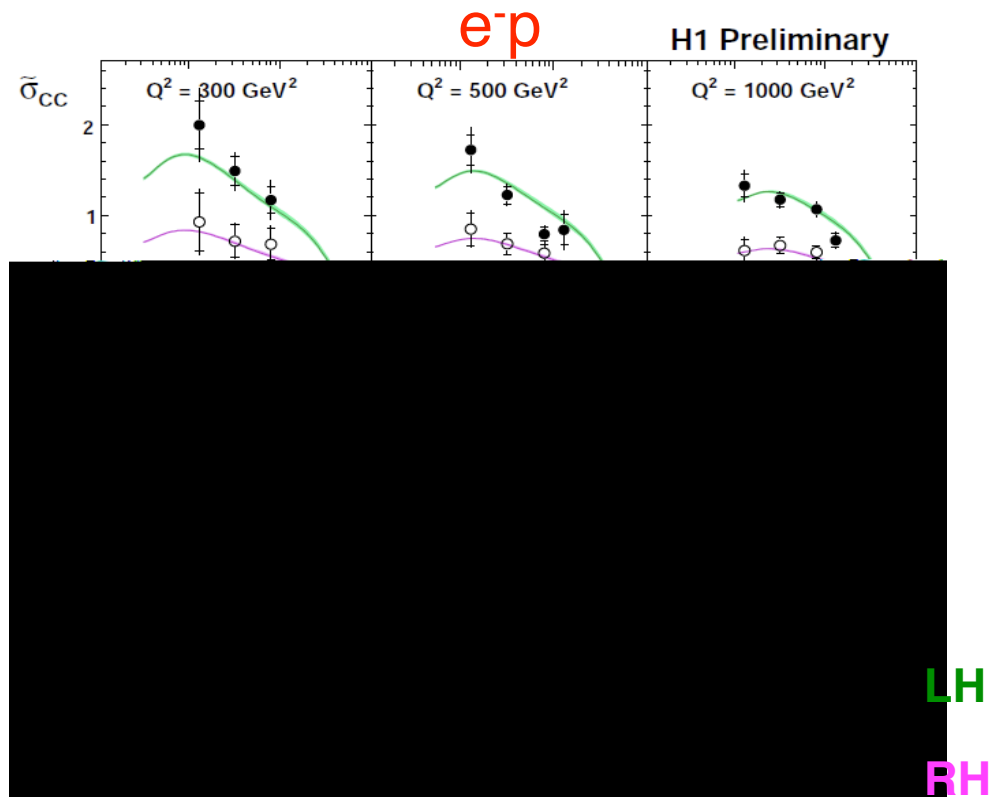
Absence of
right-handed
weak current in SM

