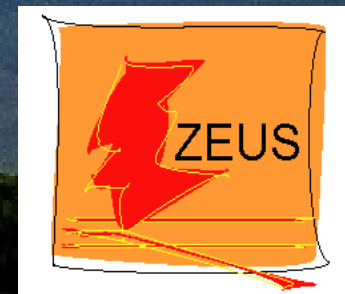


The 7<sup>th</sup> Joint International HADRON STRUCTURE'13 Conference  
30 June – 4 July, Tatranské Maltiare, Slovakia

# Precision QCD measurements at HERA: *Jet production and $\alpha_s$ determination*

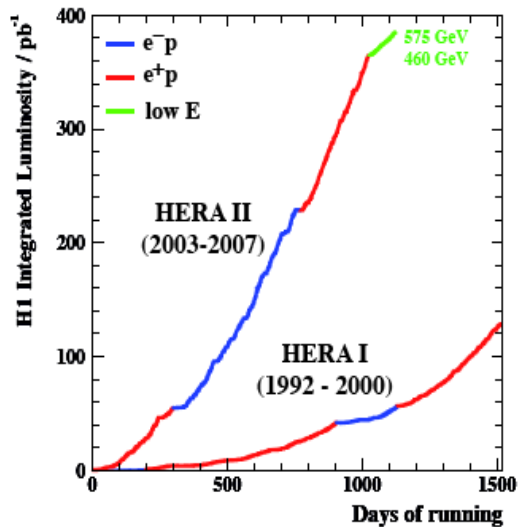
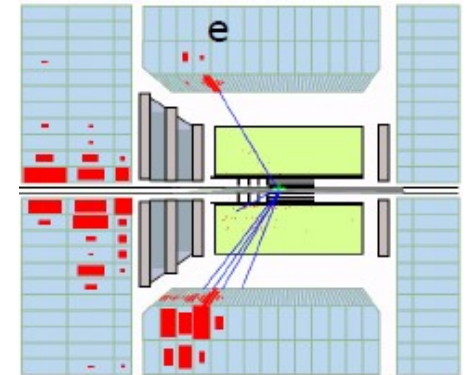
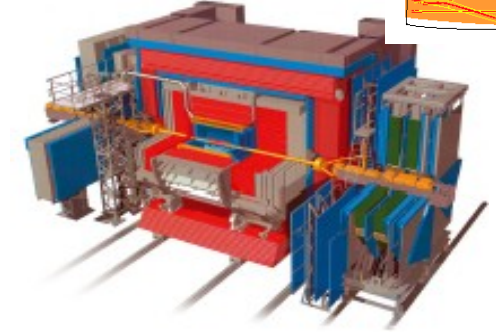
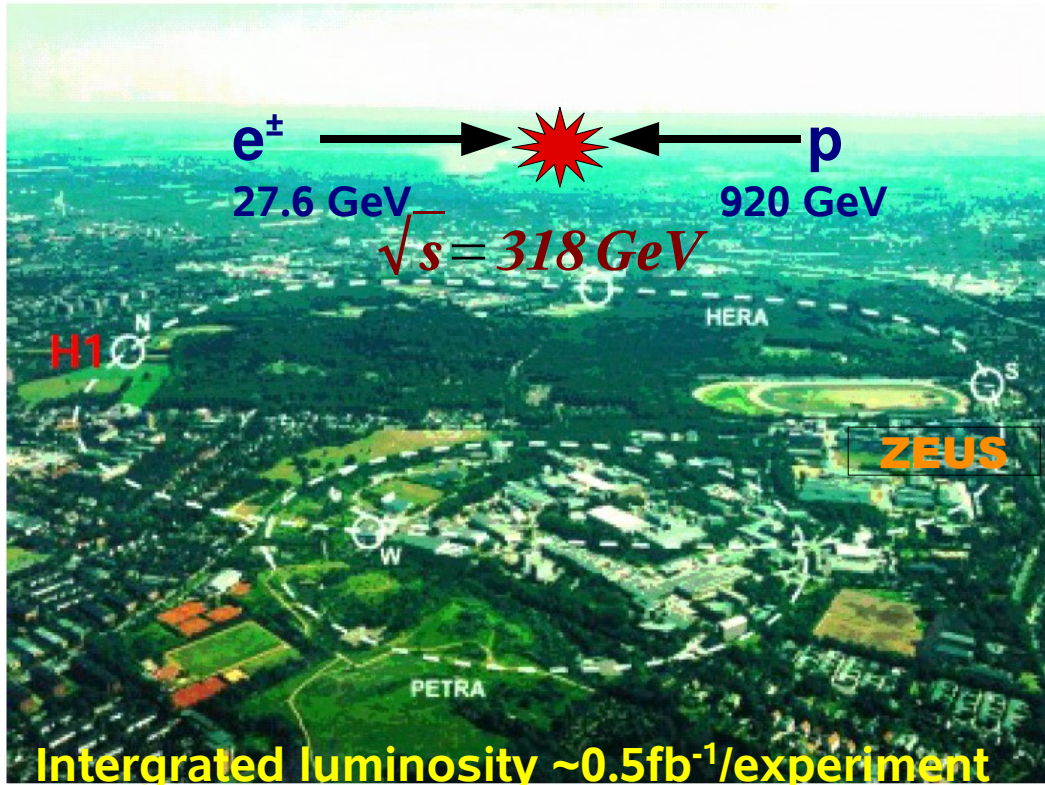
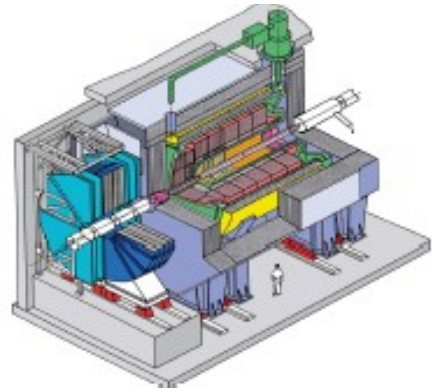
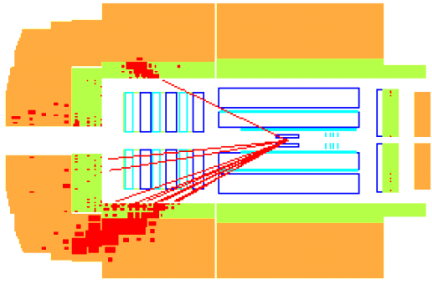
Stanislaw Mikocki

Institute of Nuclear Physics PAN, Cracow  
*on behalf of the H1 and ZEUS collaborations*





# HERA Experiments



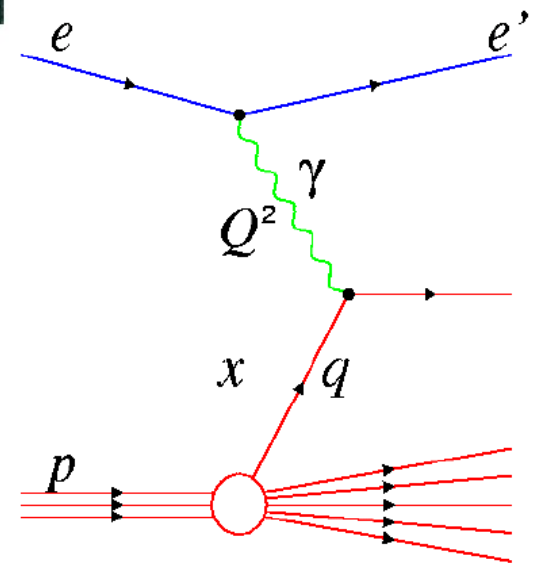
Stanislaw Mikocki

## Standard DIS variables :

- $Q^2$  virtuality of the exchanged boson
- $x$  in QPM fraction of proton momentum carried by struck quark
- $y = Q^2 / xs$  inelasticity

## Kinematics regimes:

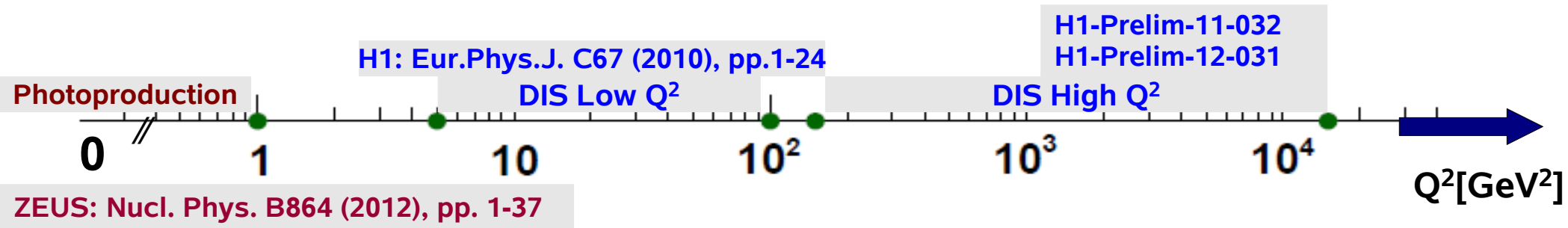
- $Q^2 \approx 0 \text{ GeV}^2$  – Photoproduction ( $\gamma p$ )
- $Q^2 > 1 \text{ GeV}^2$  - DIS



# Overview

## Part I

- jets production in large range of  $Q^2$ : inclusive jets(PHP) and multijets(DIS)
- measurements of the cross sections and extraction of  $\alpha_s(M_Z)$
- comparison with NLO prediction

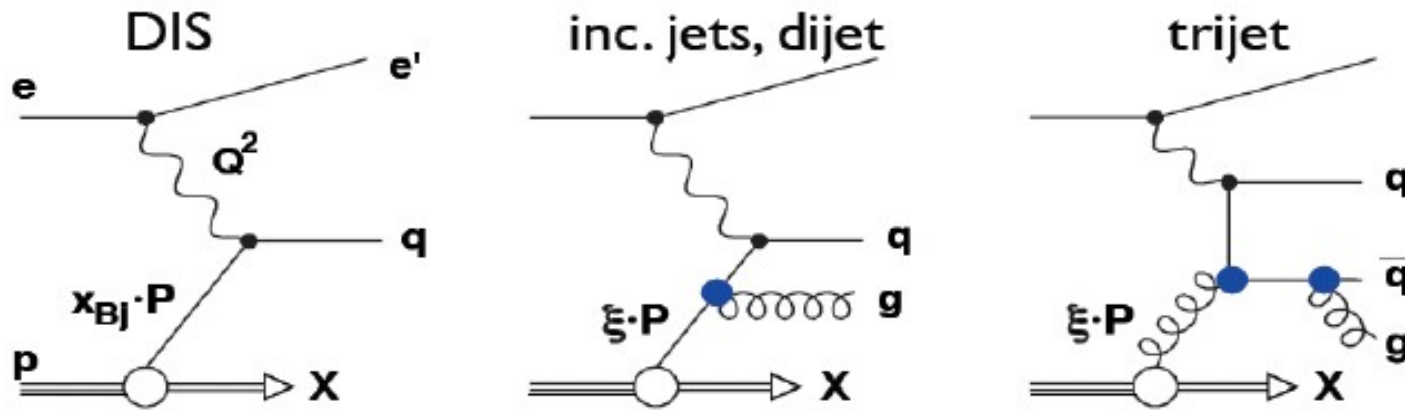


## Part II

H1: Eur. Phys. J. C72 (2012) 1910

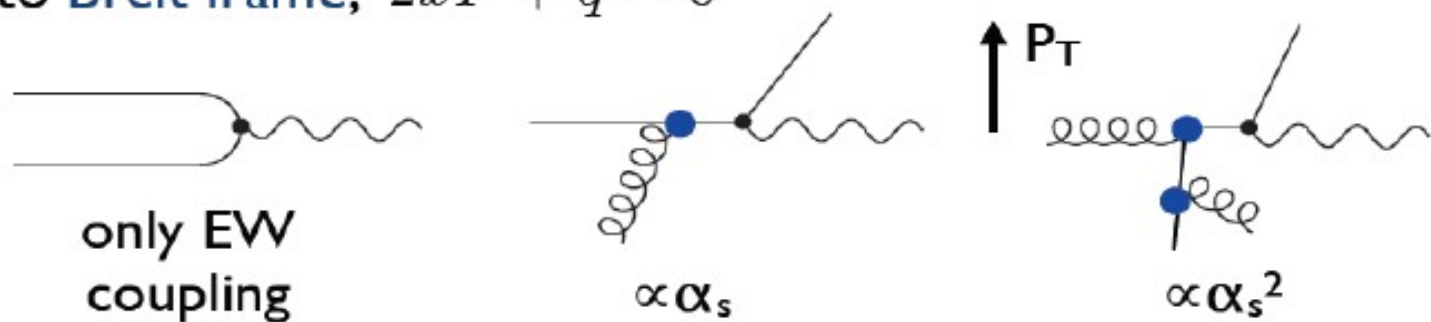
Test of QCD evolution mechanisms (*DGLAP* / *BFKL* / *CCFM*)  
using azimuthal correlation  
between the most forward jet and the scattered positron in DIS

