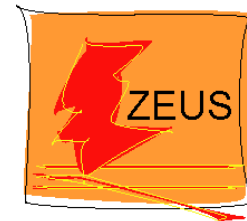


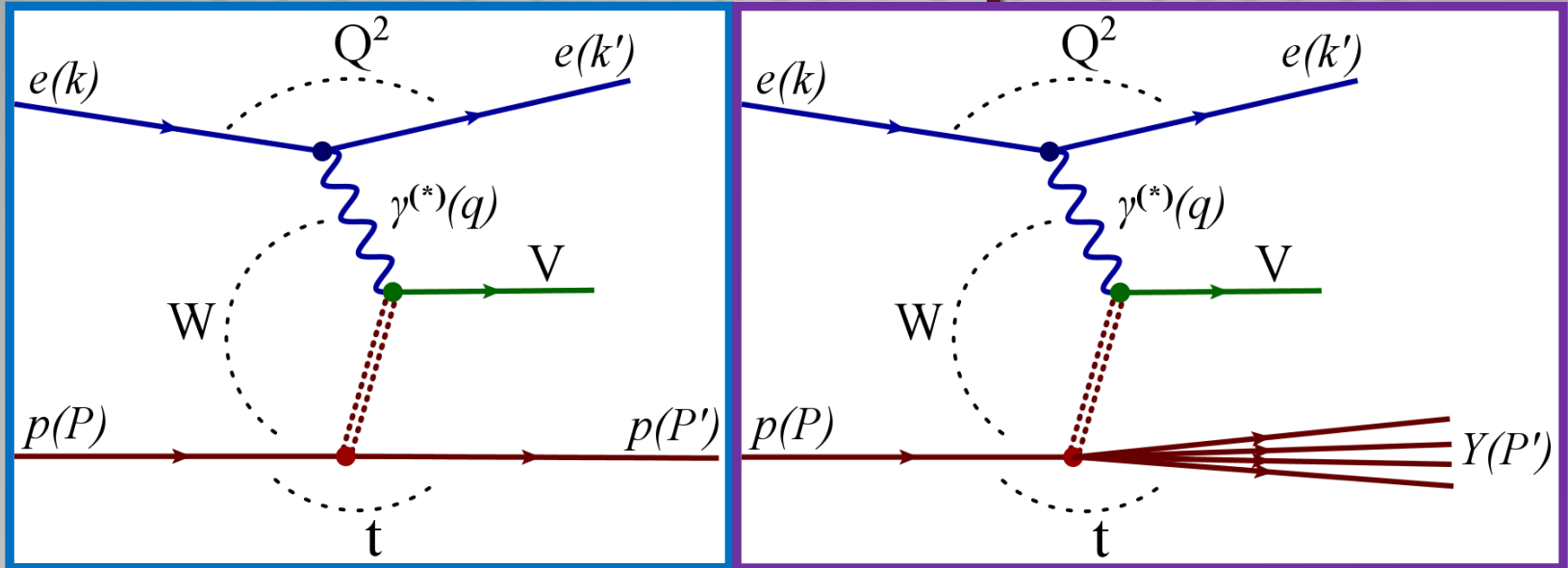
# **Vector Meson production and DVCS at HERA**

**On behalf of the H1 and ZEUS Collaborations**



***Janusz Tomasz Malka  
University of Łódź***

# Exclusive diffractive processes



**Diffractive** - no quantum numbers are exchanged in the interaction btw  $\gamma$  and  $p \rightarrow$  no colour flux  $\rightarrow$  large rapidity gap

$V$  - Vector Meson ( $\rho, \omega, \phi, J/\psi, \psi', \Upsilon$ ) photon ( $\gamma$ )

$Q^2$  - photon virtuality  $Q^2 = -q^2 = -(k - k')^2$

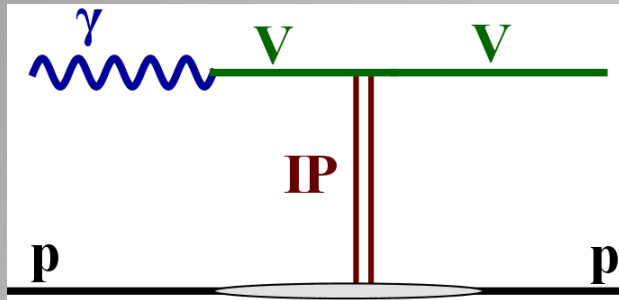
$(Q^2 \approx 0$  - photoproduction,  $Q^2 > 0$  - electroproduction)

$W$  - c.m. energy of  $\gamma p$  system  $W^2 = (q + p)^2$

$t$  - (4-mom. transfer) $^2$  at  $p$ -vertex  $t = (P - P')^2$

The proton can stay intact  $p(P')$  or dissociate  $Y(P')$

# VDM & Regge theory (Soft physics)



**Vector Dominance Model (VDM):**

- The photon fluctuates into a vector meson,  $V$ , which carries the same quantum numbers as the photon ( $\gamma p \rightarrow Vp$ )
- The vector meson scatters elastically off the incoming proton ( $Vp \rightarrow Vp$ )

$$\frac{d\sigma(\gamma p \rightarrow Vp)}{dt} \propto e^{-b_0 t} \left( \frac{W^2}{W_0^2} \right)^{2(\alpha(t)-1)}$$

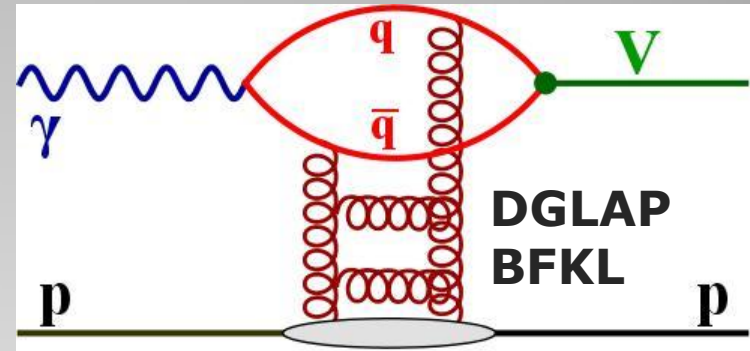
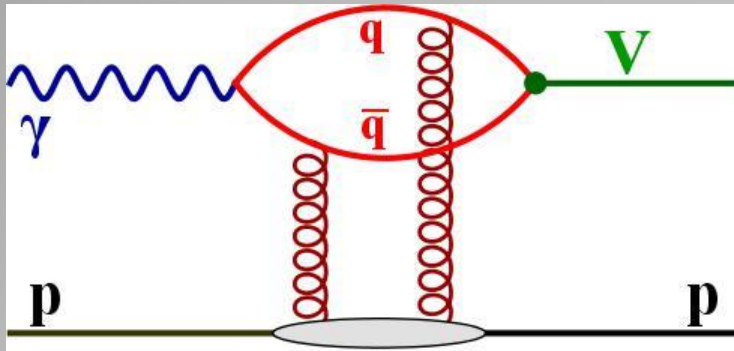
**Predictions:**

**Experimental observations:**

- $\alpha(t) = \alpha(0) + \alpha' t$  ,  $\alpha(0) = 1.08$  ,  $\alpha' = 0.25$  ; **(Donnachie-Landschof)**
- **Shrinkage of the diffractive peak**  $b(W) = b_0 + 4\alpha' \ln \left( \frac{W}{W_0} \right)$
- **Weak energy dependence of cross section**

$$\sigma \propto W^\delta , \quad \delta \approx 4 \left[ \alpha(0) - \frac{\alpha'}{b_0} - 1 \right] , \quad \delta \sim 0.2 ;$$

# *pQCD models (Hard physics)*



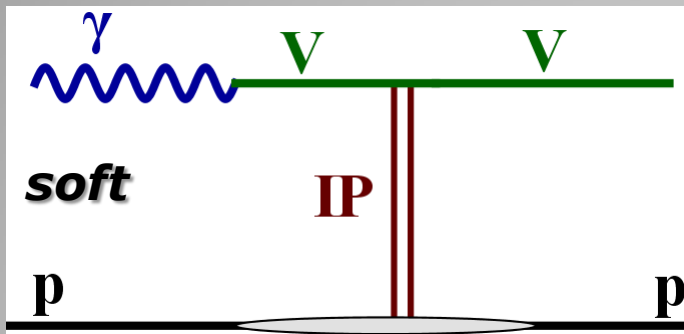
- the photon fluctuates into a  $q\bar{q}$  state
- the  $q\bar{q}$  pair scatters off the proton target
- the scattered  $q\bar{q}$  pair turns into a vector meson

## Predictions:

- $\sigma \propto [xg(x, \mu^2)]^2$  where  $\mu^2 = f(Q^2, M_V^2, t)$
- **Increasing  $W$  corresponds to going to small  $x$**
- **fast increase of the  $\gamma^* p \rightarrow Vp$  cross section with energy  $W$**
- **universal exponential  $t$  dependence**
- $b \sim 4 - 5 \text{ GeV}^{-2} \Rightarrow \alpha' \approx 0$

$$x \approx \frac{Q^2}{W^2} \text{ at low } x$$

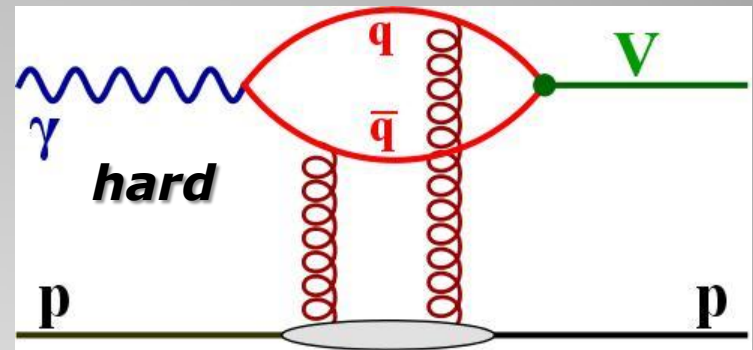
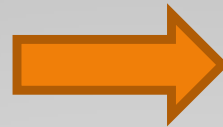
# Transition from soft to hard physics