Polarized CC reduced cross section measurements

double differential CC reduced cross sections as a function of x in Q^2 bins



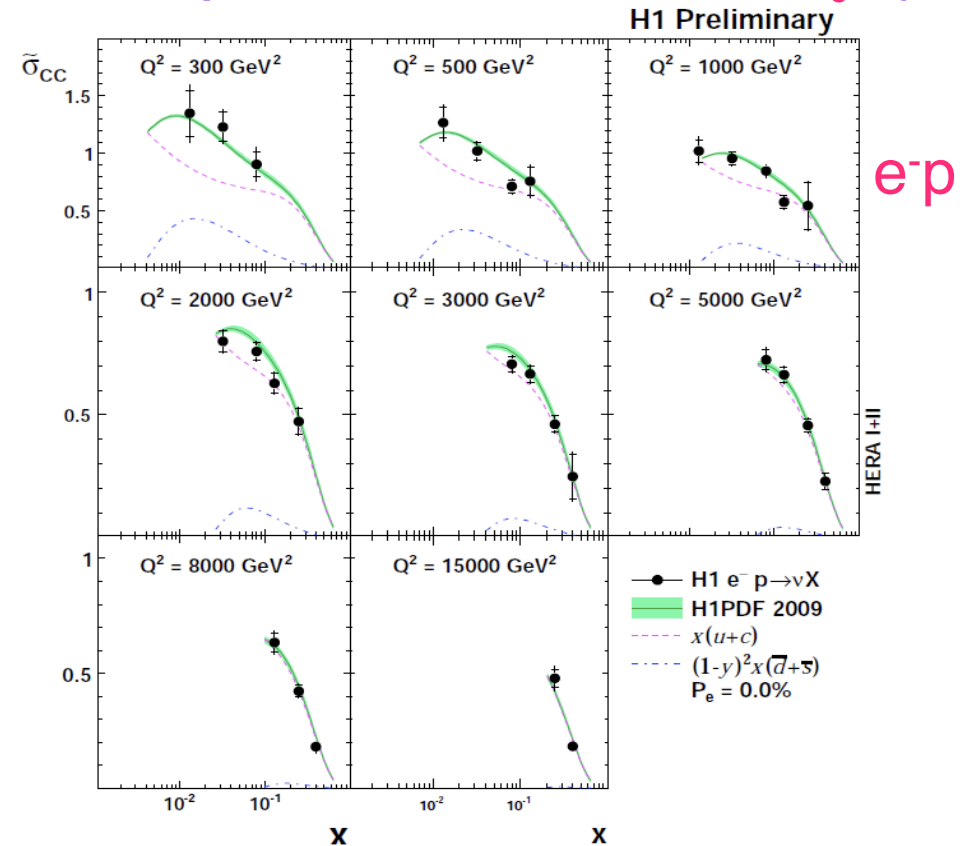
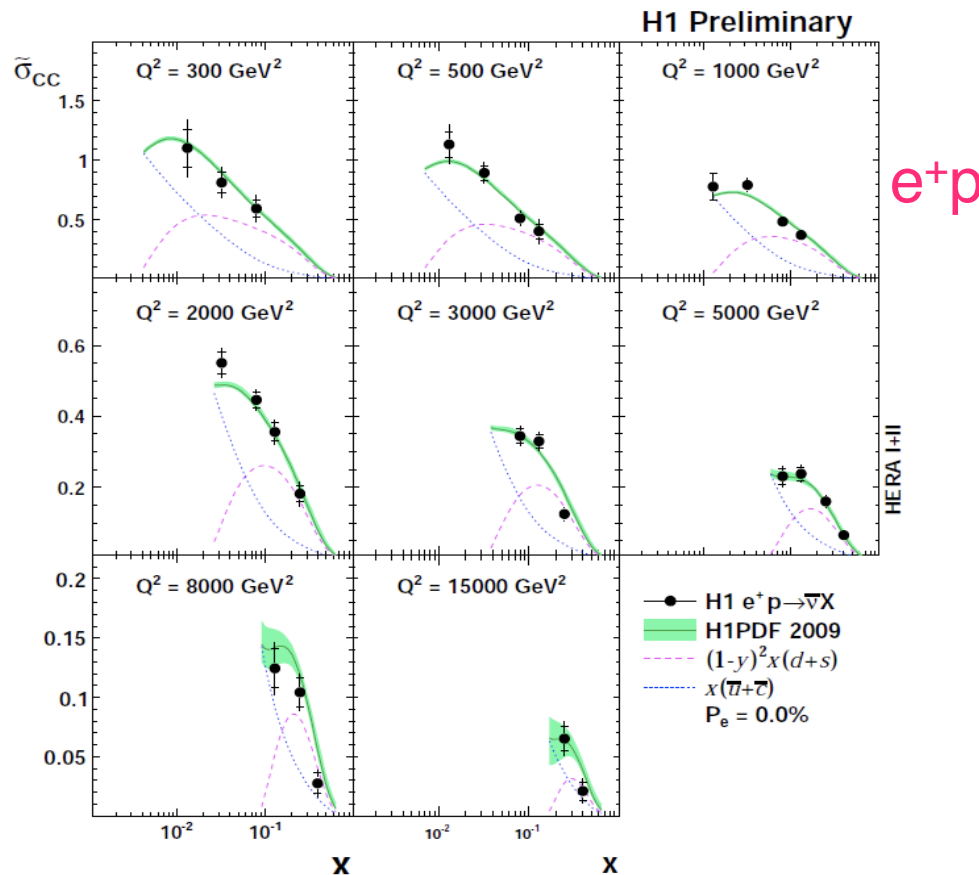
Data points shown sep. for + and - polarizations



Predictions of SM give good description of data

CC reduced cross section measurements

Data points shown for entire data set (both + and – polarizations, corrected to $P_e=0$)



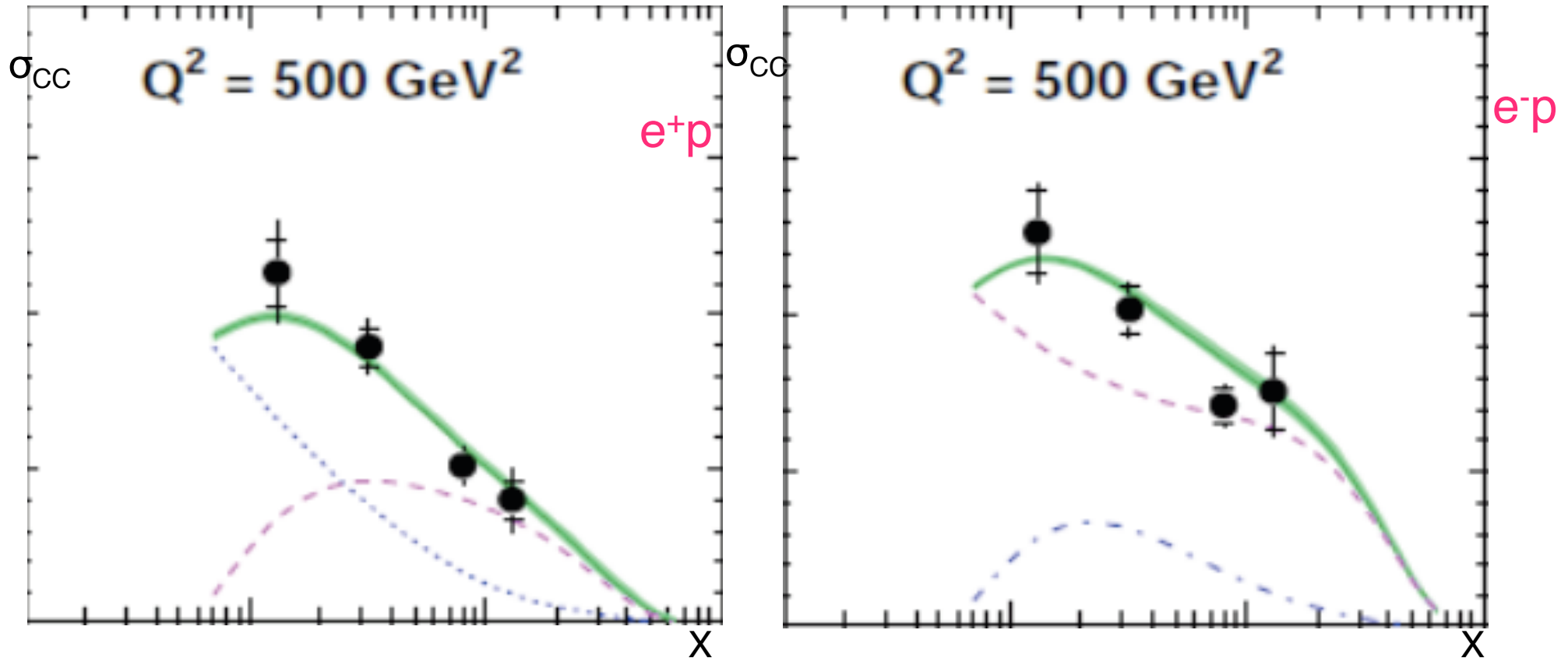
$$(1-y)^2 x(d+s) \quad x(\bar{u} + \bar{c}) \quad (1-y)^2 x(\bar{d} + \bar{s}) \quad x(u+c)$$