# Jet Production in DIS



The momentum fraction of the proton carried by the parton entering the hard subprocess :  $\xi = x_{Bj} (1 + M_{12}^2 / Q^2)$

Boost to Breit frame,  $2xP + q = 0$



Only hard QCD process generates significant  $P_T$  in Breit frame  
Direct sensitivity to  $\alpha_s$  and gluon PDF

# Multi-Jet Production in DIS



Low  $Q^2$

High  $Q^2$

## NC DIS and Jet selections

HERA-1 data  $44\text{pb}^{-1}$

HERA-2 data  $350\text{pb}^{-1}$

NC DIS Selection	$5 < Q^2 < 100 \text{ GeV}^2, 0.2 < y < 0.7$		
Inclusive jet	$P_T > 5 \text{ GeV}$		$-1.0 < \eta_{\text{Lab}}^{\text{jet}} < 2.5$
2-jet	$P_T^{\text{jet}1}, P_T^{\text{jet}2} > 5 \text{ GeV}$	$M_{12} > 18 \text{ GeV}$	
3-jet	$P_T^{\text{jet}1}, P_T^{\text{jet}2}, P_T^{\text{jet}3} > 5 \text{ GeV}$		

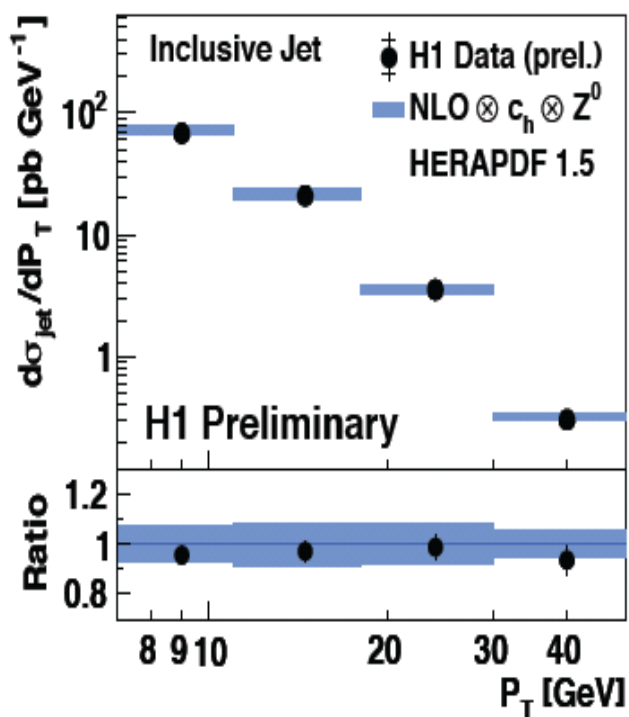
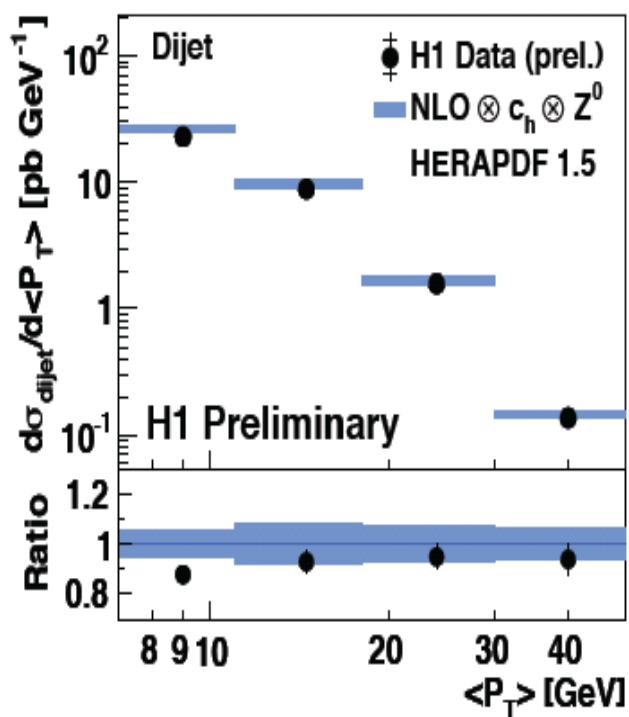
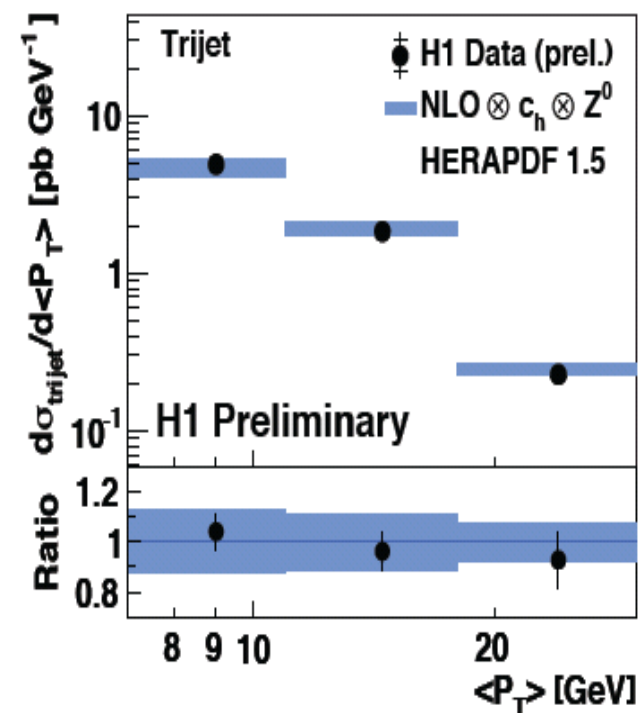
NC DIS Selection	$150 < Q^2 < 15000 \text{ GeV}^2, 0.2 < y < 0.7$		
Inclusive jet	$7 < P_T < 50 \text{ GeV}$		$-1.0 < \eta_{\text{lab}} < 2.5$
Dijet	$5 < P_T^{\text{jet}1}, P_T^{\text{jet}2} < 50 \text{ GeV}$	$M_{12} > 16 \text{ GeV}$	
Trijet	$5 < P_T^{\text{jet}1}, P_T^{\text{jet}2}, P_T^{\text{jet}3} < 50 \text{ GeV}$		

- Cross sections are measure as function of  $Q^2$ ,  $p_T(<p_T)$  and  $\xi$
- main experimental uncertainties
  - jet energy scale 2%  $\rightarrow \Delta\sigma/\sigma = 4\text{-}10\%$
  - uncertainty in acceptance  $\rightarrow \Delta\sigma/\sigma = 2\text{-}15\%$
- NLO calculation: NLOJET++
  - MSbar scheme for 5 massless quark flavors
  - PDFs: CTEQ6.5M

- Cross sections are measure as function of  $Q^2$ ,  $p_T(<p_T)$  and  $\xi$
- main experimental uncertainties
  - jet energy scale 1%  $\rightarrow \Delta\sigma/\sigma = 3\text{-}10\%$
  - uncertainty in acceptance  $\rightarrow \Delta\sigma/\sigma = 4\text{-}5\%$
- NLO calculation: NLOJET++
  - MSbar scheme for 5 massless quark flavors
  - PDFs: HERAPDF1.5, CT10



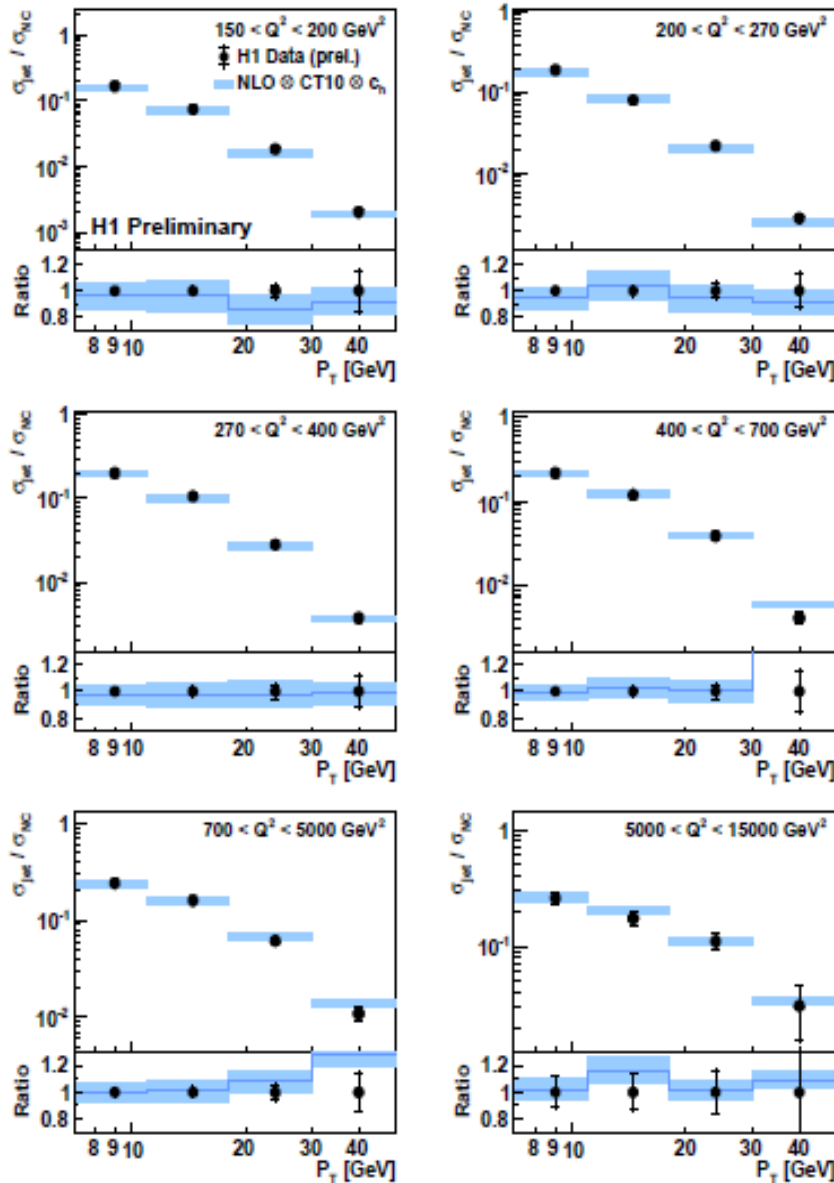
## Single Differential Cross Sections

**Incl. Jet**

**2- Jet**

**3- Jet**


NLO QCD with  $\mu_r = \sqrt{(Q^2 + P_T^2)}/2$  and HERAPDF 1.5 describes well inclusive jet, dijet and trijet single differential cross sections