$$\delta \approx 0.2$$

$$b \approx 10 \text{ GeV}^{-2}$$

$$\alpha' \approx 0.25 \text{ GeV}^{-2}$$



**hard**

$$\sigma \propto W^\delta$$

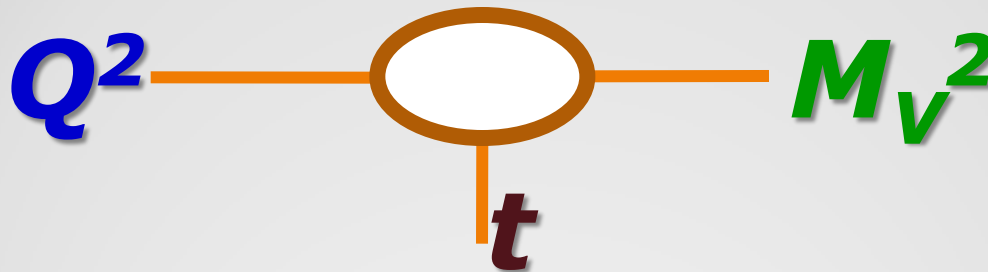
$$\frac{d\sigma}{d|t|} \propto e^{-b|t|}$$

$$\alpha(t), b(W)$$

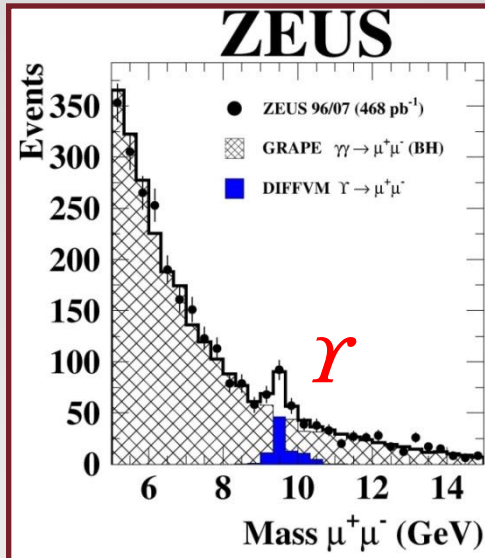
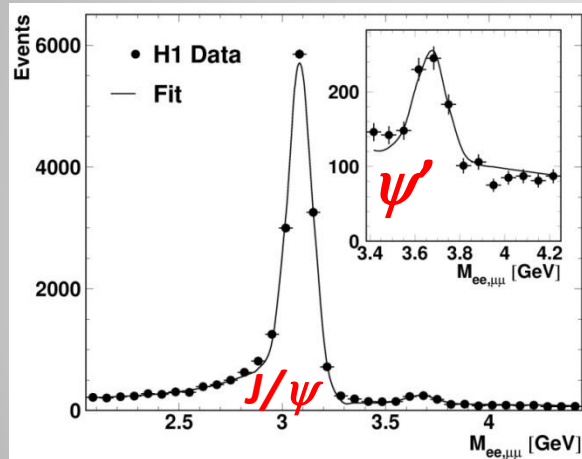
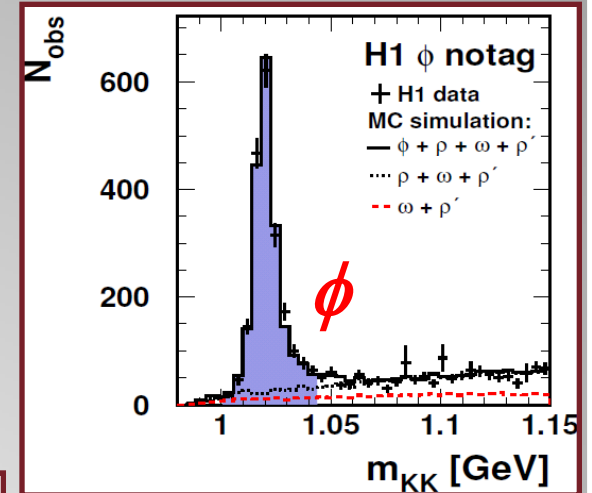
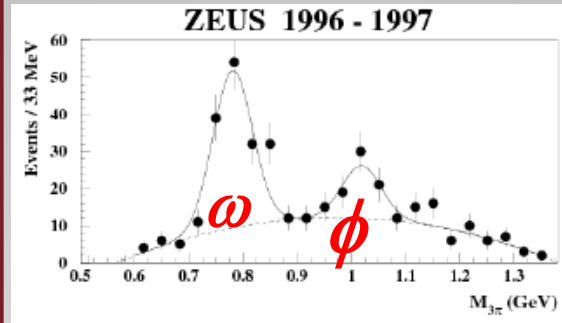
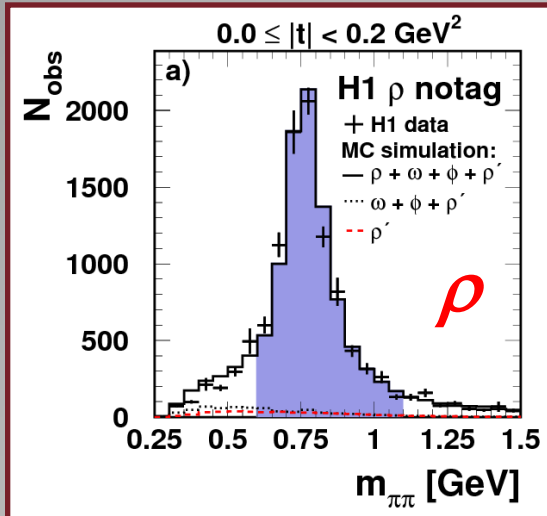
$$\delta \leq 0.7$$

$$b \approx 4\text{-}5 \text{ GeV}^{-2}$$

$$\alpha' \approx 0$$

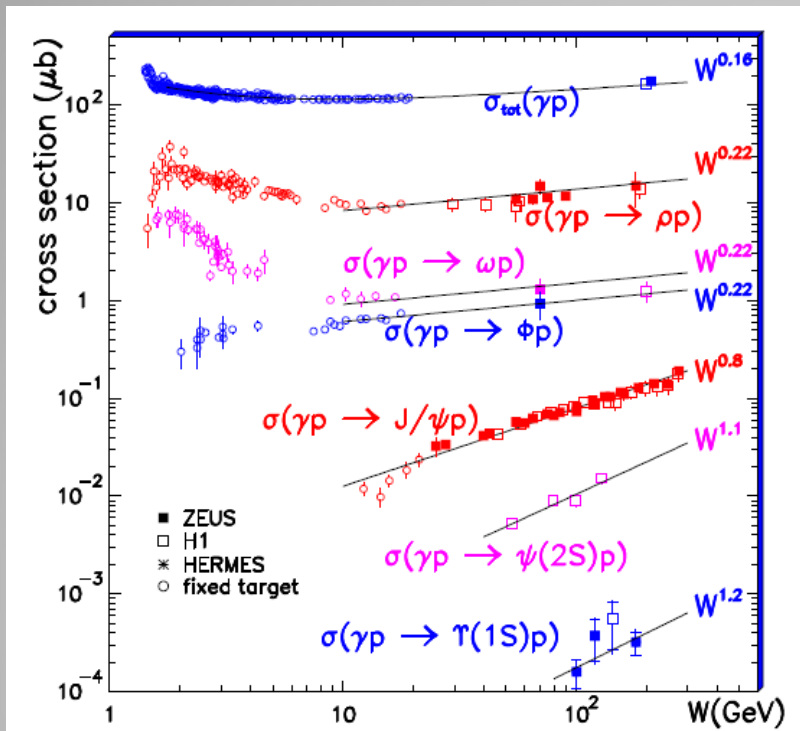


# Vector mesons mass distributions



# ***W dependence of the photoproduction of VM***

The parameterisation of the cross section:  $\sigma \propto W^\delta$



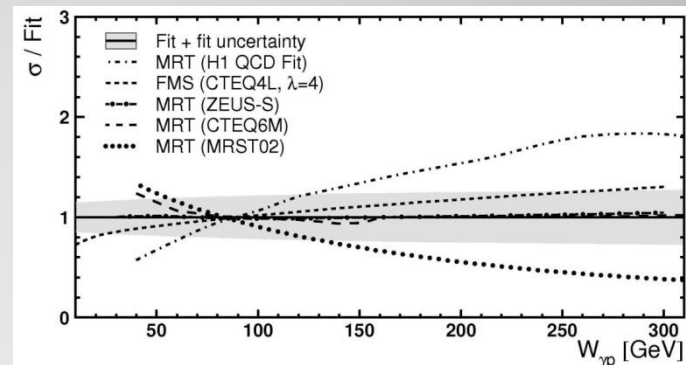
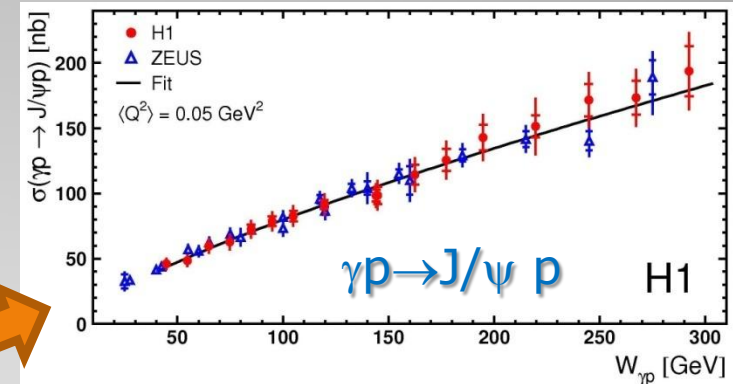
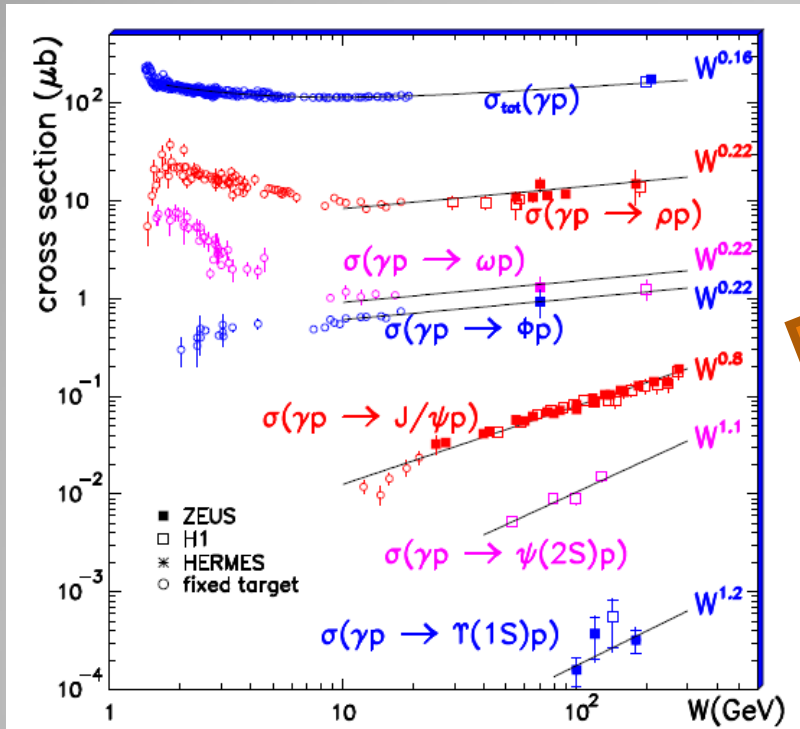
**Process becomes hard  
(steeper W dependence)  
as mass of VM becomes larger**