Blue and magenta dashed lines show the contributions of the PDF combinations

Sensitivity to different combinations of quarks than in NC

CC reduced cross section measurements

Data points shown for entire data set (both + and – polarizations, corrected to $P_e=0$)

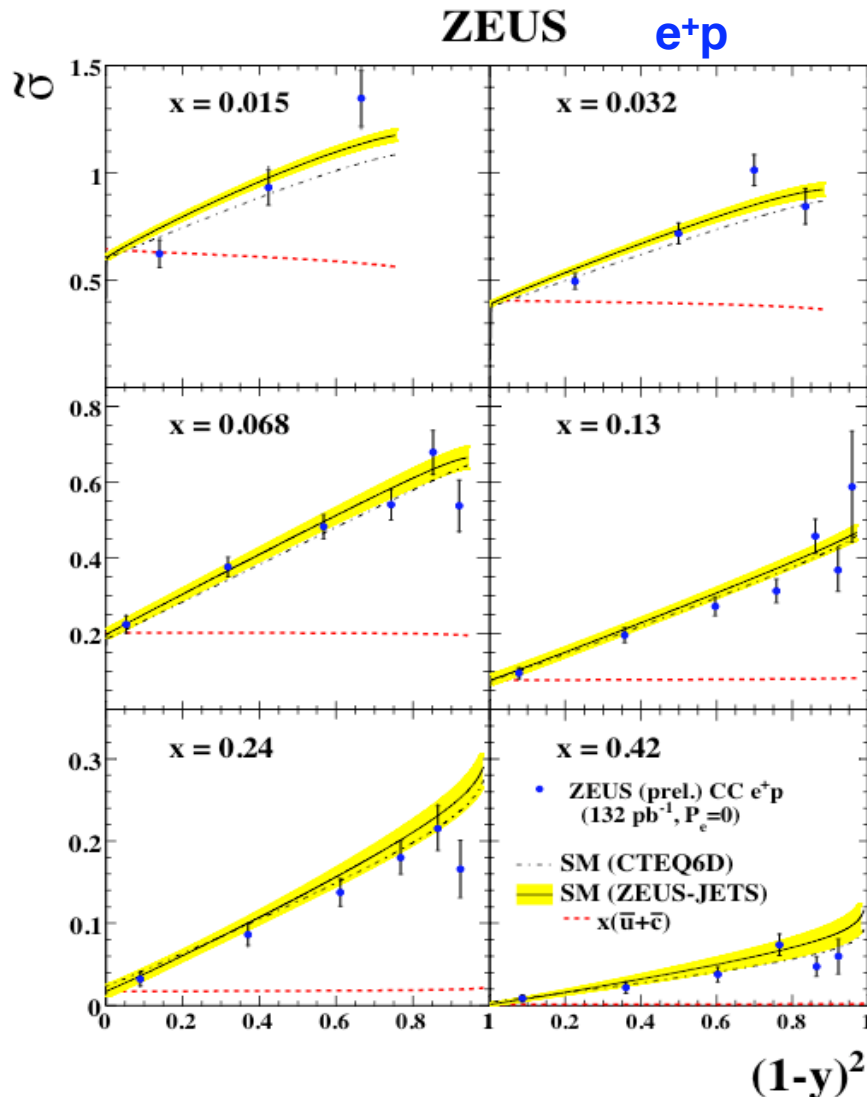


$$(1-y)^2 x(d+s) \quad x(\bar{u} + \bar{c}) \quad (1-y)^2 x(\bar{d} + \bar{s}) \quad x(u+c)$$

Blue and magenta dashed lines show the contributions of the PDF combinations

Sensitivity to different combinations of quarks than in NC

Unpolarized CC measurements



→ different way to look at sensitivity to these diff. quark species

Unpolarized reduced CC e^+p cross section in bins of x as a function of $(1-y)^2$
→ Helicity structure of CC