## Double Differential Inclusive Jet Cross Sections



### *Benefit:*

partial cancellation of  
experimental and  
theoretical uncertainties

### *Comparison with*

NLOJet++ and QCDNUM  
corrected to hadronisation  
effects

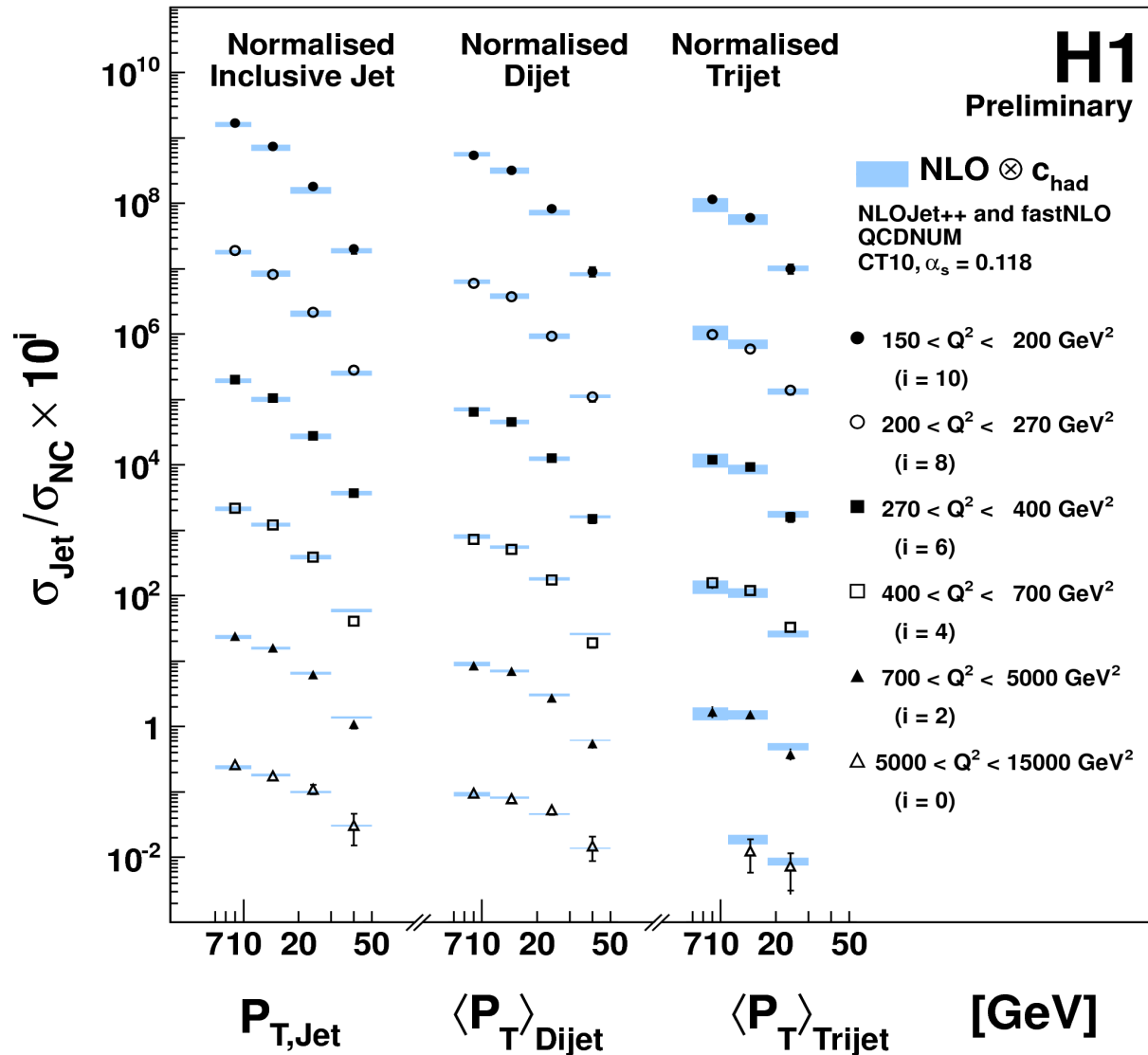
### *Scale choice:*

$$\mu_f^2 = Q^2,$$

$$\mu_r^2 = (Q^2 + P_T^2)/2$$

In all bins (besides the highest  $Q^2$  and highest  $P_T$ )  
the experimental uncertainties are smaller than the theoretical uncertainties

# Normalized Multi-Jet Cross Sections at High $Q^2$



## NLO Calculation:

NLOJet++ and QCDNUM corrected for hadronisation effects

## Scale Choice:

$$\mu_f^2 = Q^2$$

$$\mu_r^2 = (Q^2 + P_T^2)/2$$

- Small experimental uncertainties
- Good NLO description of the data





## $\alpha_s(M_Z)$ Combined Fit

### Largest benefit is from a combined fit

simultaneous fit to normalised inclusive jet, dijet and trijet cross sections  
(all correlations are included)

### Sensitive to higher orders

Theoretical uncertainties estimated by variation of scale,  
k-factor ( $k = \sigma_{NLO}/\sigma_{LO}$ ) – an estimator of higher order contributions  
reaches values up to 1.45

### Restrict analysis to $k < 1.3$

faster convergence of perturbative series  
trade-off between number of data points and smaller theoretical uncertainties

*Normalised Multijets with  $k < 1.3$*

$\chi^2/\text{ndf}: 53.2/41 = 1.30$

$$\alpha_s(M_Z) = 0.1163 \pm 0.0011(\text{exp.}) \pm 0.0008(\text{had})_{-0.0035}^{+0.0044}(\text{th.}) \pm 0.0014(\text{PDF})$$

Consistent with other  $\alpha_s(M_Z)$  measurements

Small experimental uncertainties

Theoretical uncertainties are larger than the experimental

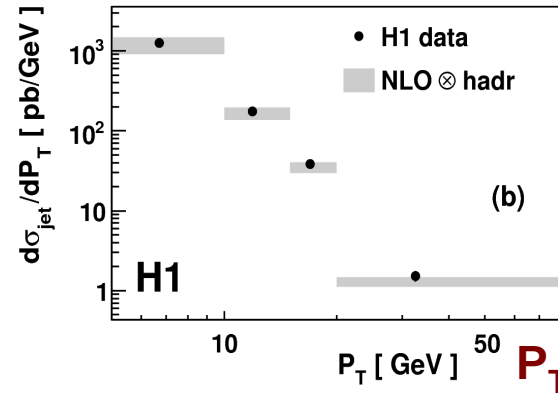
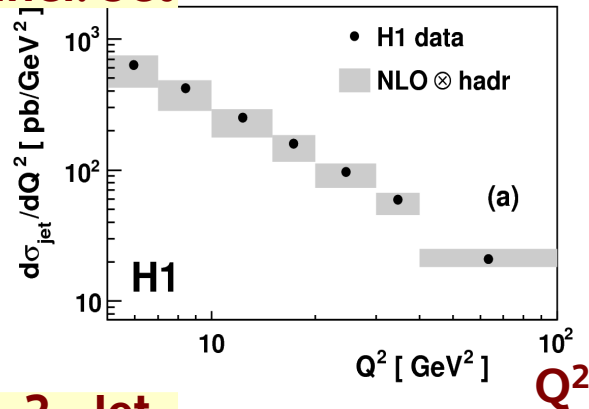
# Multi-Jet Cross Sections at Low $Q^2$

Eur.Phys.J. C67 (2010), pp.1-24

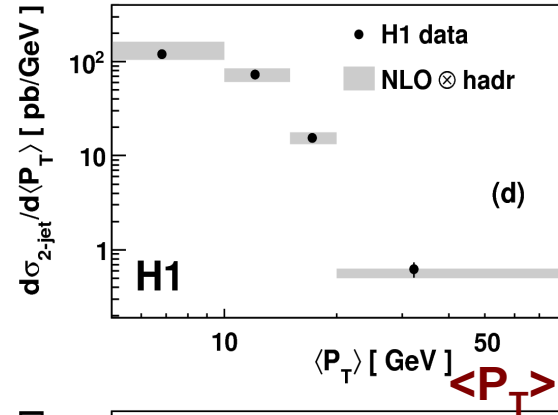
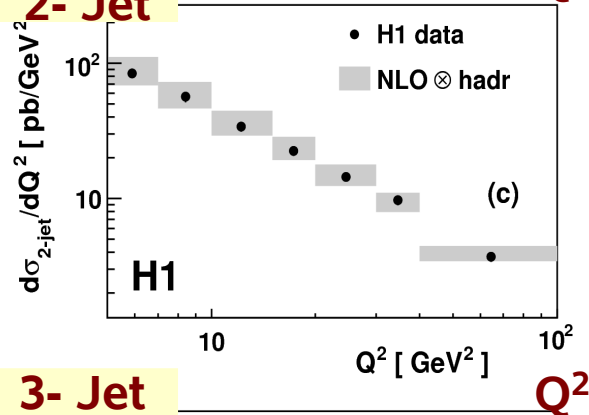


## Inclusive Jet, 2-Jet and 3-Jet Cross Sections

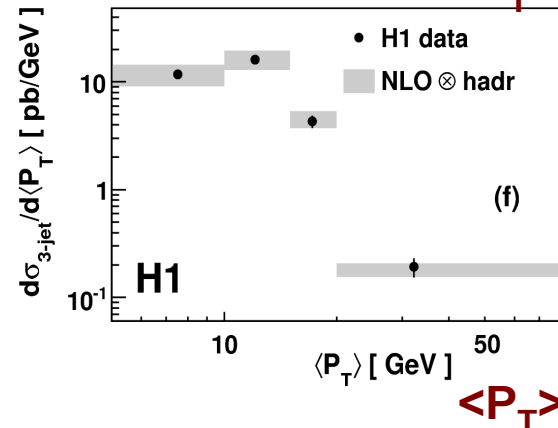
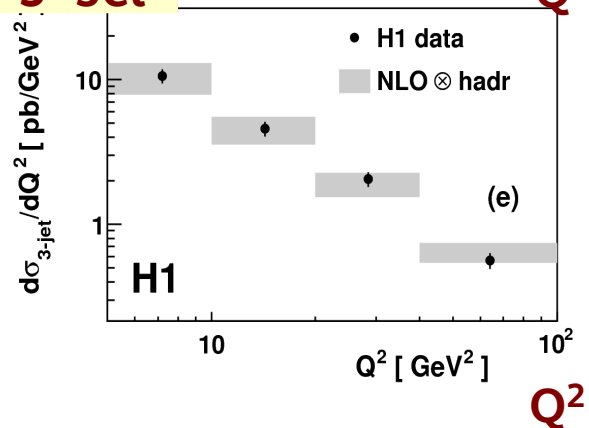
### Incl. Jet



### 2- Jet



### 3- Jet

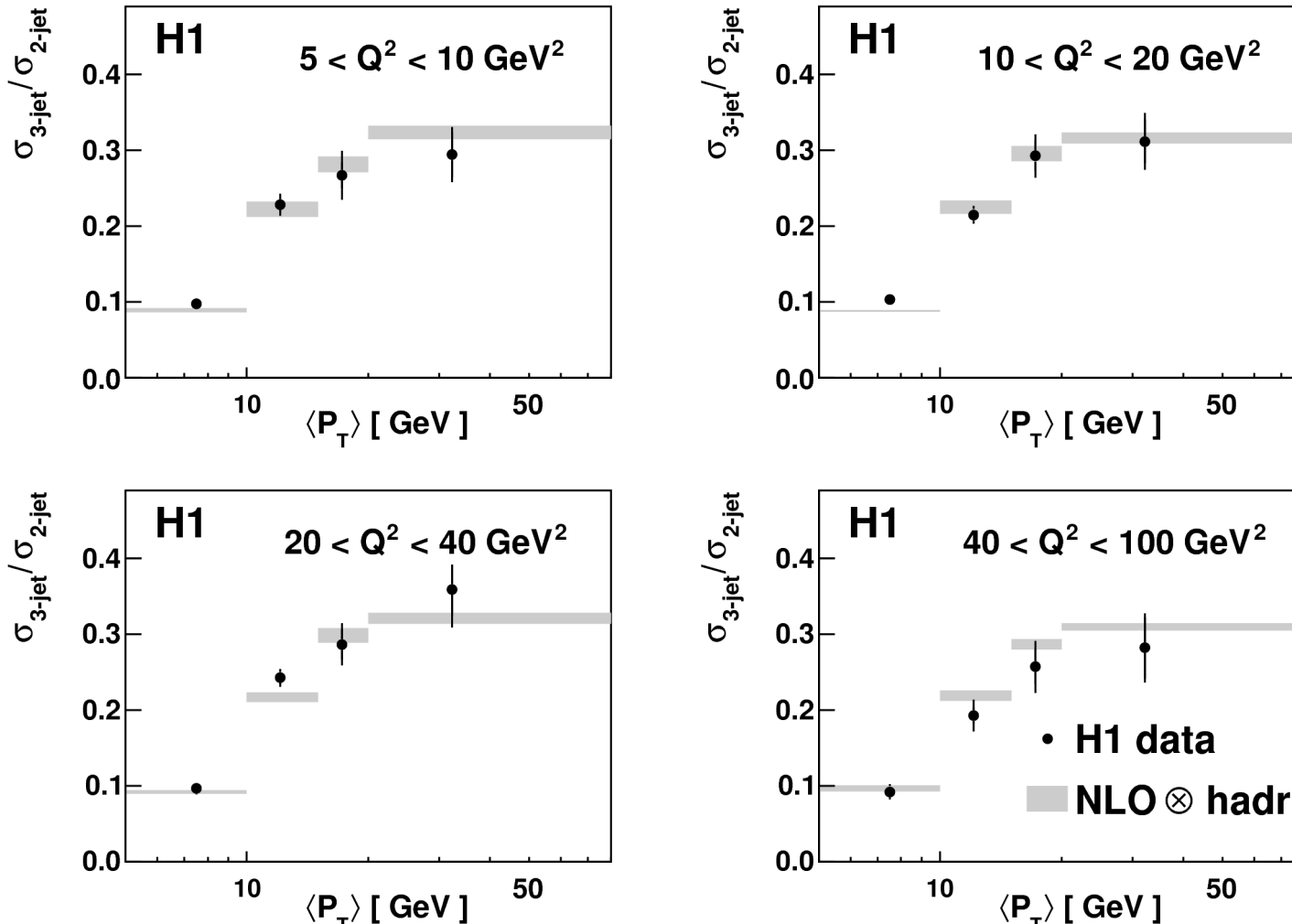


- measurements are well described by NLO
- Experimental uncertainty 6-11%
- theory uncertainty dominated by renorm. scale uncertainty: 10% (highest  $Q^2$  and  $P_T$ ) to 30% (lowest  $Q^2$  and  $P_T$ )
- pdf uncertainty 2-6%
- low predictive power of NLO at low  $Q^2$  and/or low  $P_T$
- orders beyond NLO are needed to match the precision of data