**Apply pQCD models**

# ***W dependence of the photoproduction of VM***

The parameterisation of the cross section:  $\sigma \propto W^\delta$



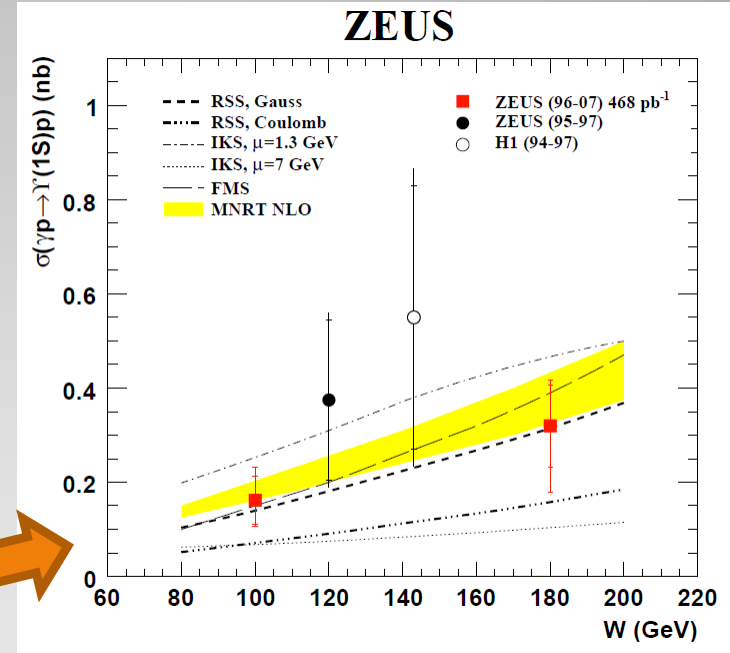
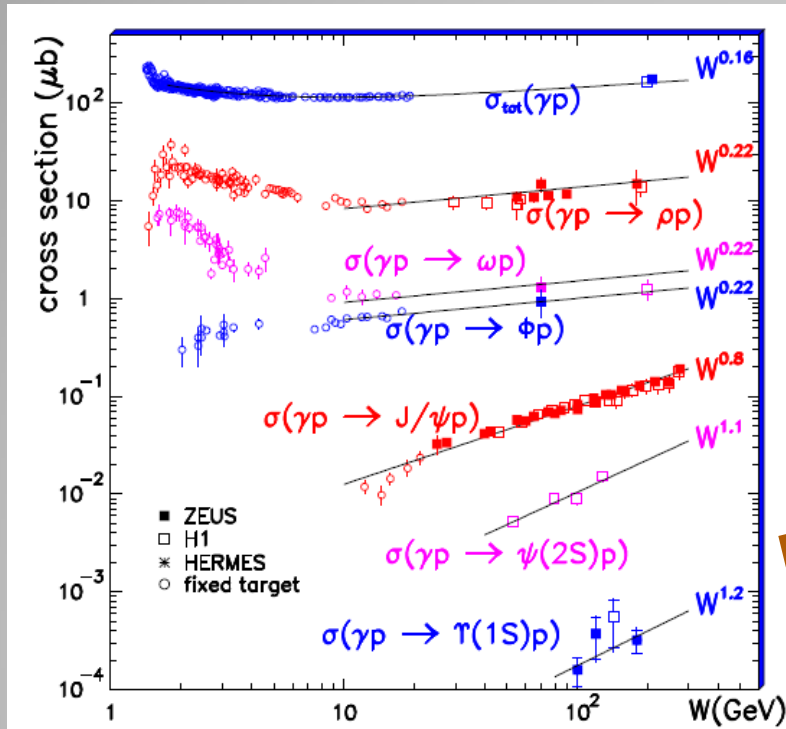
**FMS** – L. Frankfurt, M. McDermott and M. Strikman, JHEP **0103** (2001)  
**MRT** – A. D. Martin, M. G. Ryskin and T. Teubner, Phys. Rev. D **62** (2000) 014022

***Sensitivity to gluon density***



# ***W dependence of the photoproduction of VM***

The parametrisation of the cross section:  $\sigma \propto W^\delta$



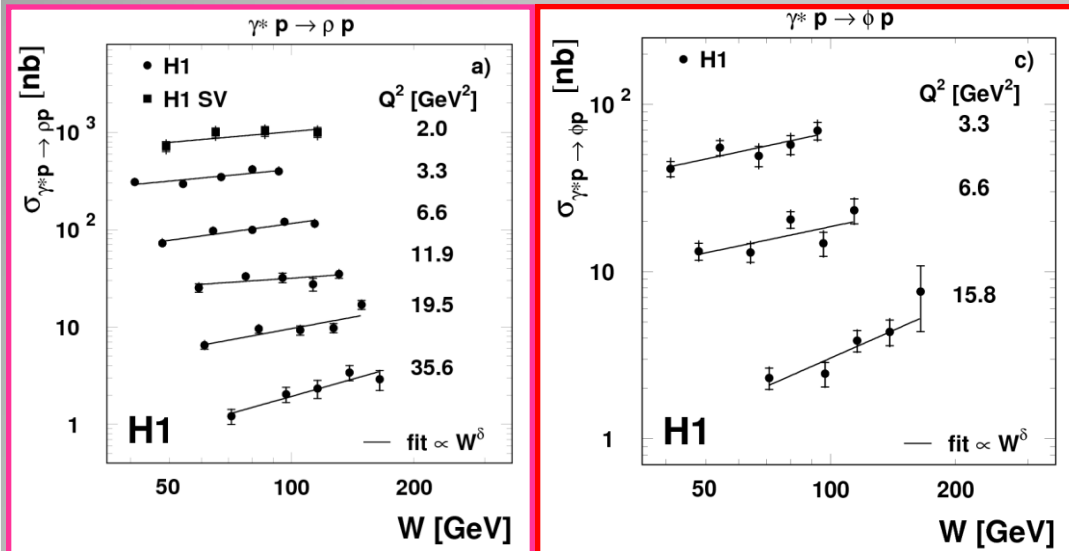
**RSS ( $k_T$ )** – A. **Rybarska**, W. **Schäfer**, A. **Szczurek**,  
Phys. Lett. B 668 (2008), p. 126.

**IKS(NLO)** – D.Yu. **Ivanov**, G. **Krasnikov**, L. **Szymanowski**,  
Nucl. Phys. B (Proc. Suppl.) 146 (2005), p. 134.

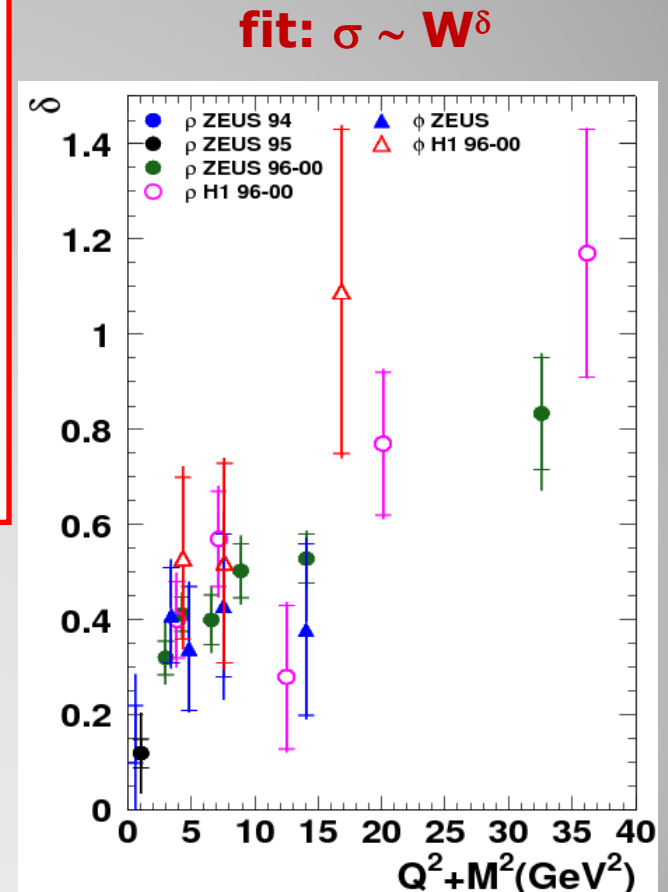
**FMS(CTEQ4L)** – L.L. **Frankfurt**, M.F. **McDermott**,  
M. **Strikman**, JHEP 9902 (1999), p. 002.