At leading order in QCD:

$$\tilde{\sigma}_{CC}^{e^+p} = x \underbrace{[\bar{u} + \bar{c}]}_{\text{intercept}} + (1-y)^2 x \underbrace{[d + s]}_{\text{slope}}$$

Data are well described by SM calculations

- ZEUS-JETS fit
- global fits CTEQ6

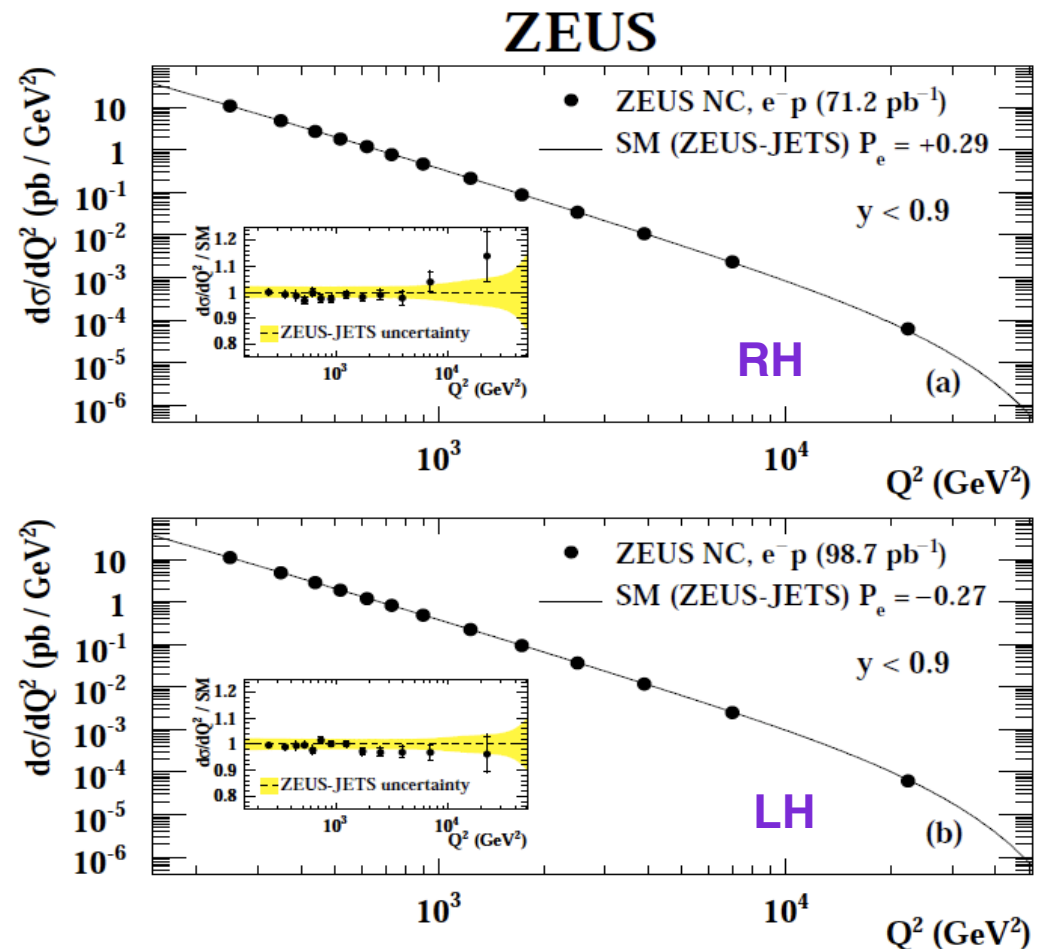
Polarized NC measurements

Single differential cross sections vs. Q^2 for polarized e

Data well described by SM

Sensitivity to polarization asymmetries

→ compare these directly (next slide)