# Multi-Jet Cross Sections at Low $Q^2$



## 3-Jet to 2-Jet Ratio



- in ratio normalisation errors cancel and other syst. Uncertainties reduced by 50%
- reduced sensitivity to renormalisation scale variation in theory
- good description of ratio by NLOjet++

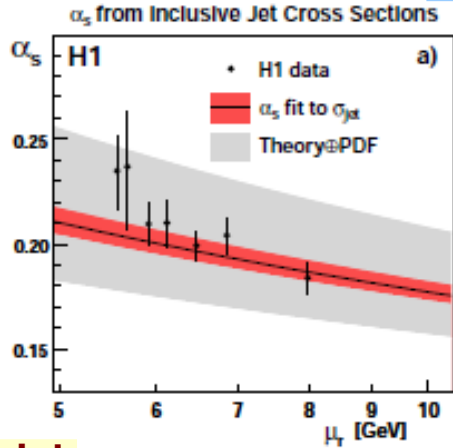
# $\alpha_s$ from Multi-Jet Cross Sections at Low $Q^2$



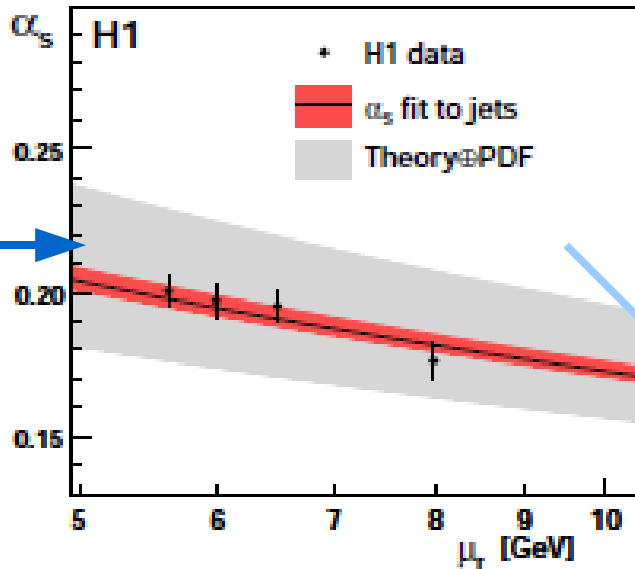
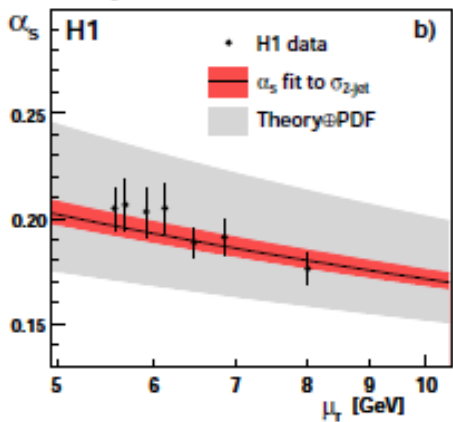
## Incl. Jet

## Combined Fit Incl.Jet+2-jet+3-jet :

$$\alpha_s(M_Z) = 0.1160 \pm 0.0014 (exp.)_{-0.0077}^{+0.0093} (th.) \pm 0.0016 (PDF)$$

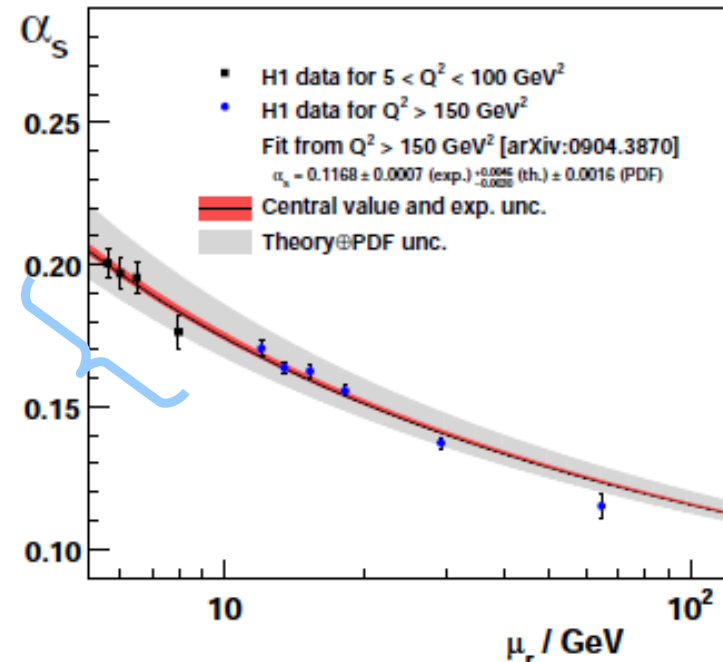


## 2- Jet



## $\alpha_s$ from Jet Cross Sections in DIS

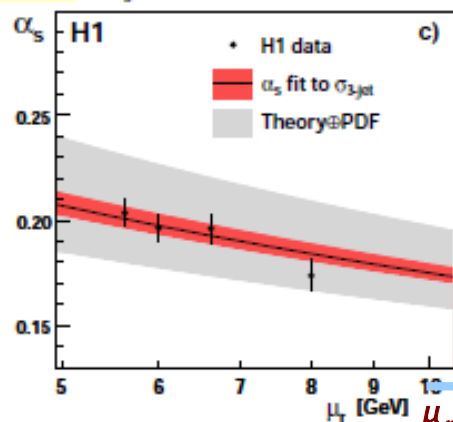
(High  $Q^2$  results: Eur.Phys.J.C65 (2010) 363, *not this talk*)



Experimental uncertainties reduced significantly wrt individual  $\alpha_s$  extraction

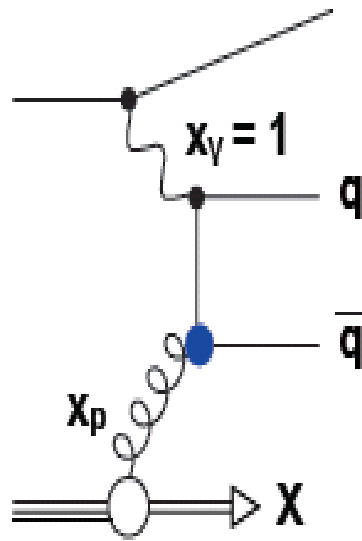
Remarkable agreement between low, high  $Q^2$   $\alpha_s$  extraction and QCD expectations

## 3- Jet

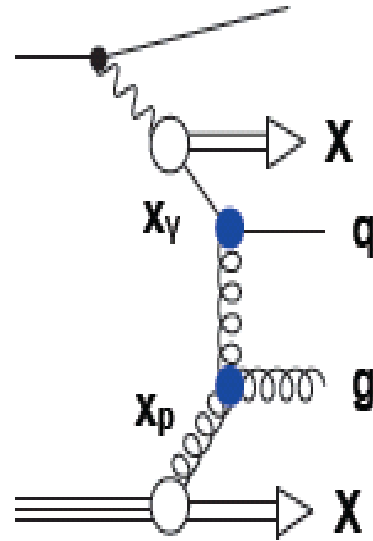


$$\mu_r = \mu_f = \sqrt{(Q^2 + p_{T,jet}^2)/2}$$

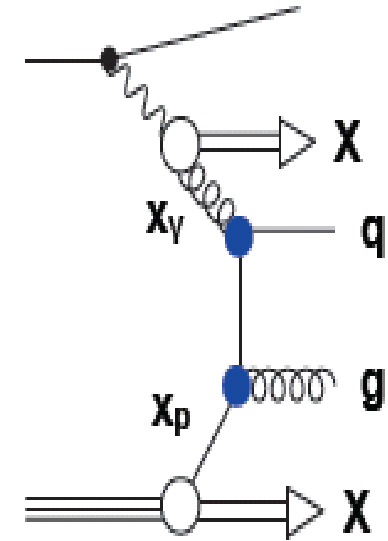
# Jet Production in Photoproduction



direct photoproduction



resolved photoproduction



Direct sensitivity to  $\alpha_s$ , gluon and photon PDFs

# Jet Production in Photoproduction



Data:  $\sim 300 \text{ pb}^{-1}$  (HERA-2)

Single and double differential inclusive jet cross sections are measured as functions of jet transverse energy  $E_T^{\text{jet}}$  and pseudorapidity  $\eta^{\text{jet}}$  for

photon virtuality:

$$Q^2 < 1 \text{ GeV}^2$$

$\gamma p$  centre-of-mass energies:  $142 < W_{\gamma p} < 293 \text{ GeV}$

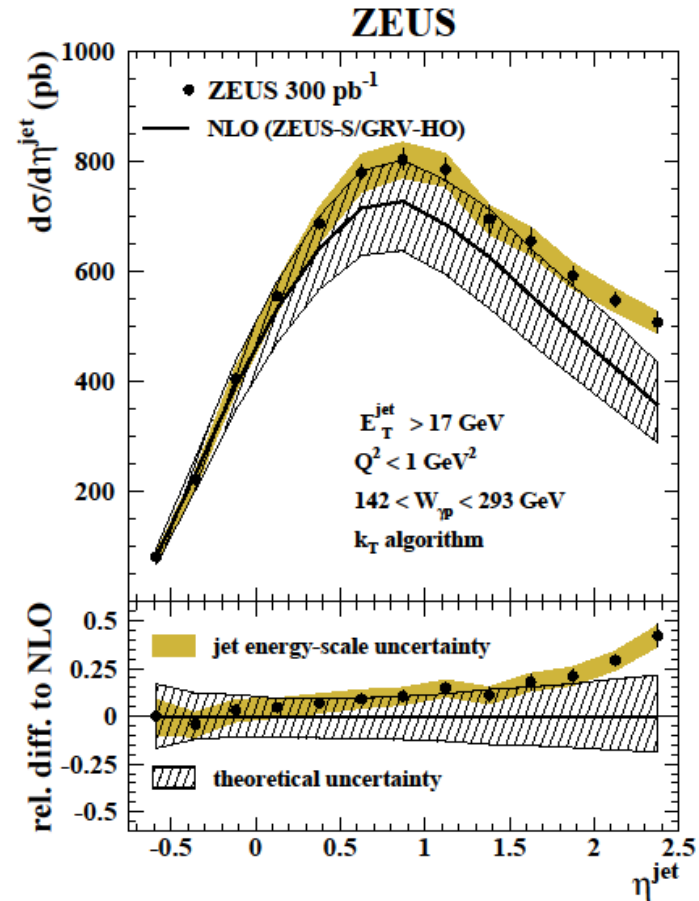
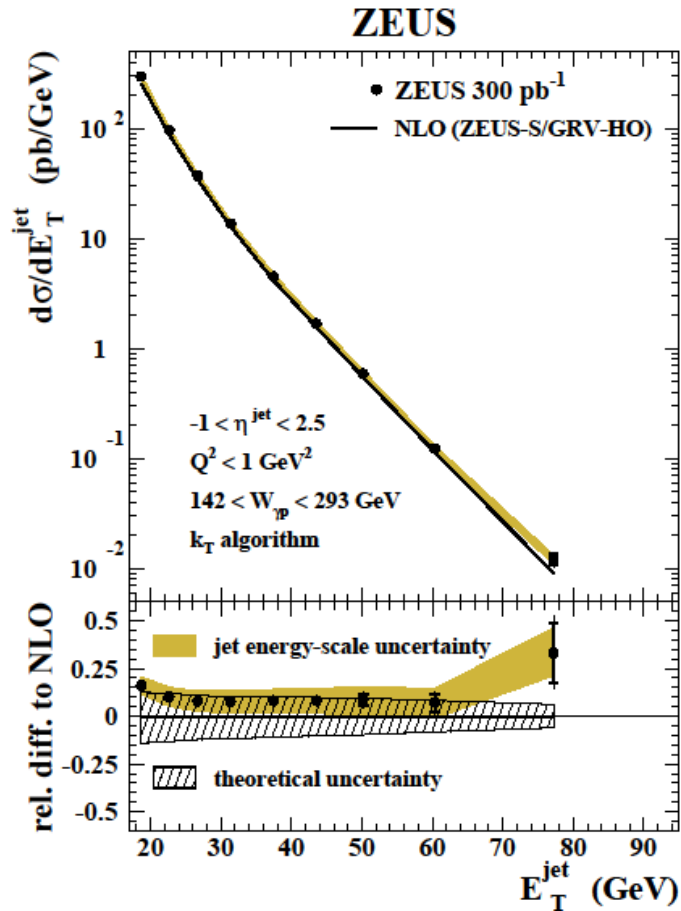
Jets in lab frame:

$$E_T^{\text{jet}} > 17 \text{ GeV}$$

$$-1 < \eta^{\text{jet}} < 2.5$$

Jets were identified using the  $k_T$ , anti- $k_T$  and SIScone jet algorithms in laboratory frame.