**MNRT(HERA J/ψ)** – A.D. **Martin**, C. **Nockles**,  
M. **Ryskin**, T. **Teubner**, Phys. Lett.B 662 (2008), p. 252.

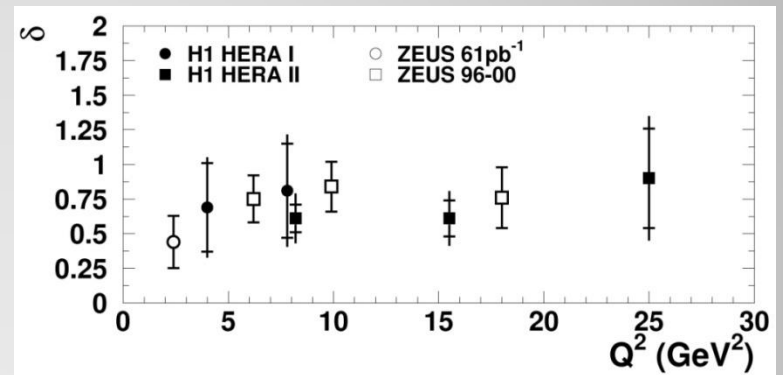
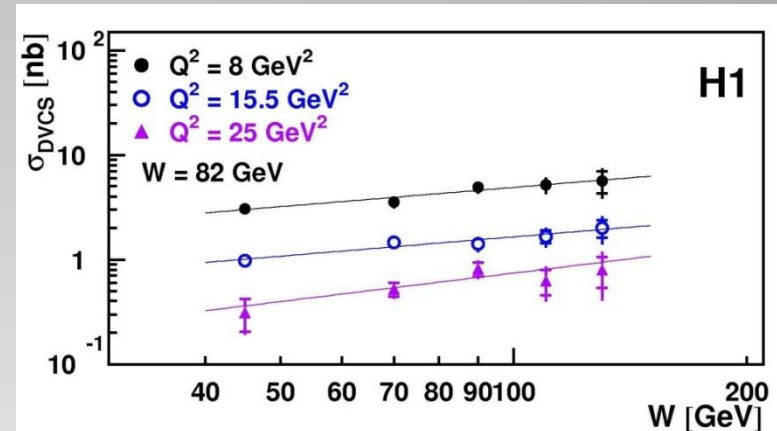
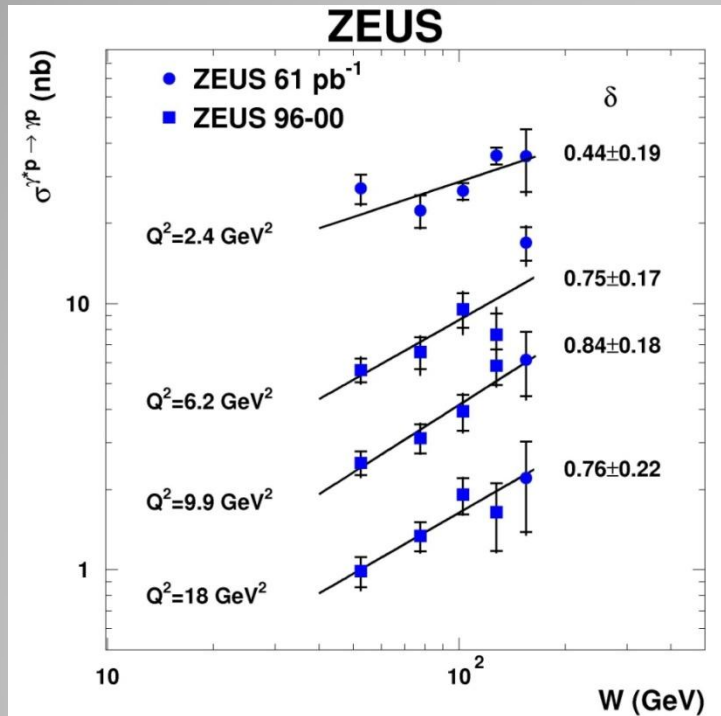
# VM electroproduction (W dependence)



cross section  $W$  dependence becomes steeper at high  $Q^2$  for  $\rho$ ,  $\phi$  from soft to hard regime

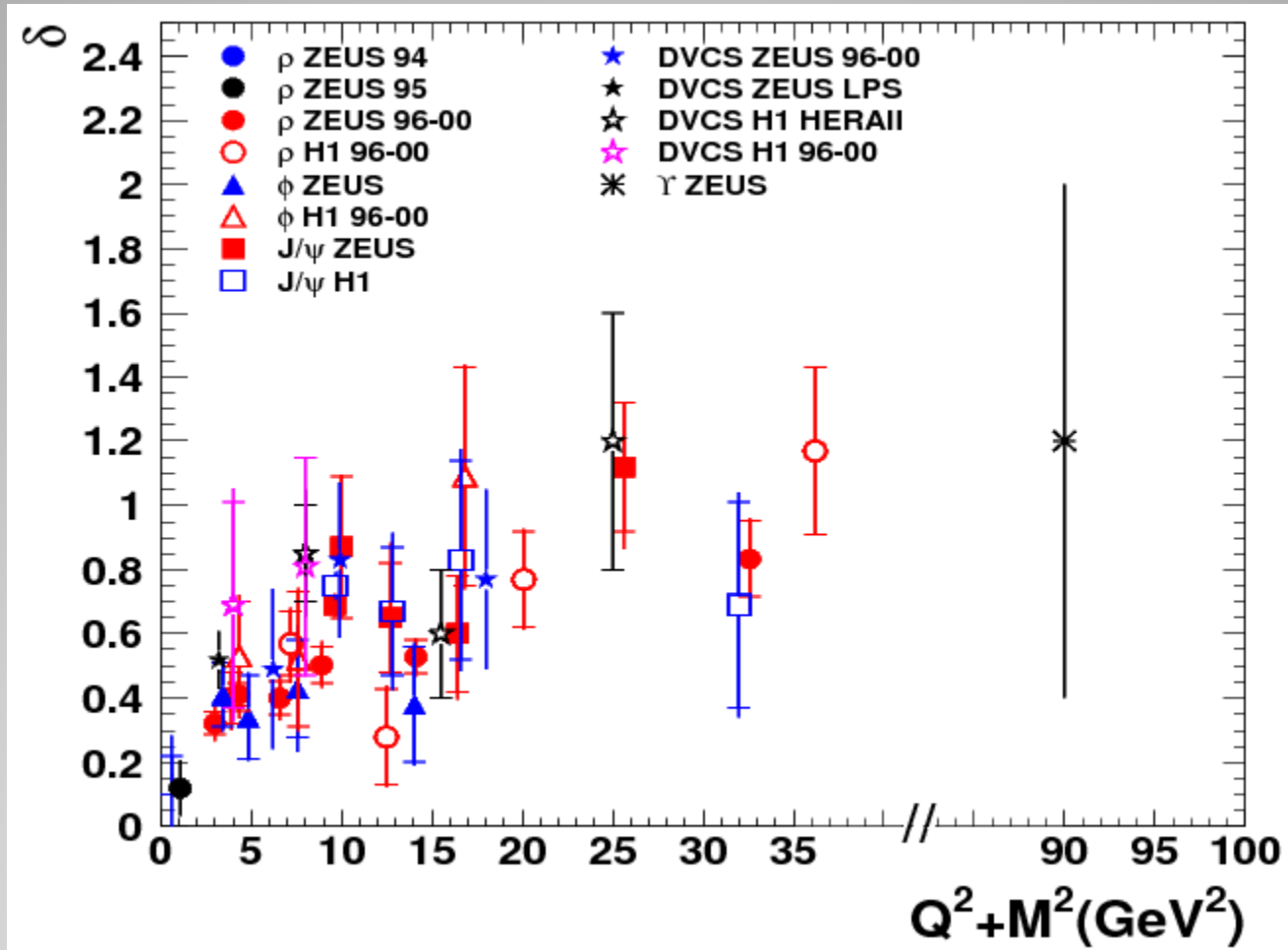


# Deeply Virtual Compton Scattering (W dependence)



- fit:  $\sigma \sim W^\delta$
- No  $\delta$  dependence on  $Q^2$  is observed
- Hard regime

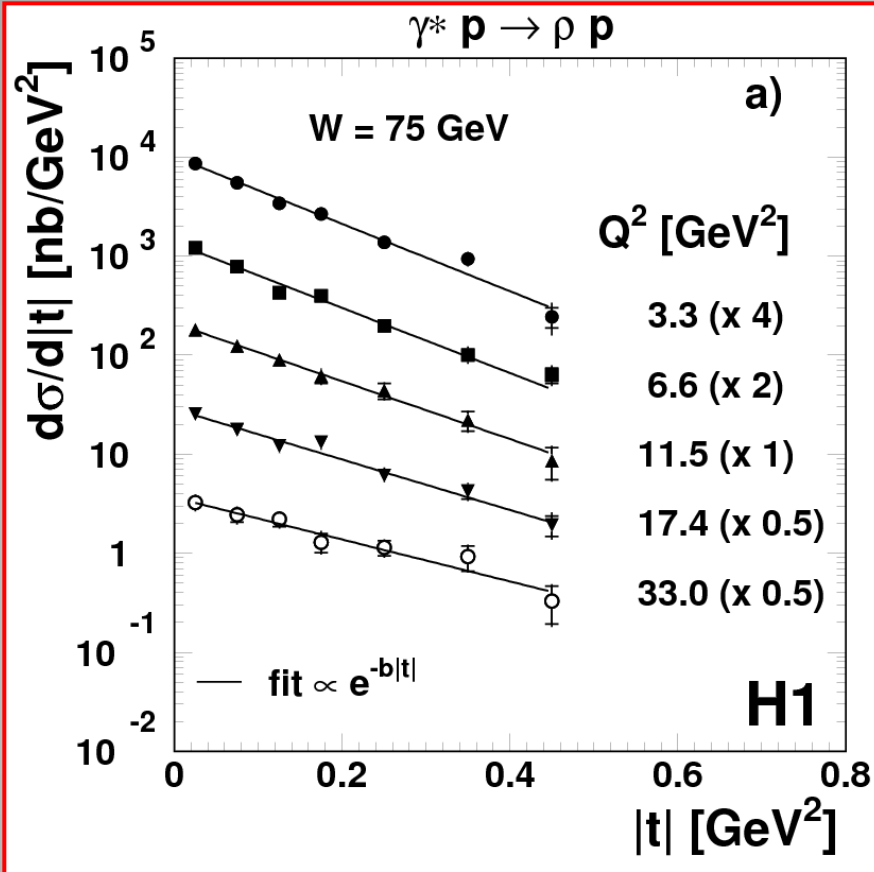
# Summary (W dependence)



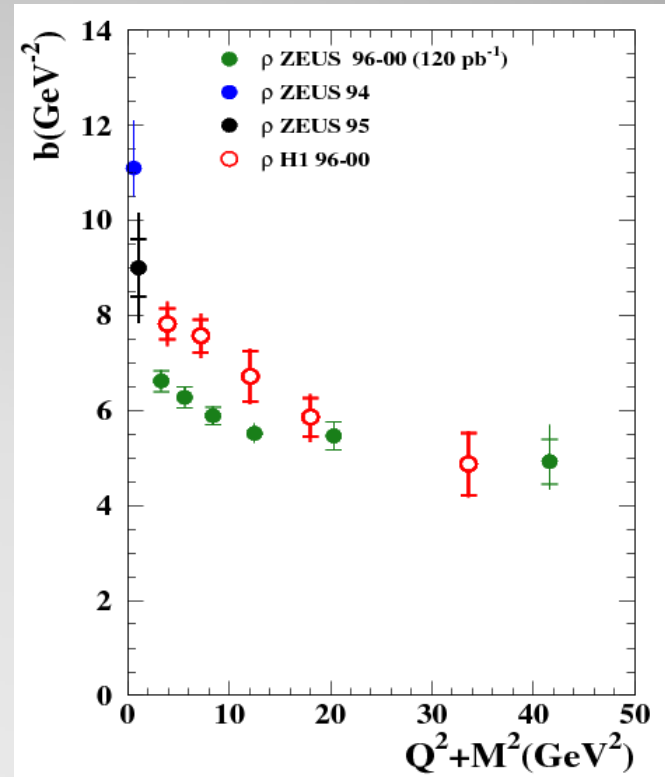
• process becomes hard as scale ( $Q^2 + M^2$ ) becomes larger