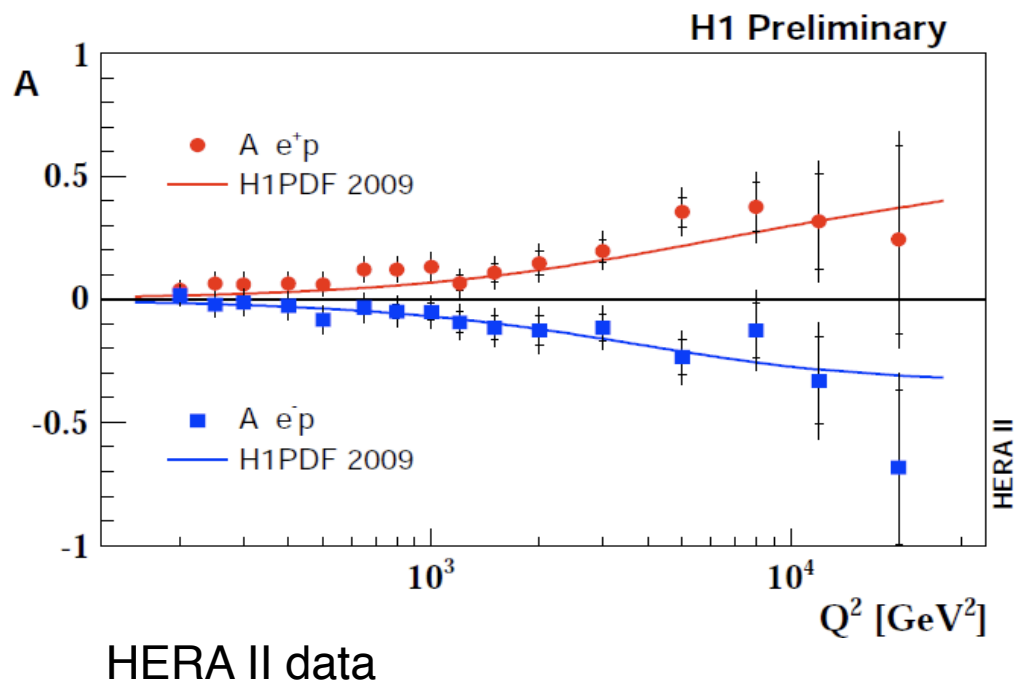
Polarized NC measurements Asymmetry

The charge dependent polarization asymmetries in neutral currents
→ direct measure of EW effects

Polarization asymmetries (A) sensitive to ratio of γZ interference term to F_2

A is proportional to $a_e v_q$ combination

$$A_{\pm} = \frac{2}{P_R - P_L} \frac{\sigma^{\pm}(P_R) - \sigma^{\pm}(P_L)}{\sigma^{\pm}(P_R) + \sigma^{\pm}(P_L)} \simeq \mp \kappa a_e \frac{F_2^{\gamma Z}}{F_2}$$



low $Q^2 \rightarrow$ contribution from Z negligible
cross sections overlap
high $Q^2 \rightarrow$ Z becomes important, A increases

Data well described by SM

recall: neglecting Z term,
generalized structure function F_2 is expressed:

$$\tilde{F}_2^{\pm} \approx F_2^{\gamma} + \kappa(-v_e \pm P_e a_e) F_2^{\gamma Z}$$

$$\text{At LO: } F_2^{\gamma Z} = x \sum_q 2e_q v_q (q + \bar{q})$$

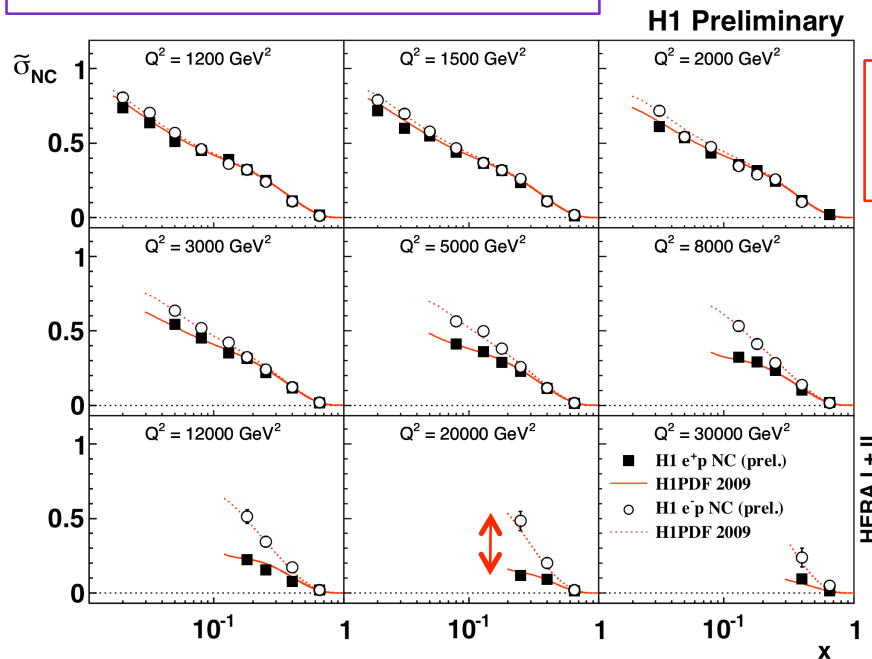
xF_3 measurements in NC

recall: NC cross section:

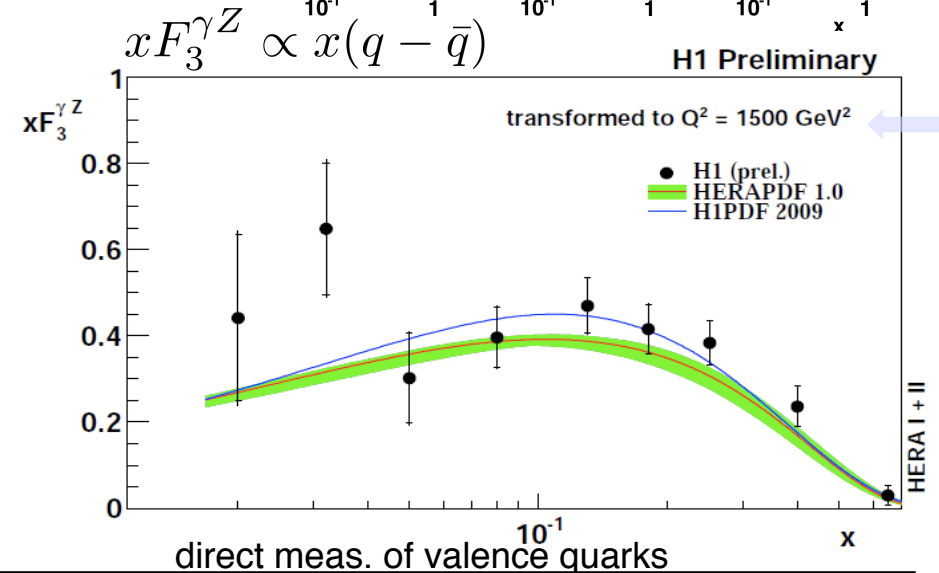
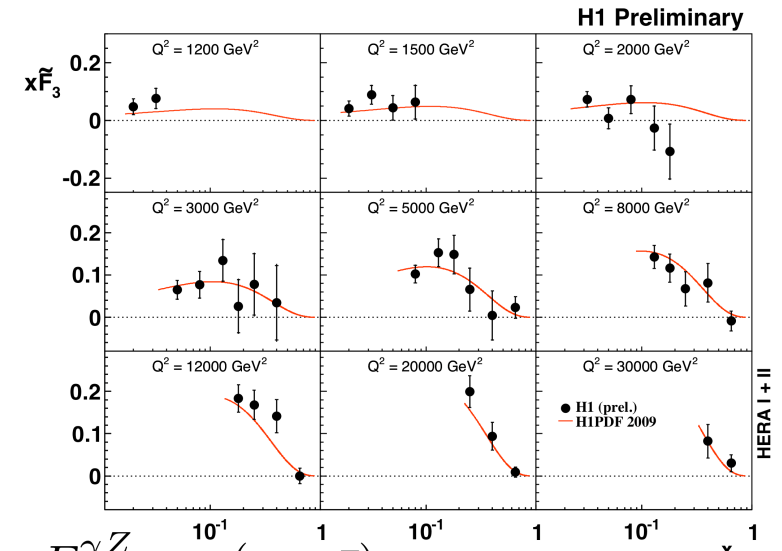
$$\tilde{\sigma}^{\pm} = \frac{d^2\sigma^{NC}(e^{\pm}p)}{dx dQ^2} \frac{xQ^4}{2\pi\alpha^2} \frac{1}{Y_{\pm}} = \tilde{F}_2 \mp \frac{Y_{-}}{Y_{+}} x\tilde{F}_3 - \frac{y^2}{Y_{+}} \tilde{F}_L$$

$$x\tilde{F}_3 = \frac{Y_{+}}{2Y_{-}} [\tilde{\sigma}^{-} - \tilde{\sigma}^{+}]$$

comparing reduced xsec
for e- and e+ can extract $x\tilde{F}_3$



dominant
contribution to $x\tilde{F}_3$: $x\tilde{F}_3^{\gamma Z} \simeq x\tilde{F}_3 \frac{(Q^2 + M_Z^2)}{\alpha^2 \kappa Q^2}$



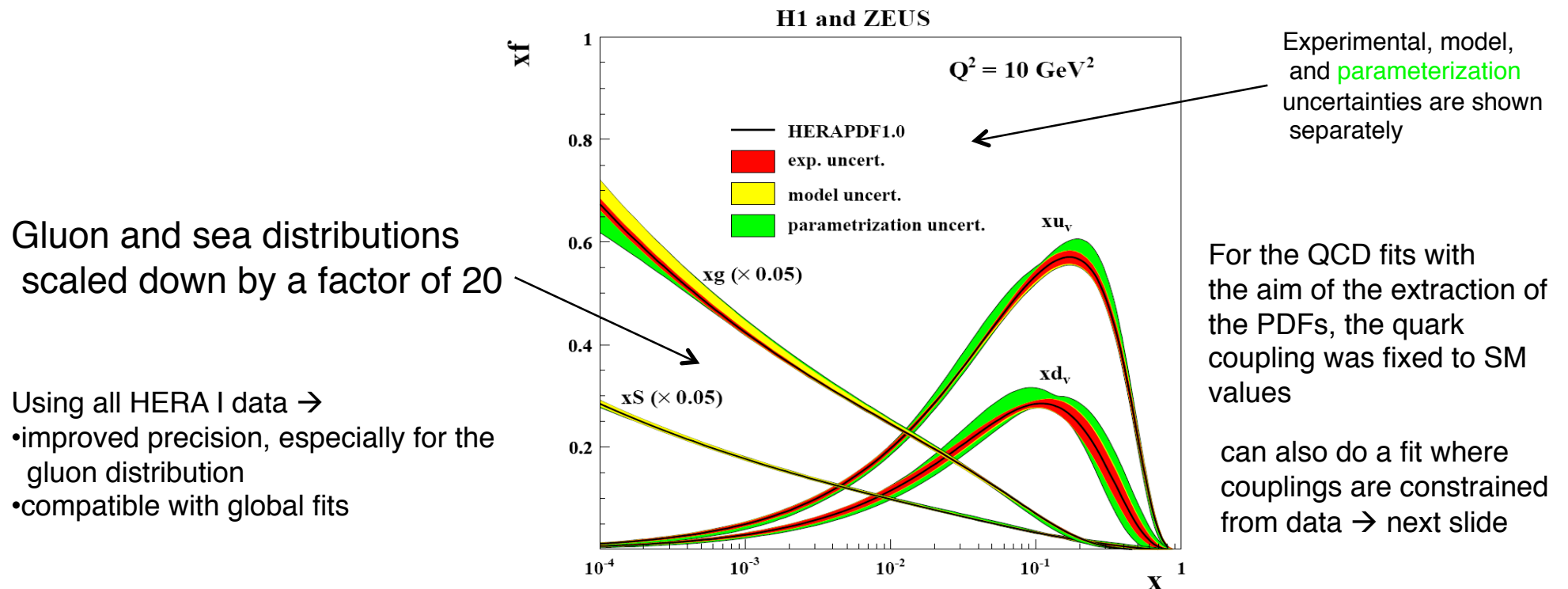
HERAPDF 1.0 fit

Seen several ways that HERA is sensitive to pdfs.