# Inclusive Jets in Photoproduction

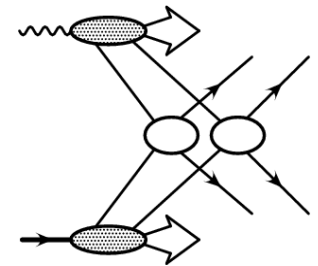
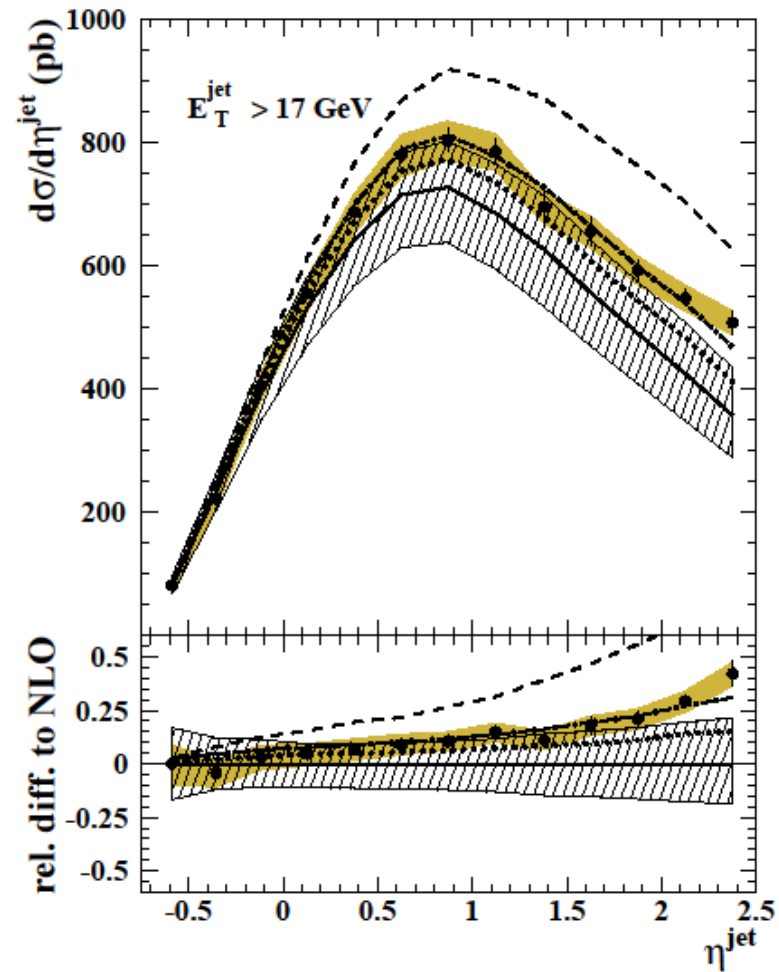
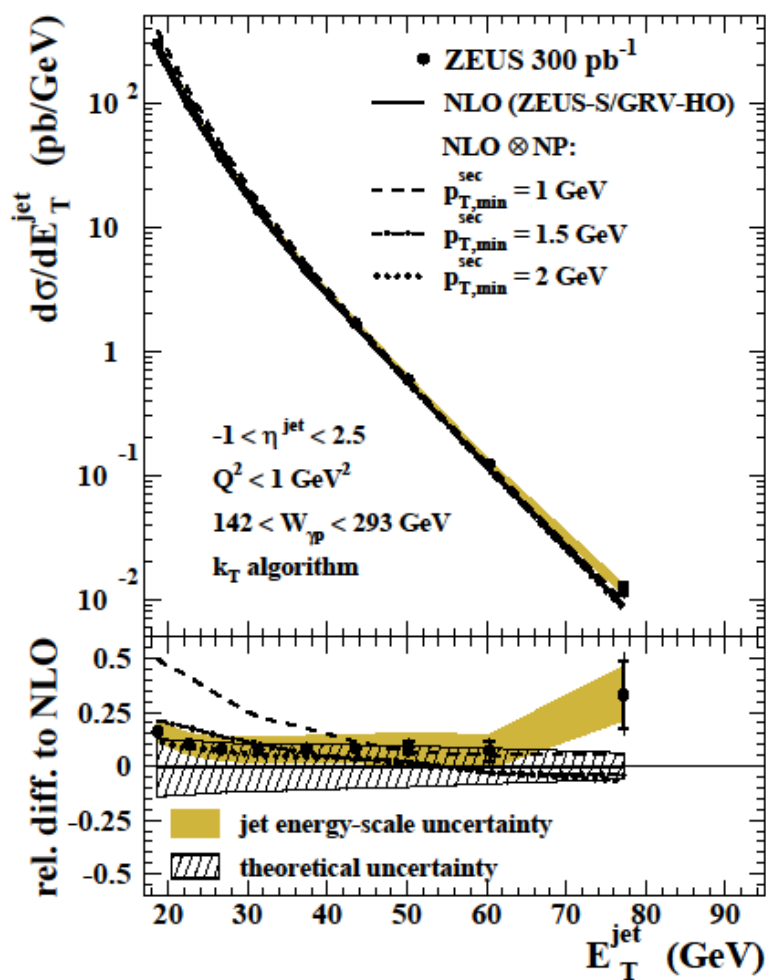


Data compared to NLO QCD ( $O(\alpha_s^2)$ ):

$$\mu_R = \mu_F = \mu = E_T^{\text{jet}}$$

PDFs: proton PDF -ZEUS-s, photon PDF – GRV-HO,  $\alpha_s = 0.118$

The NLO QCD calculation reproduce  $d\sigma/dE_T^{\text{jet}}$  well,  $d\sigma/d\eta^{\text{jet}}$  is well described for  $\eta^{\text{jet}} < 2$



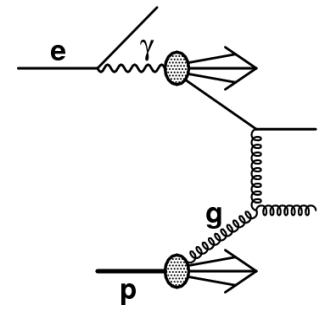
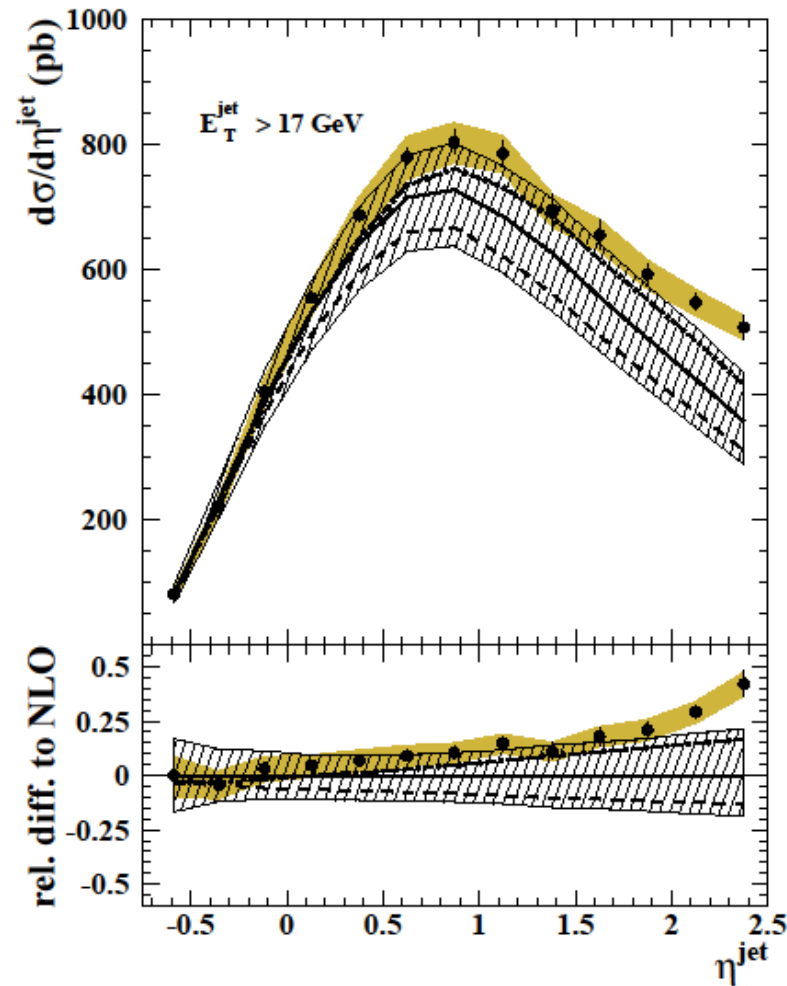
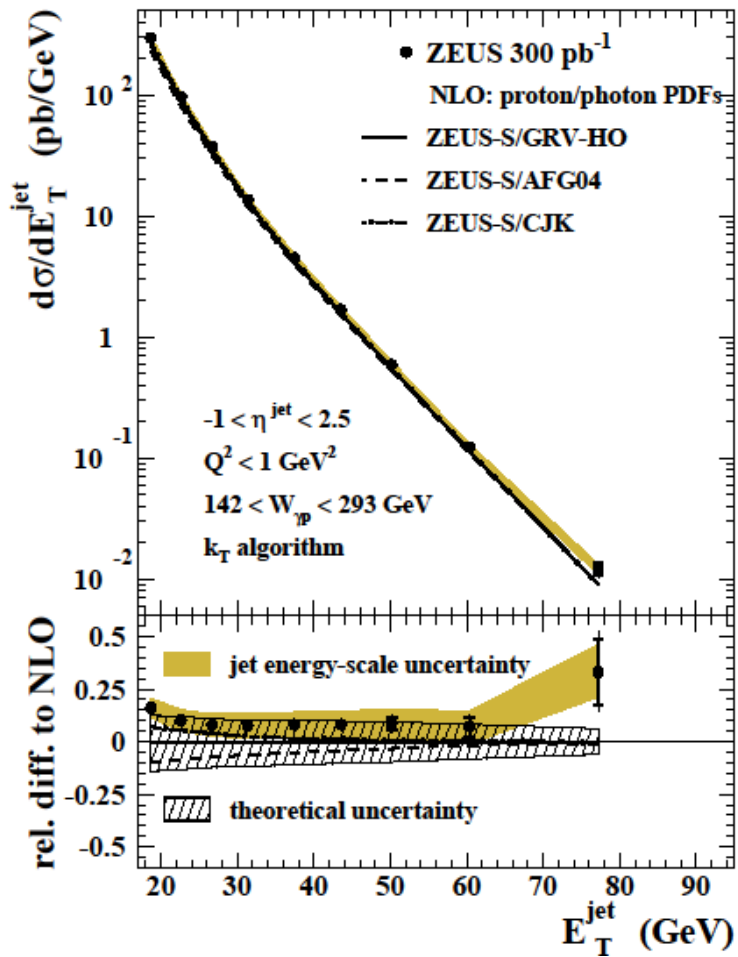
Multi-parton interactions not included in NLO calculation

Data comparison to the NLO QCD calculation including an estimation of non-perturbative effects from underlying events (not related to hadronisation)

Possible presence of effects in the data, which are not included in the NLO QCD calculation