# $|t|$ dependence ( $\rho$ )

$d\sigma/d|t| \sim \exp(-b|t|)$  in bins of  $Q^2$



$b$  describes the transverse size of the interaction region  $b \propto R^2$

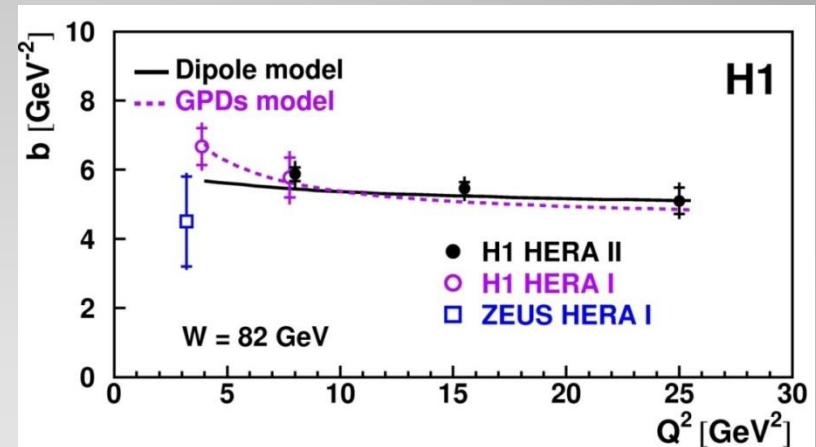
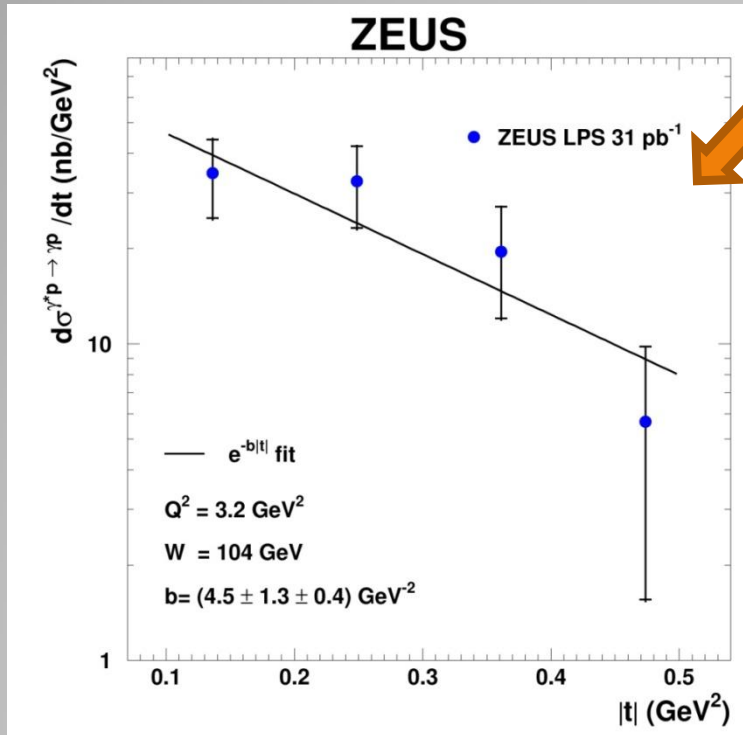


Transverse size of dipole decreases with  $Q^2$

$b$  decreases from soft values to pQCD expected values ( $\sim 4-5$  GeV<sup>-2</sup>)

# $|t|$ dependence (DVCS)

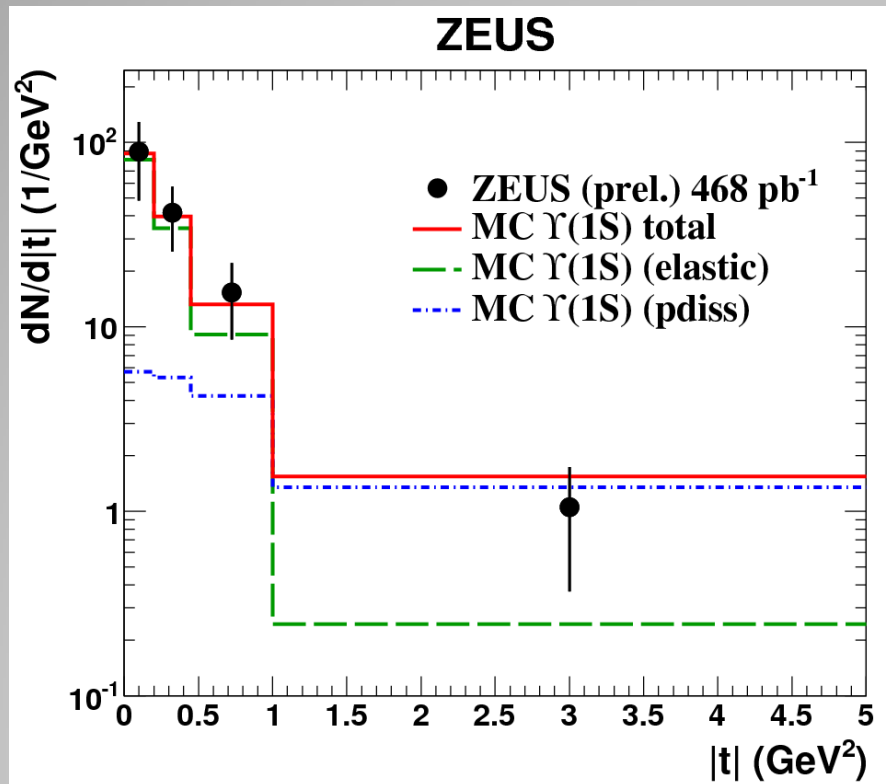
direct  $t$  measurement using Leading Proton Spectrometer



**H1:  $b = 5.41 \pm 0.14 \pm 0.31$  GeV<sup>-2</sup>**  
**ZEUS:  $b = 4.5 \pm 1.3 \pm 0.4$  GeV<sup>-2</sup>**

- DVCS point don't show  $Q^2$  dependence

# $|t|$ dependence (*Upsilon*) photoproduction



2 times far from present results  
at  $Q^2 + M^2 = 89.5 \text{ GeV}^2$

**Elastic:**  $b = 4.3^{+1.7}_{-1.1} \quad +0.5_{-0.5} [\text{GeV}^{-2}]$

# $|t|$ dependence

$$\frac{d\sigma}{dt} \sim e^{-b|t|}$$

transverse size:

$$b = b_V + b_p$$

Vector Meson  
transverse size

$$b_V \sim \frac{1}{Q^2 + M_V^2}$$

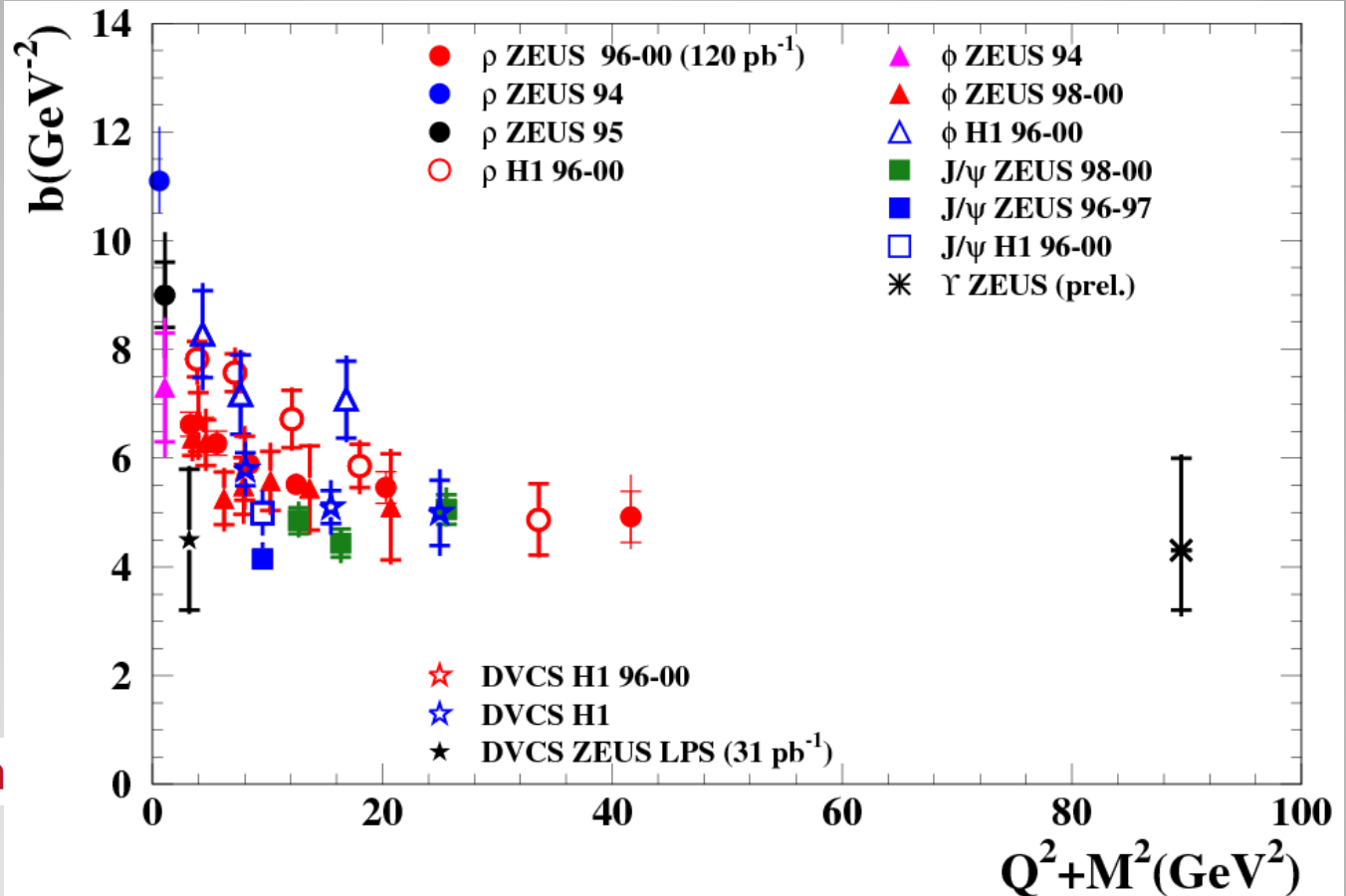
Target size:

$$b_p \approx 5 \text{ GeV}^{-2}$$

corresponds to

$$r_{\text{gluons}} \approx 0.5 \text{ fm}$$

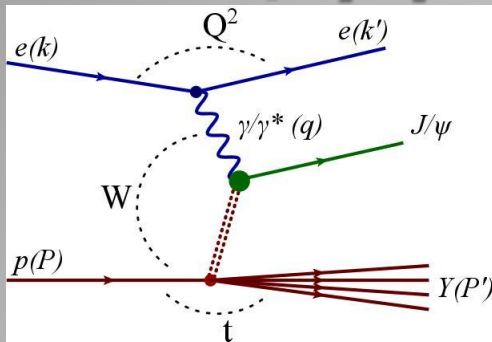
$$r_{\text{em}} \approx 0.8 \text{ fm}$$