to extract best knowledge of PDFs: QCD fit using all of H1 + ZEUS combined data from HERA 1

→ New set of parton distributions HERAPDF1.0

Parton distributions unfolded in NLO QCD fit using the HERA $e^\pm p$ data only



JHEP 1001:109,2010 / arXiv:0911.0884 [hep-ex]

Combined QCD + EW fit results

Simultaneous EW+PDF analysis of NC and CC data

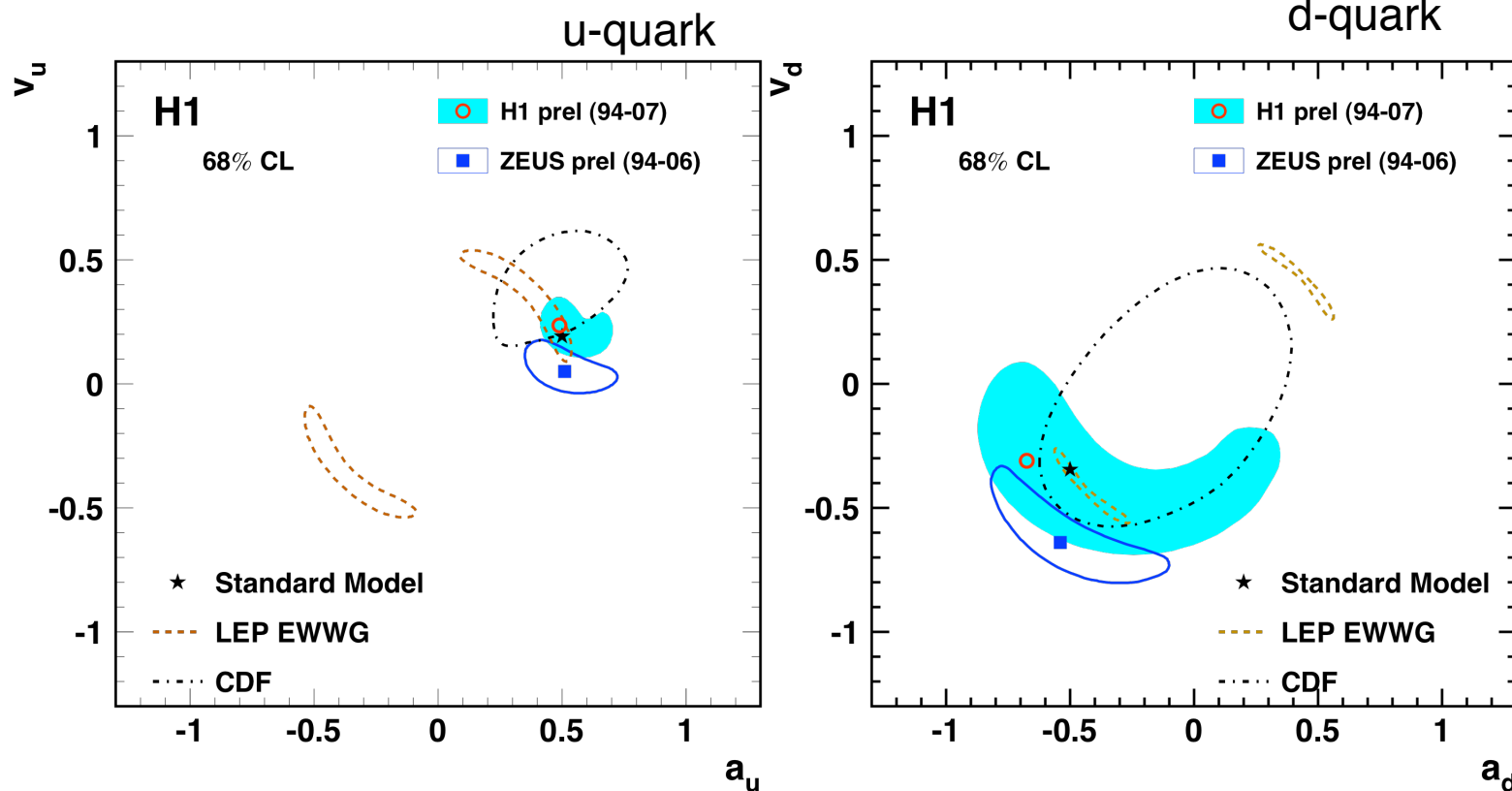
Now have precise enough data

→ can precisely fit the axial and vector coupling to u and d quarks

allowed parameter space:

combined HERA I & II

previous pub. HERA I



LEP
CDF

precision at HERA
can solve ambiguity
in LEP fit

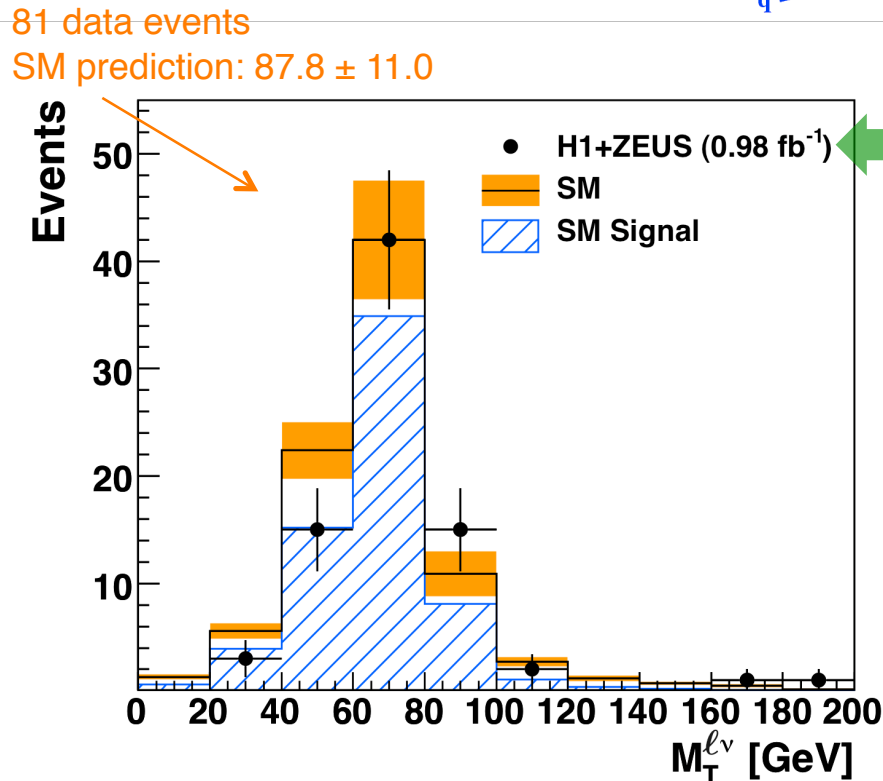
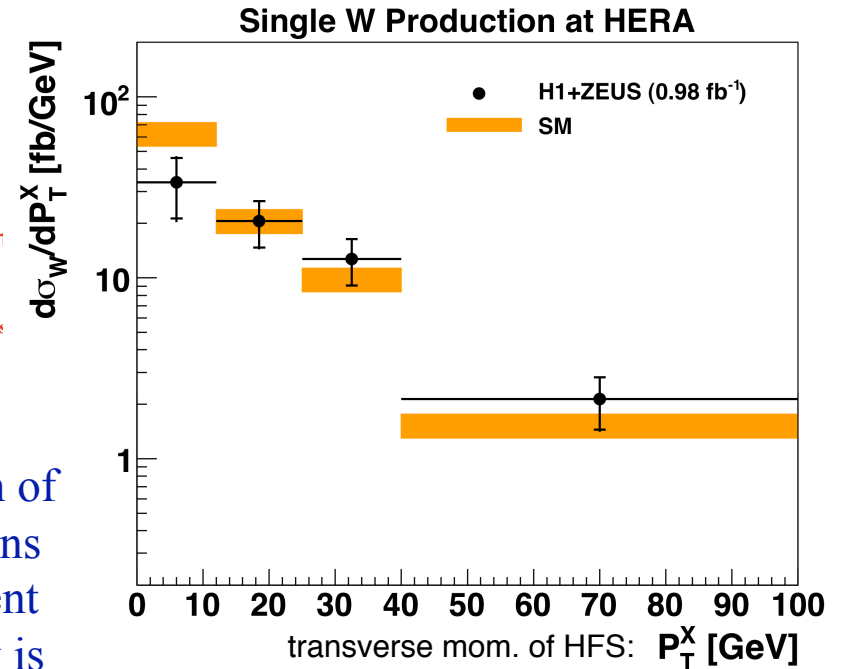
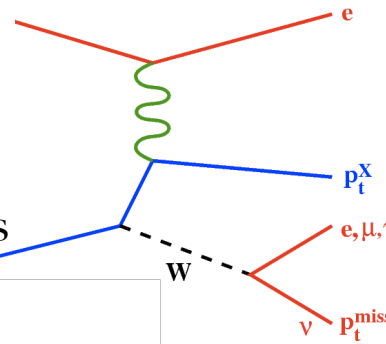
HERA II improves
precision, esp. for
u-quark

Weak couplings v_q much improved as expected

SM is ★ HERA data is consistent

Direct W production

- $\ell=e,\mu$ and ν with high P_T
- complete HERA data, $\sim 1 \text{ fb}^{-1}$
- combine W cross sections
electron and muon channels



The production of single W bosons with subsequent leptonic decay is the main signal contribution to the SM expectation

total single W cross section
 $\sigma = 1.06 \pm 0.16 \pm 0.07 \text{ pb}$

SM prediction
 $\sigma = 1.26 \pm 0.19 \text{ pb}$

Jacobian peak as expected from single W prod, all kinematic dists. show good agreement w/ SM

Summary

- NC and CC cross sections measured at high Q^2 for HERA I & HERA II
 - polarized and unpolarized
 - 1st combined results for the two experiments: ZEUS and H1
- extracted a set of PDFs: HERAPDF1.0
 - using combined H1 and ZEUS HERA I data
- EW fits to extract axial and vector couplings of quarks
- EW effects at HERA are well described by SM