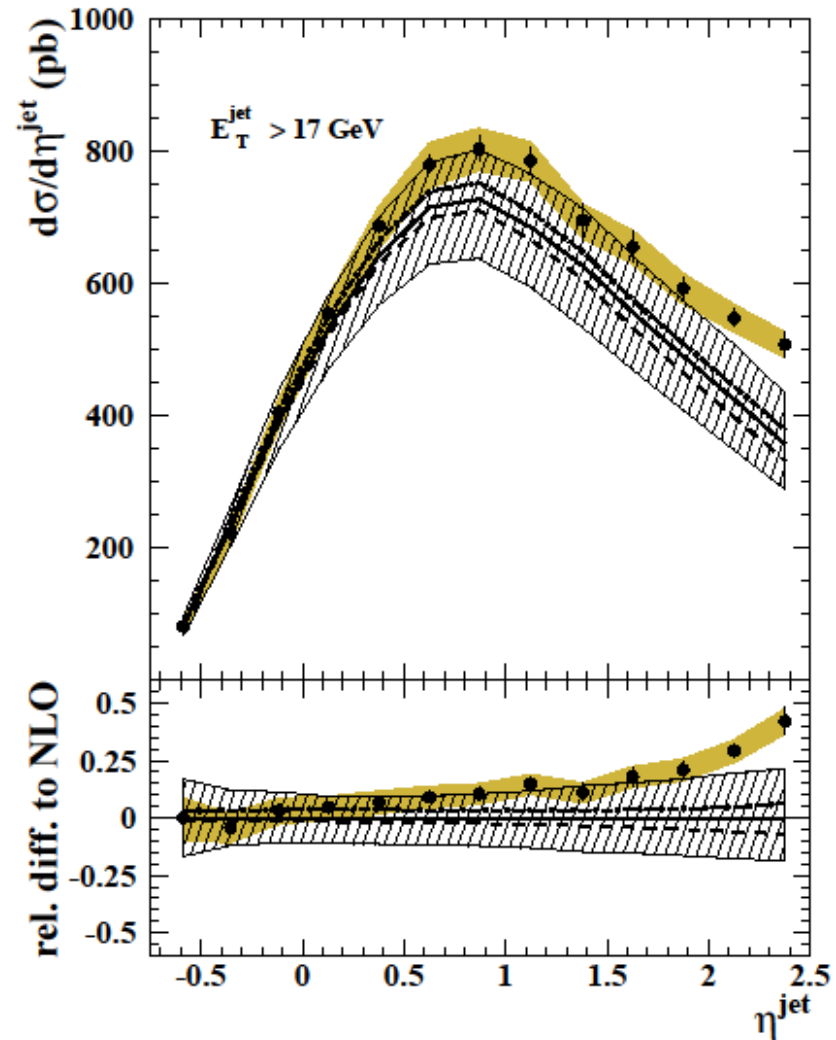
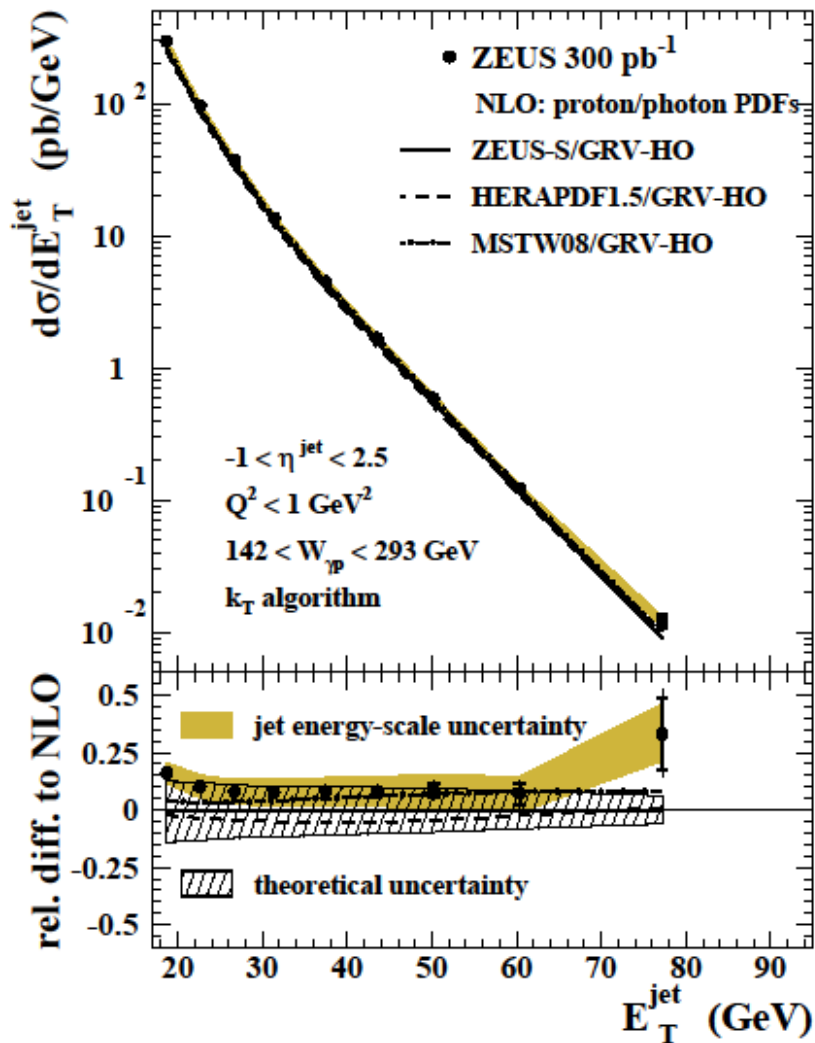
## Dependence on photon PDFs



Some difference between three predictions, especially at low  $E_T^{\text{jet}}$  and high  $\eta^{\text{jet}}$

Potential to constrain photon PDFs

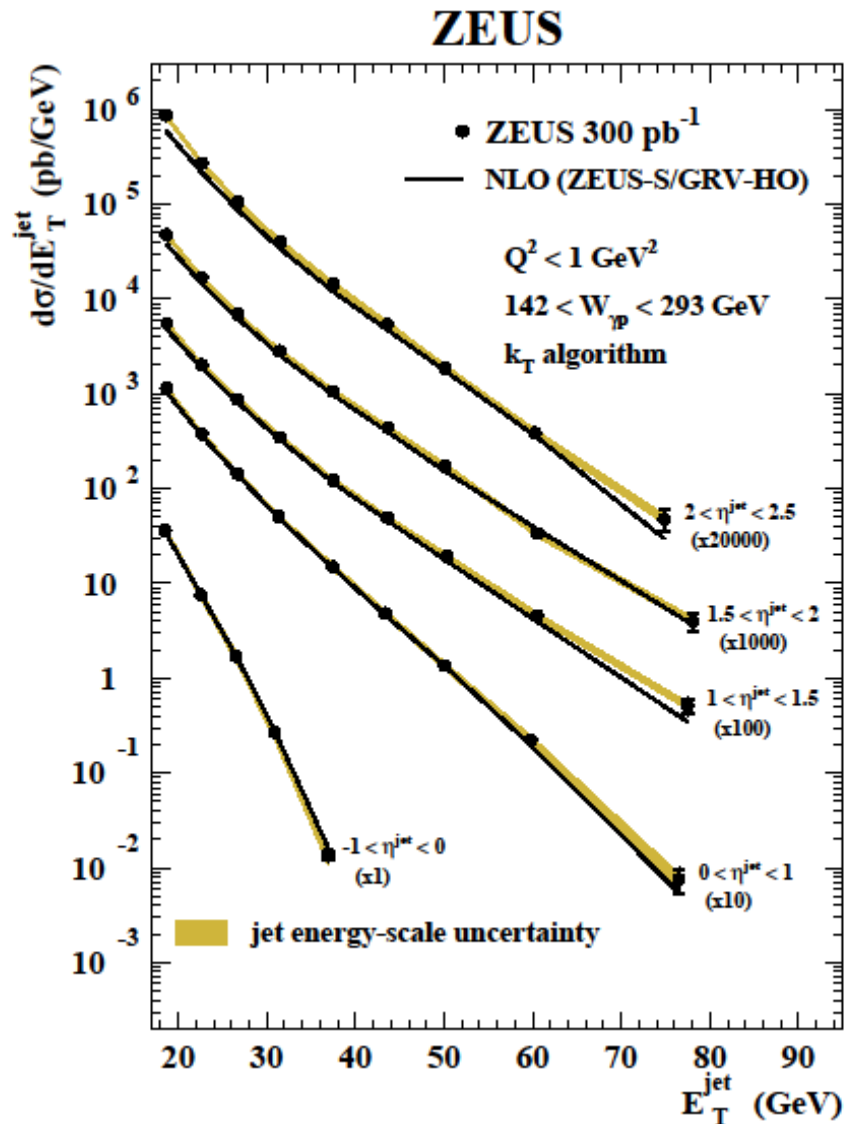
## Dependence on proton PDFs



Small difference between three predictions

Low sensitivity to proton PDFs

# Inclusive Jets in Photoproduction

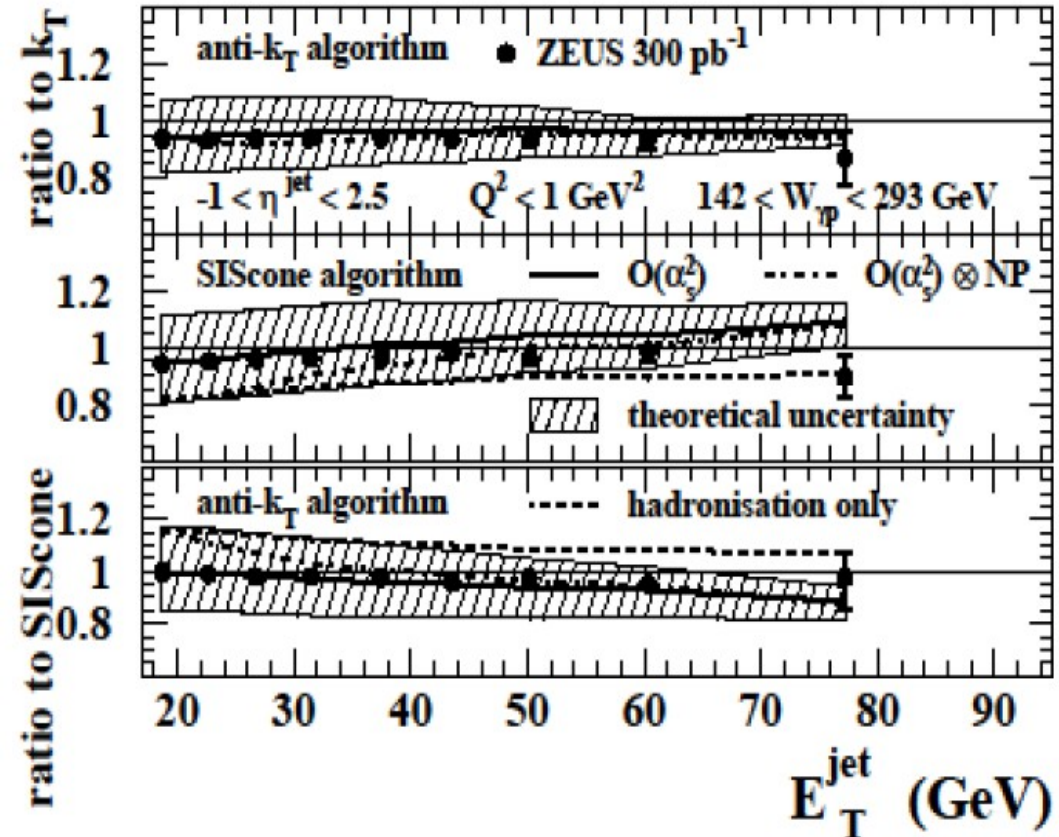
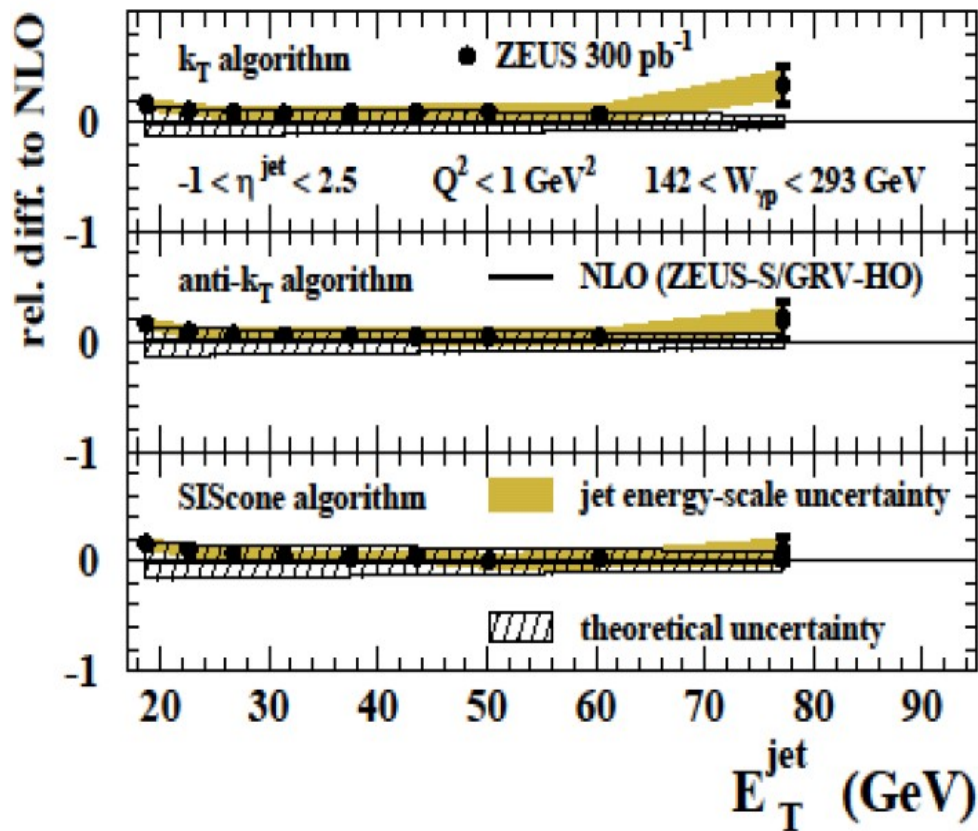


Differential cross section based on  $k_T$  jet algorithm for inclusive jet photoproduction with  $E_T^{\text{jet}} > 17 \text{ GeV}$  in different  $\eta^{\text{jet}}$  regions.