**b decreases with from  $\sim 10 \text{ GeV}^{-2}$  (soft) to  $\sim 5 \text{ GeV}^{-2}$  (hard)**  
**size of scattered VM getting smaller with scale**



# ***$J/\psi$ photoproduction at high $|t|$***

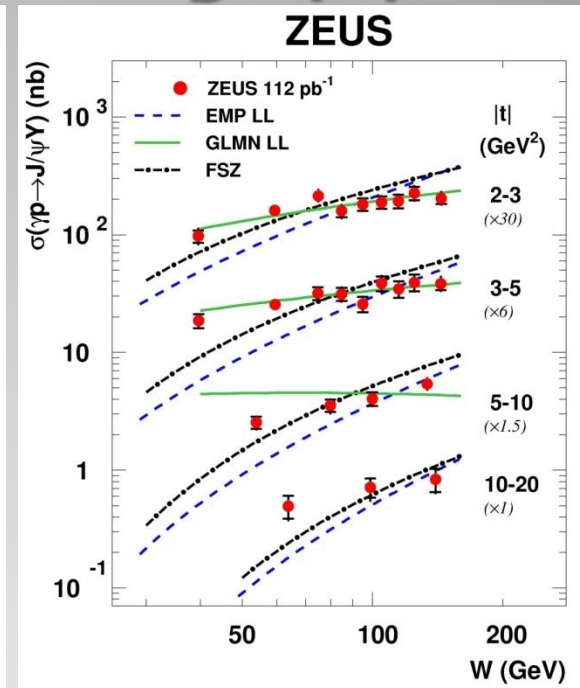
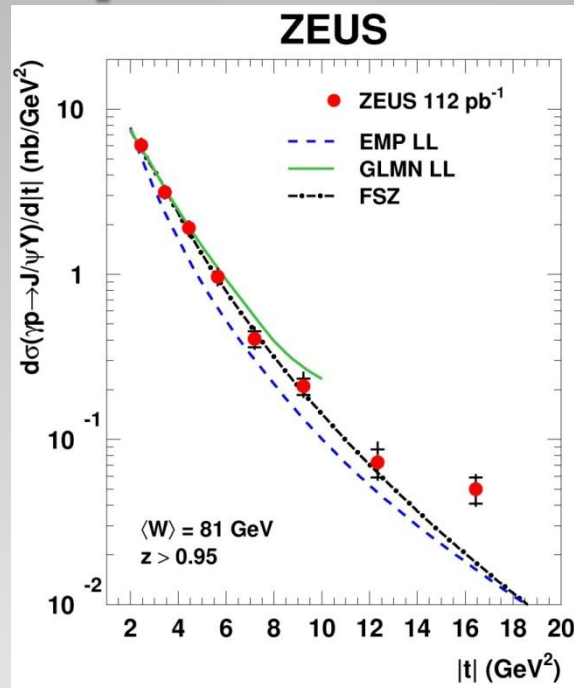


Hard scale provided by  $|t|$   
and mass  
 $t$  dependence no longer  
exponential

$$\frac{d\sigma}{d|t|} \propto t^n$$

$$n = -1.9 \pm 0.1, \quad 2 < |t| < 4 \text{ GeV}^2$$

$$n = -3.0 \pm 0.1, \quad 4 < |t| < 16 \text{ GeV}^2$$



***$\sigma$  vs  $W$  in  $t$  ranges: data rise with  $W$  for all  $t$***

***EMP (BFKL) below data***

***GLMN (DGLAP) fails at  $|t| > 5 \text{ GeV}^2$***

***FSZ gives good description up  $|t| = 12 \text{ GeV}^2$***

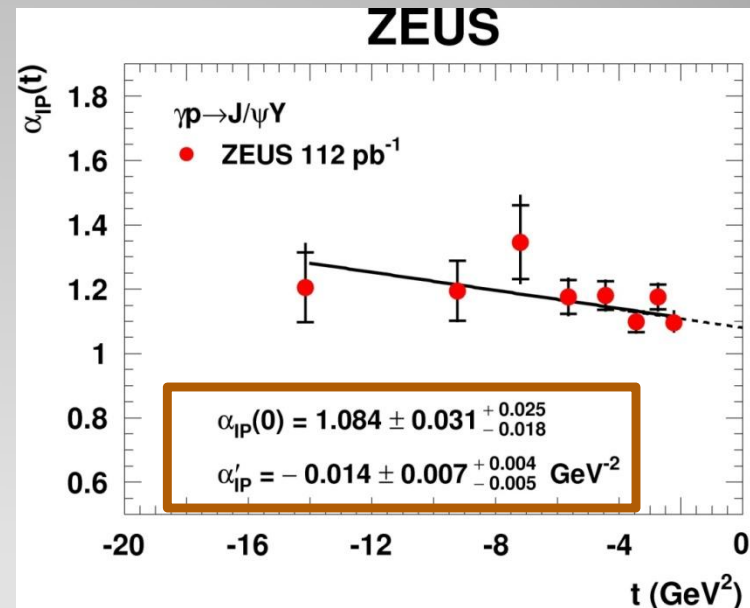
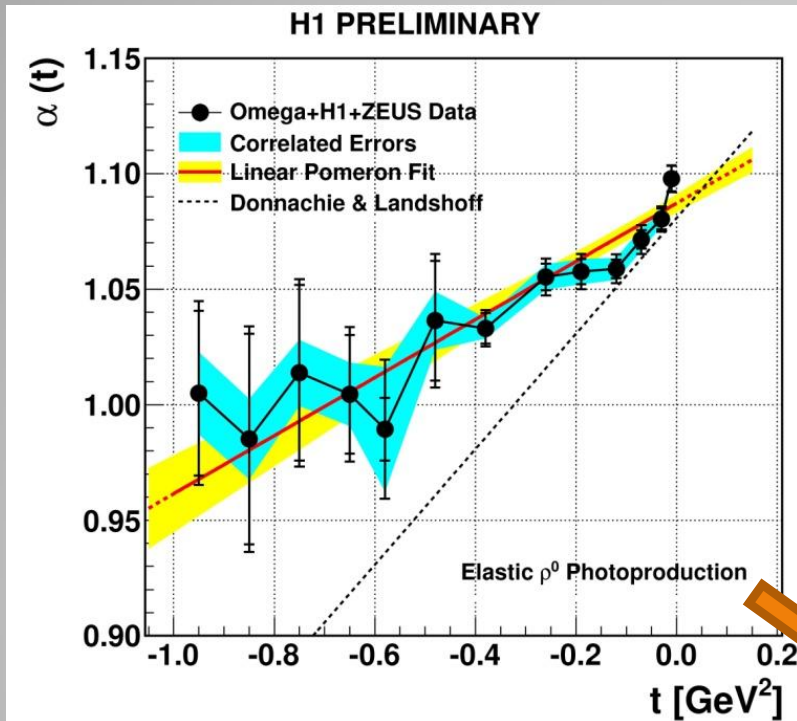
***BFKL LL and FSZ are too steep ( $W$  dependence)***

***None of the models describes the data over the full phase space***

# effective Pomeron trajectory

$$\frac{d\sigma}{d|t|} \propto \left(\frac{W}{W_0}\right)^{4(\alpha(t)-1)}$$

$$\alpha(t) = \alpha(0) + \alpha' t$$



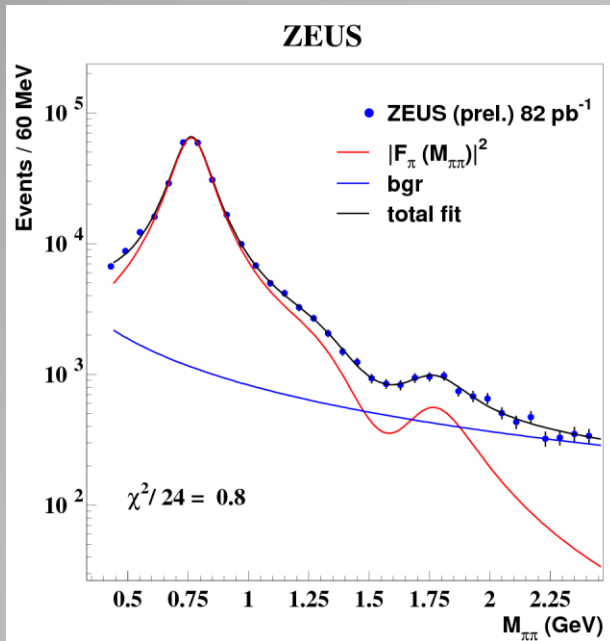
$$DL: \alpha(t) = 1.08 + 0.25t$$

$$\alpha(0) = 1.0871 \pm 0.0026 \pm 0.0030$$

$$\alpha' = 0.126 \pm 0.013 \pm 0.012 \text{ GeV}^{-2}$$

**$\alpha(0)$  consistent with 1.08 from soft pp scattering**  
 **$\alpha' \sim$  twice smaller than 0.25 GeV<sup>-2</sup>**