Difference between data and NLO at large  $\eta^{\text{jet}}$  and low  $E_T^{\text{jet}}$  could be from photon PDFs or non-perturbative effects

The NLO QCD predictions give a good description of the data , except at (low  $E_T^{\text{jet}}$  and high  $\eta^{\text{jet}}$  )

## NLO QCD and Jet Algorithms Comparison



- the agreement of the data to the NLO prediction is the same for all three jet algorithms
- no sensitivity of the result on the choice of the jet algorithm used

## Determination of $\alpha_s(M_Z)$ and Energy scale dependence

The measured single differential cross sections based on the three jet algorithms were used to determine  $\alpha_s(M_Z)$  values.

To minimise the effects of a non-perturbative contributions and reduce uncertainties coming from proton PDFs only the measurements for  $21 < E_T^{\text{jet}} < 71$  GeV were used in the fit.



$\alpha_s(M_Z)$  obtained from presented data are:

$$\alpha_s(M_Z)|_{k_T} = 0.1206^{+0.0023}_{-0.0022} \text{ (exp.) } ^{+0.0042}_{-0.0035} \text{ (th.)},$$

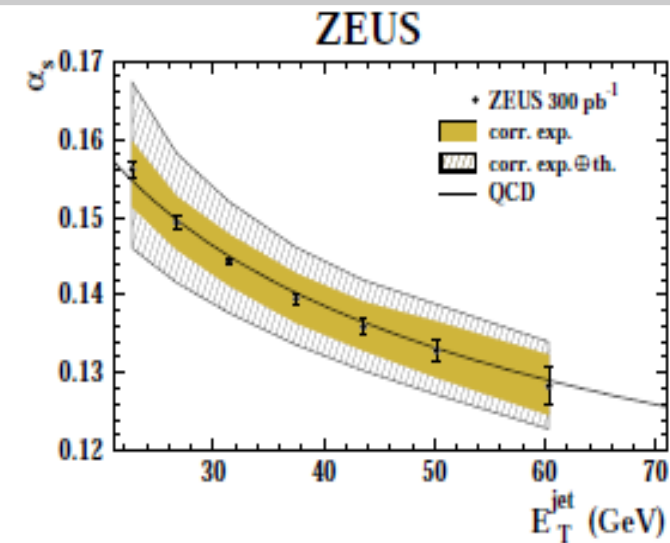
$$\alpha_s(M_Z)|_{\text{anti-}k_T} = 0.1198^{+0.0023}_{-0.0022} \text{ (exp.) } ^{+0.0041}_{-0.0034} \text{ (th.)},$$

$$\alpha_s(M_Z)|_{\text{SIScone}} = 0.1196^{+0.0022}_{-0.0021} \text{ (exp.) } ^{+0.0046}_{-0.0043} \text{ (th.)}.$$

The value of  $\alpha_s(M_Z)$  determined from the  $k_T$ , anti- $k_T$  and SIScone measurements are nicely agreeing

These determinations are consistent with previous determinations of  $\alpha_s(M_Z)$  and have a precision comparable to those obtained from  $e^+e^-$  experiments

Energy-scale dependence of  $\alpha_s$



Running of  $\alpha_s$  in single experiment in good agreement with Renormalisation Group Equations prediction at 2-loops

# Comparison of $\alpha_s(M_Z)$ values

Uncertainties: exp. ——— theo. - - - - -

**EW Fit, Z decays, 4NLO**

Gfitter Group, EPJC 72, 2003 (2012)

**H1+ZEUS NC, CC and jet QCD fits**

H1-prelim-11-034, ZEUS-Prel-11-001

**H1 multijets at low  $Q^2$**

H1, EPJC 67, 1 (2010)

**H1 norm. multijets at high  $Q^2$  (unfold)**

H1-prelim-12-031

**ZEUS inclusive jets in  $\gamma^*p$**

ZEUS, Nucl. Phys. B 864, 1 (2012)

**D0 incl. jets, approx. NNLO**

D0, PRD 80, 111107 (2009)

**D0 angular correlations, NLO**

D0, Phys. Lett. B 718, 56 (2012)

**ATLAS incl. jets, NLO**

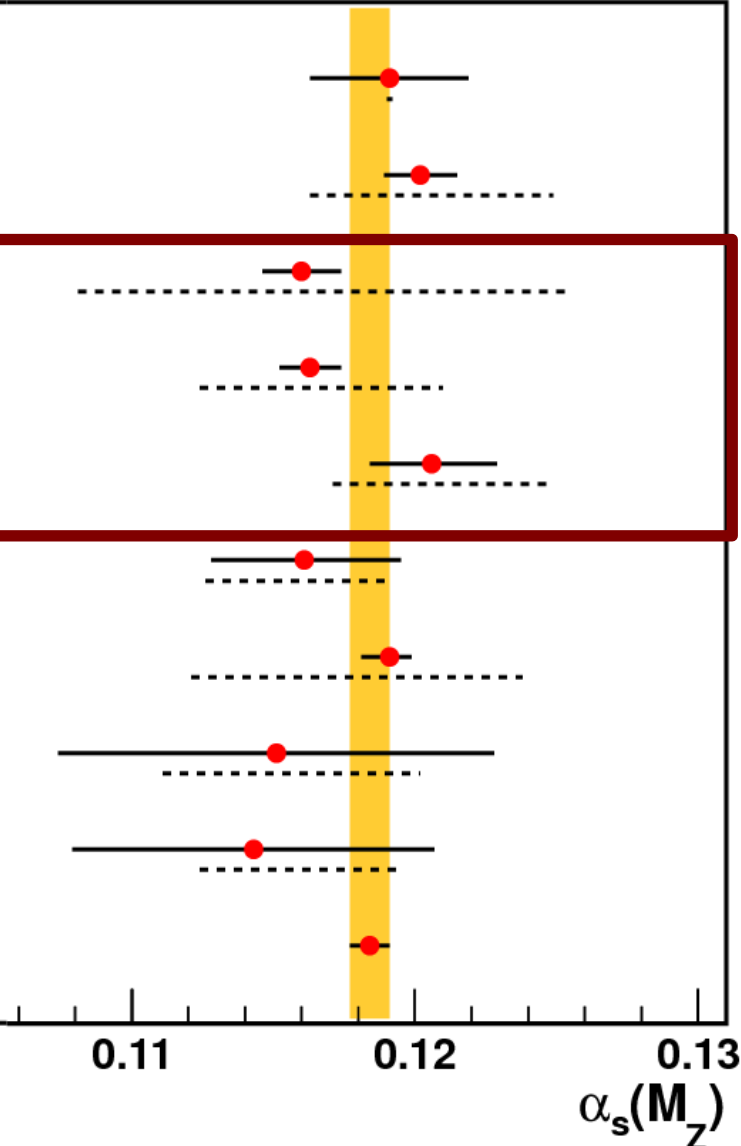
B. Malaescu et al., EPJC 72, 2041 (2012)

**CMS R3/2, NLO**

CMS PAS QCD-11-003 (2013)

**World average**

J. Beringer et al. (PDG), PRD 86 010001 (2012)



Experimental precision  
as good as  
or better than  
others measurements

Theory uncertainty  
dominates  
and  
NNLO is needed

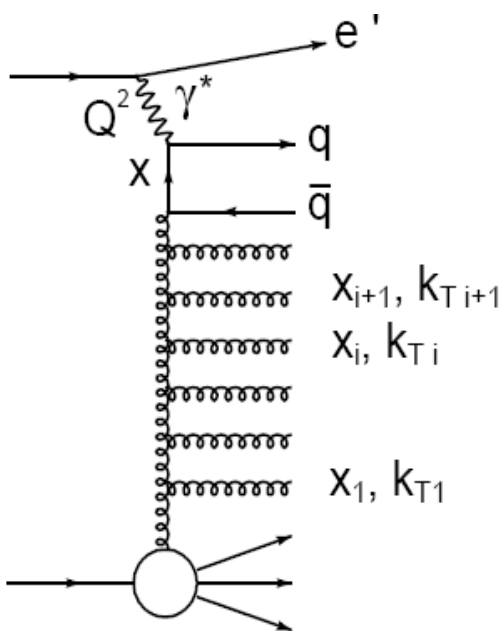
## Part II

**Test of QCD evolution mechanisms (*DGLAP* / *BFKL* / *CCFM*)  
using azimuthal correlation  
between the most forward jet and the scattered positron in DIS**

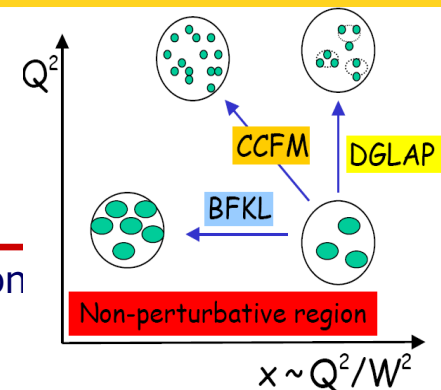
# QCD dynamics at low Bjorken-x

HERA : DIS at low Bjorken-x down to  $10^{-5}$   $\rightarrow$  energy in  $\gamma^*p$  cms is large ( $W_{\gamma^*p} \approx Q^2 / x$ )

- long gluon cascades exchanged between the proton and the photon
- pQCD – multiparton emissions described only with approximations :



- **DGLAP** (Dokshitzer-Gribov-Lipatov-Altarelli-Parisi) evolution  
resums terms  $\sim (\alpha_s \ln Q^2)^n$   
Assumes strong ordering of parton  $k_T$
- **BFKL** (Balitsky-Fadin-Kuraev-Lipatov) evolution:  
resums terms  $\sim (\alpha_s \ln(1/x))^n$   
No ordering in  $k_T$ , strong ordering in  $x_i$   
Transition from DGLAP to BFKL scheme expected at low  $x$
- **CCFM** (Ciafaloni-Catani-Fiorani-Marchesini) evolution:  
emitted partons are ordered in angles  
reproduces DGLAP at large  $x$  and BFKL at  $x \rightarrow 0$

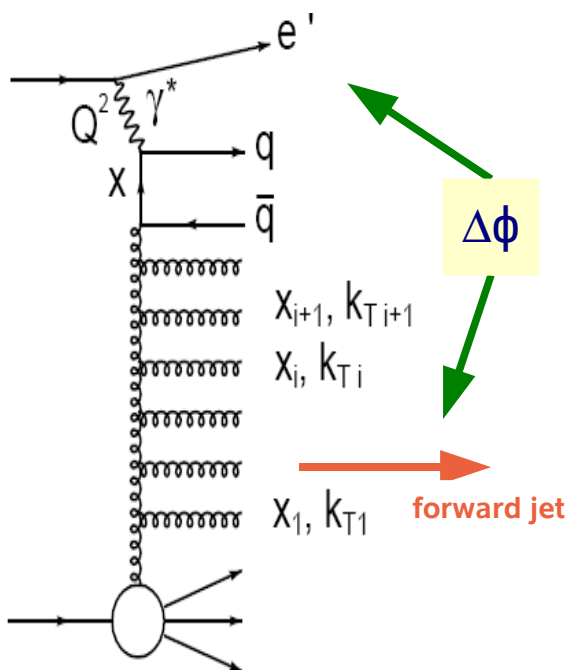


## Search at HERA for effects of parton dynamics beyond the standard DGLAP approach

- Strong rise of the proton structure function  $F_2(x, Q^2)$  with decreasing  $x$   
– well described by NLO DGLAP over a large range of  $Q^2$   
 $F_2$  measurement too inclusive to discriminate between different QCD evolution schemes
- Look at **hadronic final states** – reflecting kinematics, structure of gluon emissions



# Forward Jet Azimuthal Correlations



## Forward Jets in DIS (*Mueller – Navelet jets*) :

BFKL – more hard partons emitted close to the proton

Study high transverse momentum and high energy jets produced close to the proton ( forward region in LAB )