# Two pion diffractive electroproduction



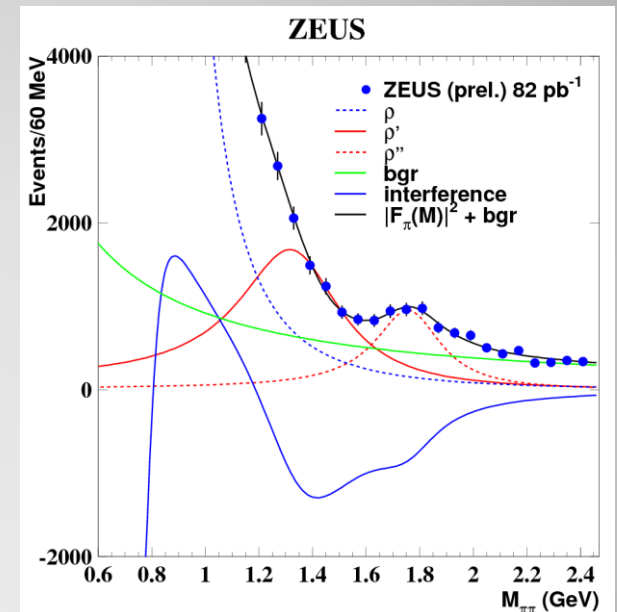
Measure two pion mass distribution  
 $0.4 < M_{\pi\pi} < 2.4 \text{ GeV}$  in  $2 < Q^2 < 80 \text{ GeV}^2$

$$\frac{dN(M_{\pi\pi})}{dM_{\pi\pi}} = N \left[ |F_{\pi}(M_{\pi\pi})|^2 + \frac{B}{M_{\pi\pi}^n} \right]$$

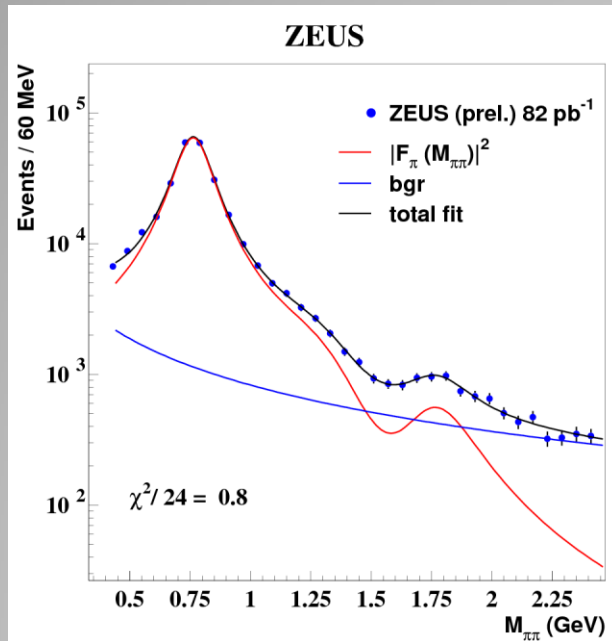
$$F_{\pi}(M_{\pi\pi}) = \frac{BW(\rho) + \beta BW(\rho') + \gamma BW(\rho'')}{1 + \beta + \gamma}$$

**11 parameters fit:**

$N, B, n, M(\rho), \Gamma(\rho), M(\rho'), \Gamma(\rho'), M(\rho''), \Gamma(\rho''), \beta, \gamma$ ,  
 where  $N$  - total normalization factor,  $B, n$  - described  
 background,  $M, \Gamma$  - masses and widths of vector mesons  
 and  $\beta, \gamma$  - relative amplitudes



# Two pion diffractive electroproduction



**Measure two pion mass distribution**  
 **$0.4 < M_{\pi\pi} < 2.4$  GeV in  $2 < Q^2 < 80$  GeV<sup>2</sup>**

$$\frac{dN(M_{\pi\pi})}{dM_{\pi\pi}} = N \left[ |F_{\pi}(M_{\pi\pi})|^2 + \frac{B}{M_{\pi\pi}^n} \right]$$

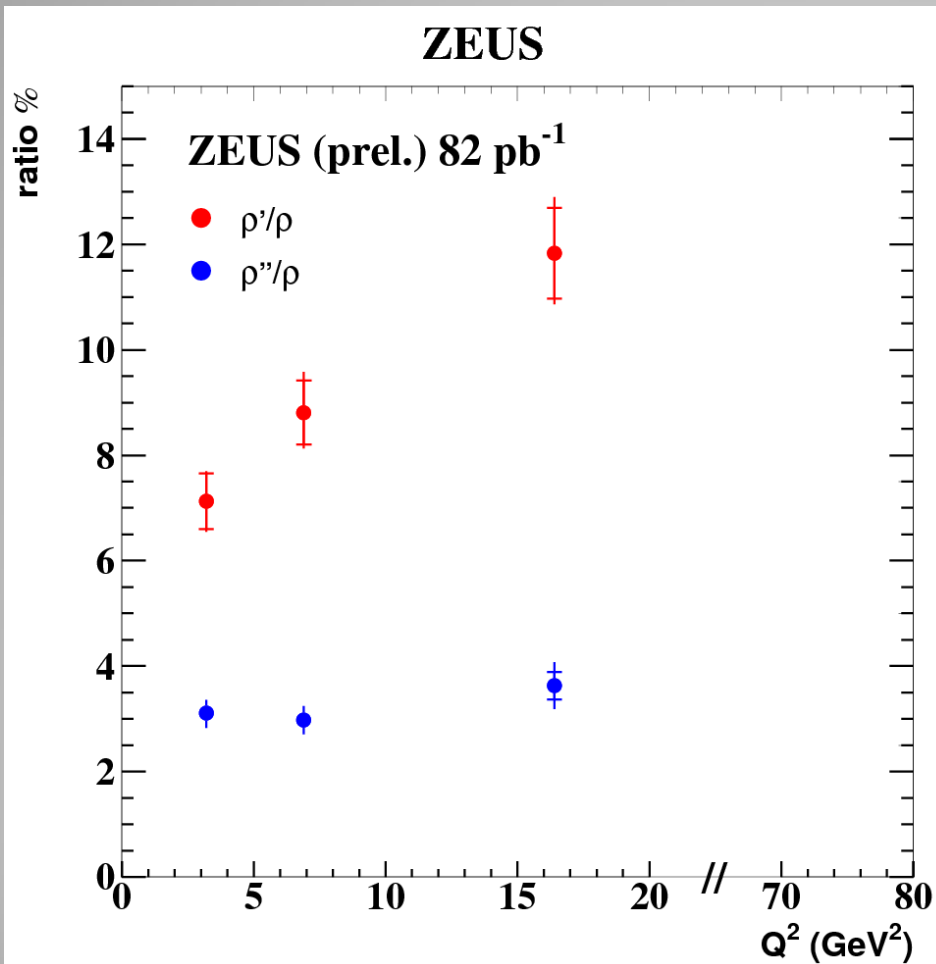
$$F_{\pi}(M_{\pi\pi}) = \frac{BW(\rho) + \beta BW(\rho') + \gamma BW(\rho'')}{1 + \beta + \gamma}$$

**11 parameters fit:**

**$N, B, n, M(\rho), \Gamma(\rho), M(\rho'), \Gamma(\rho'), M(\rho''), \Gamma(\rho''), \beta, \gamma$ ,  
 where  $N$  - total normalization factor,  $B, n$  - described  
 background,  $M, \Gamma$  - masses and widths of vector mesons  
 and  $\beta, \gamma$  - relative amplitudes**

Par.	ZEUS (prel.)	PDG
$M_{\rho}$	$772 \pm 2^{+2}_{-1}$	$775.49 \pm 0.34$
$\Gamma_{\rho}$	$155 \pm 5 \pm 2$	$149.4 \pm 1$
$M_{\rho'}$	$1360 \pm 20^{+20}_{-30}$	$1465 \pm 25$
$\Gamma_{\rho'}$	$460 \pm 30^{+40}_{-45}$	$400 \pm 60$
$\beta$	$-0.27 \pm 0.02 \pm 0.02$	
$M_{\rho''}$	$1770 \pm 20^{+15}_{-20}$	$1720 \pm 20$
$\Gamma_{\rho''}$	$310 \pm 30^{+25}_{-35}$	$250 \pm 100$
$\gamma$	$0.10 \pm 0.02^{+0.02}_{-0.01}$	

# $\rho, \rho', \rho''$ electroproduction (cross section ratios)



**$\rho'/\rho$ —increases with  $Q^2$**

the **anomalous behaviour of  $\rho'/\rho$  production** ratio with  $Q^2$  was predicted in works of J. Nemchik, B. Kopeliovich, N. Nikolaev, B. Zakharov, (see hep-ph/9605208)

**$\rho''/\rho$ —constant with  $Q^2$**

# Summary

***New measurements of  $\rho$ ,  $\phi$  and  $\Upsilon$  at HERA.***

***The measurements allow us to study the transition from the soft to hard regime.***

***Effective Pomeron trajectory has smaller slope than that extracted from soft hadron-hadron scattering.***

***DVCS asymmetries and  $|t|$  slope measurements provide access to GPDs.***

***Measurement of  $\rho$ ,  $\rho'$ ,  $\rho''$  production shows anomalous behaviour of  $\rho'/\rho$  cross section ratio***

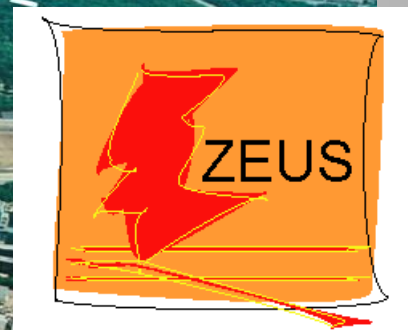
***pQCD expectations are in general compatible with the data, but still a lot to be understood.***

# ***Backup***



# ***HERA experiments***

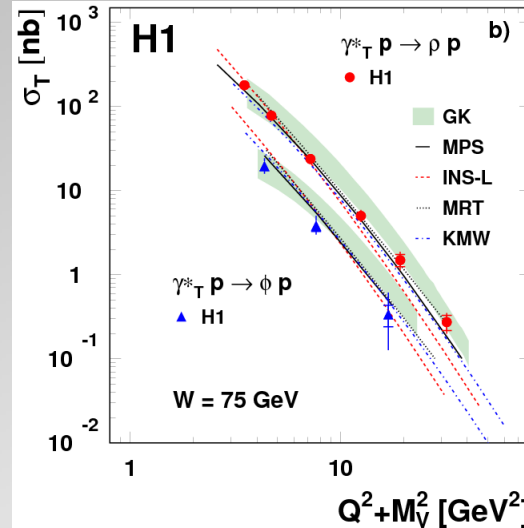
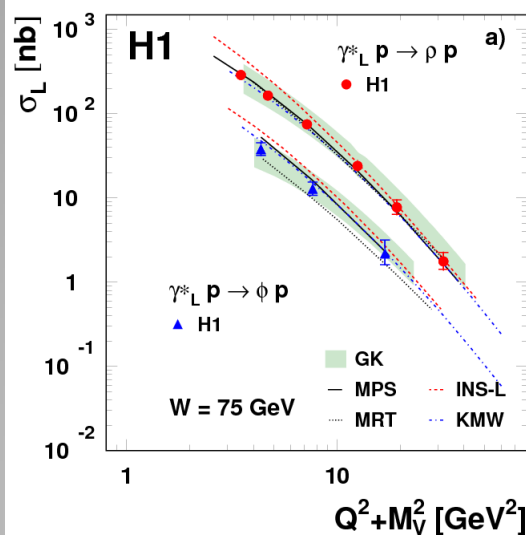
*27.5 GeV electrons/positrons with 920 GeV protons beam*  
*CMS energy = 318 GeV*





# $\sigma_L, \sigma_T$ cross sections

- Unique opportunity to extract  $R = \sigma_L / \sigma_T$   $\sigma = \sigma_T + \epsilon \sigma_L$   $\epsilon \approx 0.996$
- $R$  measured from angular distributions  $f(\cos \theta_H, r^{04}_{00})$ , in SCHC approximation



**expectation:  $R = \sigma_L / \sigma_T \sim Q^2 / M^2$**   
**as the scale gets harder  $\sigma_L$  dominates**

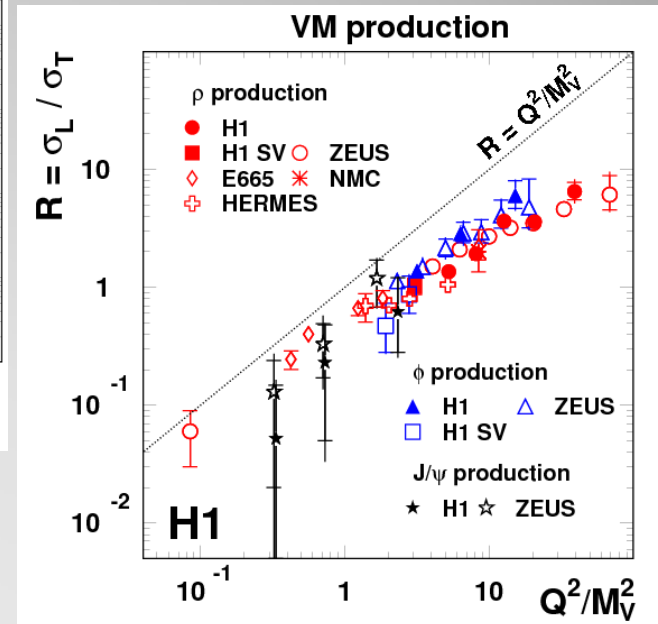
$\sigma_L$  and  $\sigma_T$  have different  $Q^2 + M^2$  dependence

Models based on pQCD describe well  $\sigma_L$ , but not  $\sigma_T$

**INS** - I.P. Ivanov, N.N. Nikolaev and A.A. Savin, Phys.Part.Nucl. 37 (2006) 1

**GK** - S.V. Goloskokov and P. Kroll, arXiv:hep-ph/0708.3569 (2007)

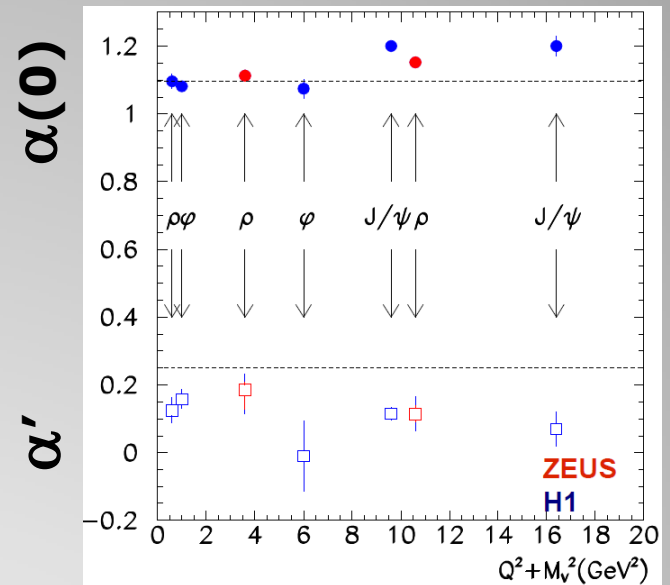
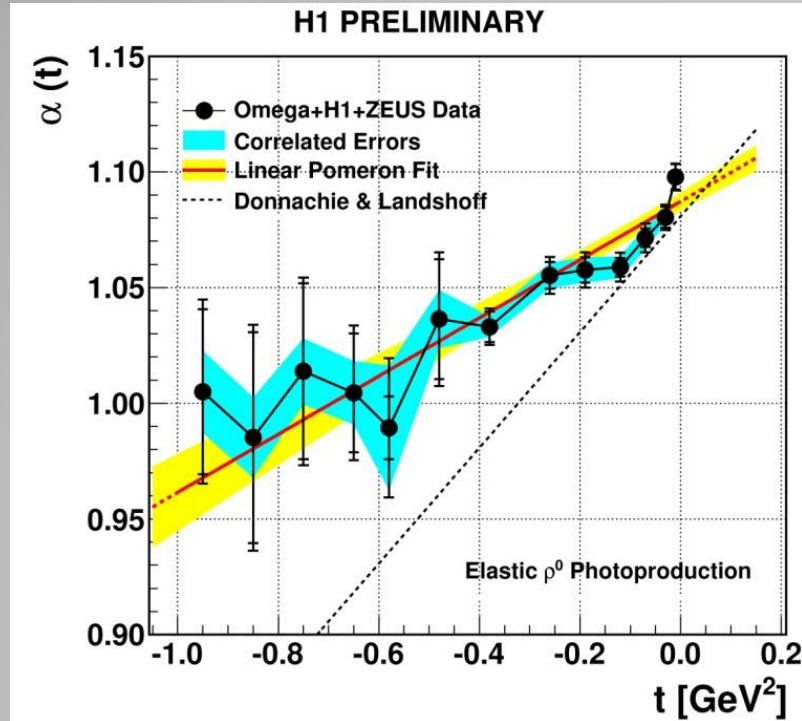
**MPS** - C. Marquet, R. Peschanski and G. Soyez, Phys. Rev. D76 (2007) 034011



# effective Pomeron trajectory

$$\frac{d\sigma}{d|t|} \propto \left(\frac{W}{W_0}\right)^{4(\alpha(t)-1)}$$

$$\alpha(t) = \alpha(0) + \alpha' t$$



$$DL: \alpha(t) = 1.08 + 0.25t$$

$$\alpha(0) = 1.0871 \pm 0.0026 \pm 0.0030$$

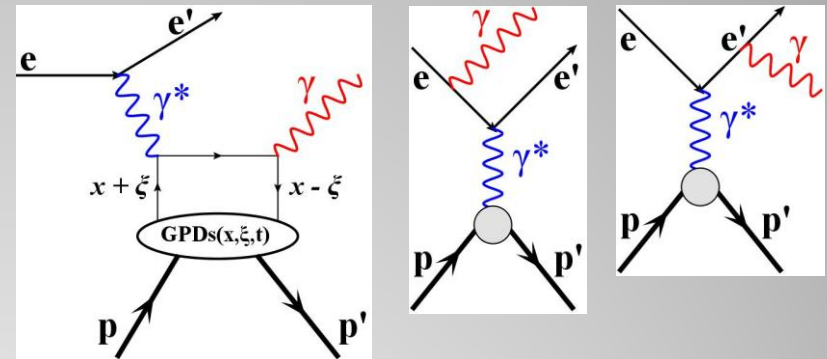
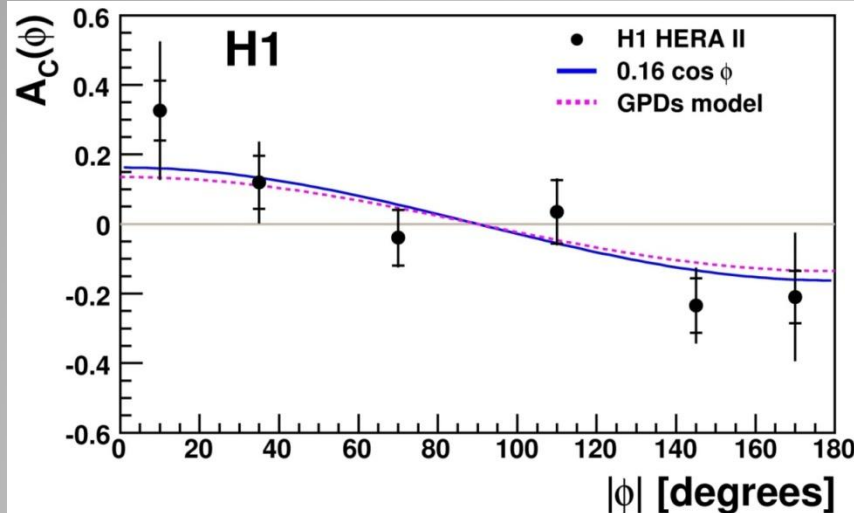
$$\alpha' = 0.126 \pm 0.013 \pm 0.012 \text{ GeV}^{-2}$$

**$\alpha(0)$  consistent with 1.08 from soft pp scattering**  
 **$\alpha' \sim$  twice smaller than  $0.25 \text{ GeV}^{-2}$**

# DVCS – Beam Charge Asymmetry

DVCS and Bethe-Heitler have the same initial and final states and they are indistinguishable. The QCD-QED interference term is sensitive to the real part of the QCD amplitude. It changes sign with lepton beam charge:

$$\sigma = \sigma_{DVCS} + \sigma_{BH} \pm \sigma_{interf.}$$



$x_{1,2} = x \pm \xi$  – parton longitudinal momentum fraction,  $\xi$  – fraction of the momentum transfer

$$A_C = \frac{\sigma^+ - \sigma^-}{\sigma^+ + \sigma^-} = p_1 \cos(\phi) + \dots$$

$$p_1 = 0.16 \pm 0.04 \pm 0.06$$

$\phi$  is the angle between two planes defined by incoming and outgoing electron and  $\gamma^*$  and outgoing proton

DVCS gives access to Generalized Parton Distributions (GPD), which describe the correlations between two partons ( $x_1, x_2$ ) which differ by longitudinal ( $x_1 \neq x_2$ ) and transverse ( $t$ ) momentum at given  $Q^2$

**GPD's based model compatible with data**