Suppress standard DGLAP evolution in  $Q^2$  :

$$p_{T, \text{fwdjet}}^2 \approx Q^2$$

Enhance BFKL evolution in  $x$  :

$$x_{\text{fwdjet}} = E_{\text{fwdjet}} / E_p \gg x_{\text{Bjorken}}$$



Data:H1, L=38.2 pb<sup>-1</sup>

### DIS selection

$$0.1 < y < 0.7$$

$$5 < Q^2 < 85 \text{ GeV}^2$$

$$0.0001 < x < 0.004$$

### Jets reconstructed in the Breit frame and boosted to LAB, all cuts in LAB

$$p_{T, \text{fwdjet}} > 6 \text{ GeV},$$

$$1.73 < \eta_{\text{fwdjet}} < 2.79$$

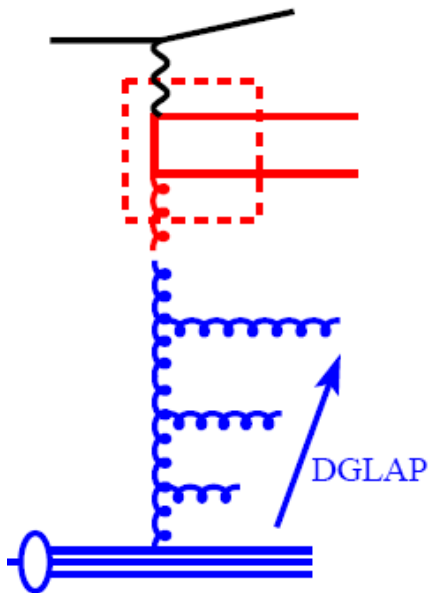
$$x_{\text{fwdjet}} = E_{\text{fwdjet}} / E_p > 0.035$$

$$0.5 < p_{T, \text{fwdjet}}^2 / Q^2 < 6.0$$

Measurement of the azimuthal angle difference  $\Delta\phi$  between the scattered positron and the forward jet as a function of the rapidity distance  $Y$  between them.

## RAPGAP - DGLAP

LO QCD matrix elements  
+ HO modelled by leading  
log parton showers

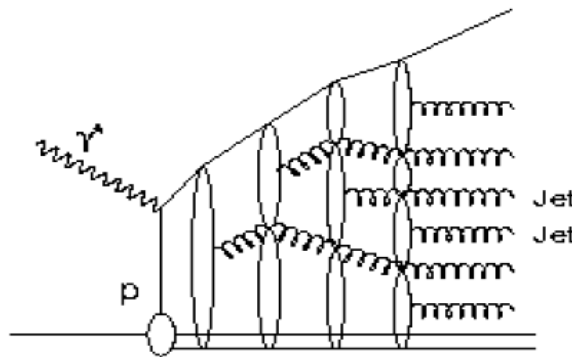


Single DGLAP ladder  
with  
strong ordering in  $k_T$

## ARIADNE Colour Dipole Model

CDM: QCD radiation  
from the colour dipole  
formed by the struck  
quark and  
the proton remnant.

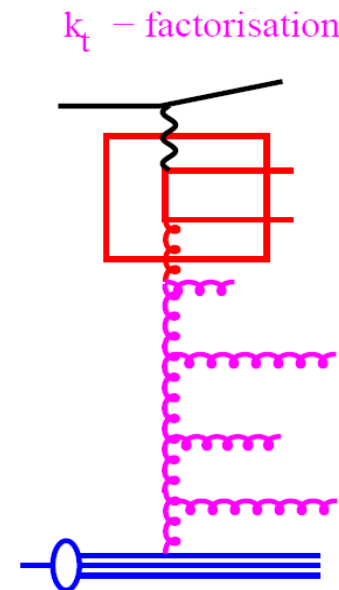
Chain of independently  
radiating dipoles formed  
by the emitted gluons.



BFKL- like Monte  
Carlo :  
random walk in  $k_T$

## CASCADE - CCFM

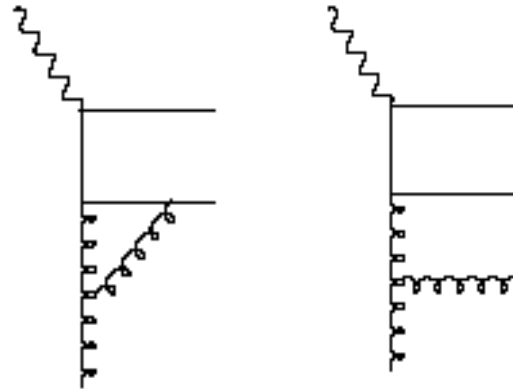
Off-shell QCD ME  
+ parton emissions  
based on the CCFM  
equation



Angular ordering of parton  
emissions

# Fixed order NLO DGLAP predictions

Forward jet cross sections – comparison with the predictions of pQCD at NLO ( $\alpha_s^2$ ) accuracy



- Forward jet analysis – reconstruction of jets in the Breit frame → at least dijet topology

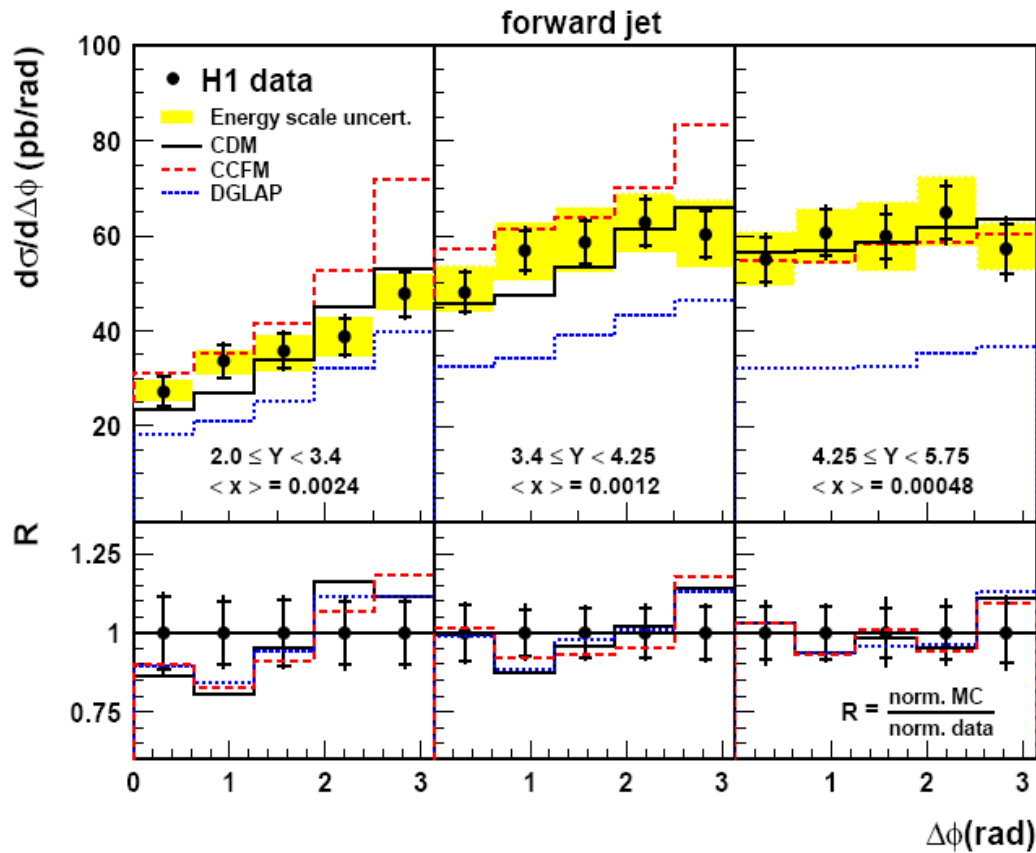
**NLOJET ++** program ( Nagy & Trocsanyi, 2001 ) :  
dijet production at parton level in DIS at NLO ( $\alpha_s^2$ )

- PDF : CTEQ6.6,  $\alpha_s(M_Z) = 0.118$
- parton level cross sections corrected for hadronisation effects using the RAPGAP model

# Forward Jet Azimuthal Correlations



At higher Y corresponding to lower x the forward jet is more decorrelated from the scattered positron



Cross sections best described by BFKL-like model CDM

- DGLAP predictions below the data
- CCFM (set A0) as good description as CDM at large Y

The shape of  $\Delta\phi$  distributions are similarly well described by all MC models

$Y = \ln(x_{\text{fwdjet}} / x)$  rapidity distance between the most forward jet and the scattered positron

$$R = \left( \frac{1}{\sigma^{\text{MC}}} \frac{d\sigma^{\text{MC}}}{d\Delta\phi} \right) / \left( \frac{1}{\sigma^{\text{data}}} \frac{d\sigma^{\text{data}}}{d\Delta\phi} \right)$$

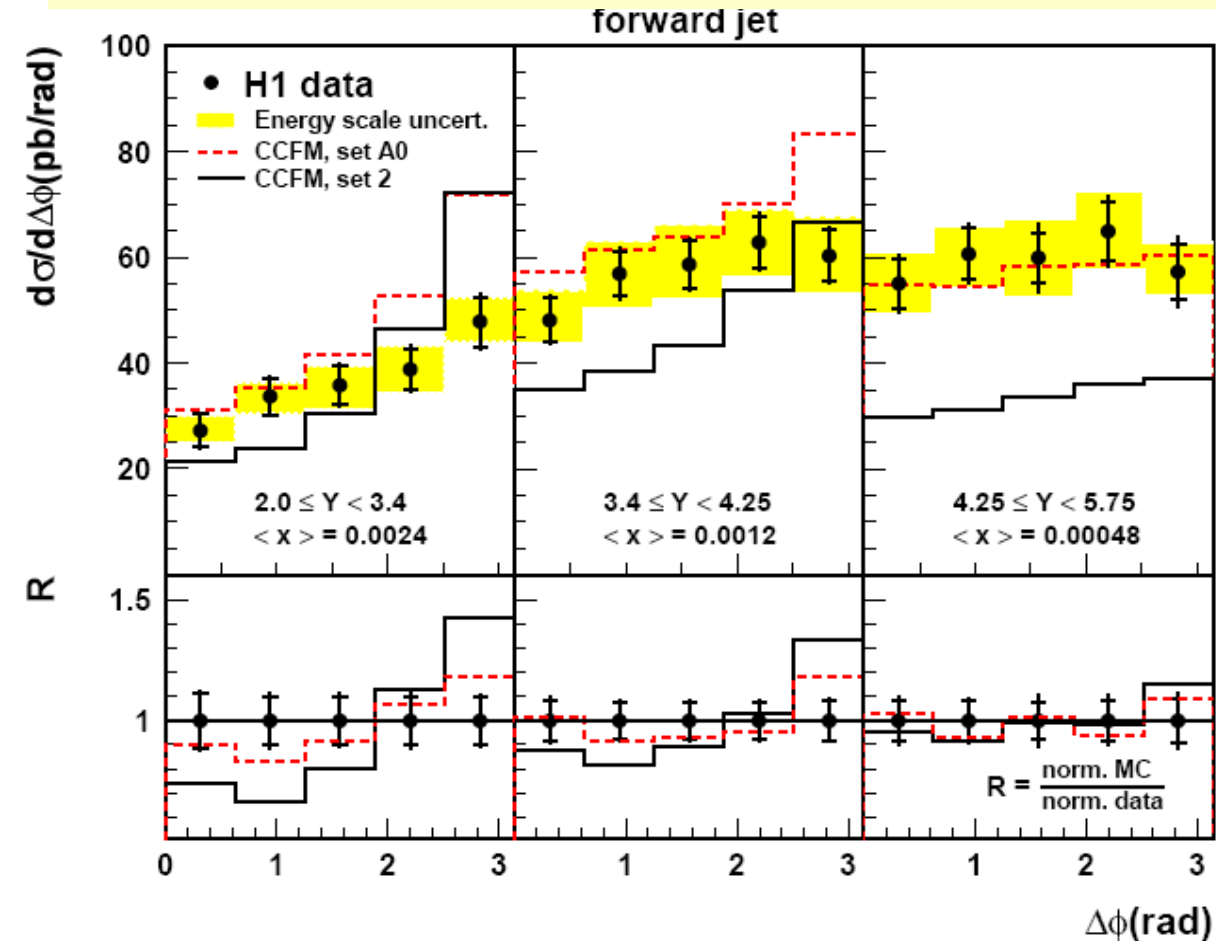
# Forward Jet Azimuthal Correlations



Different splitting functions used in unintegrated gluon density function (uPDF):

set A0 – only singular terms of the gluon splitting function

set 2 – includes also non-singular terms



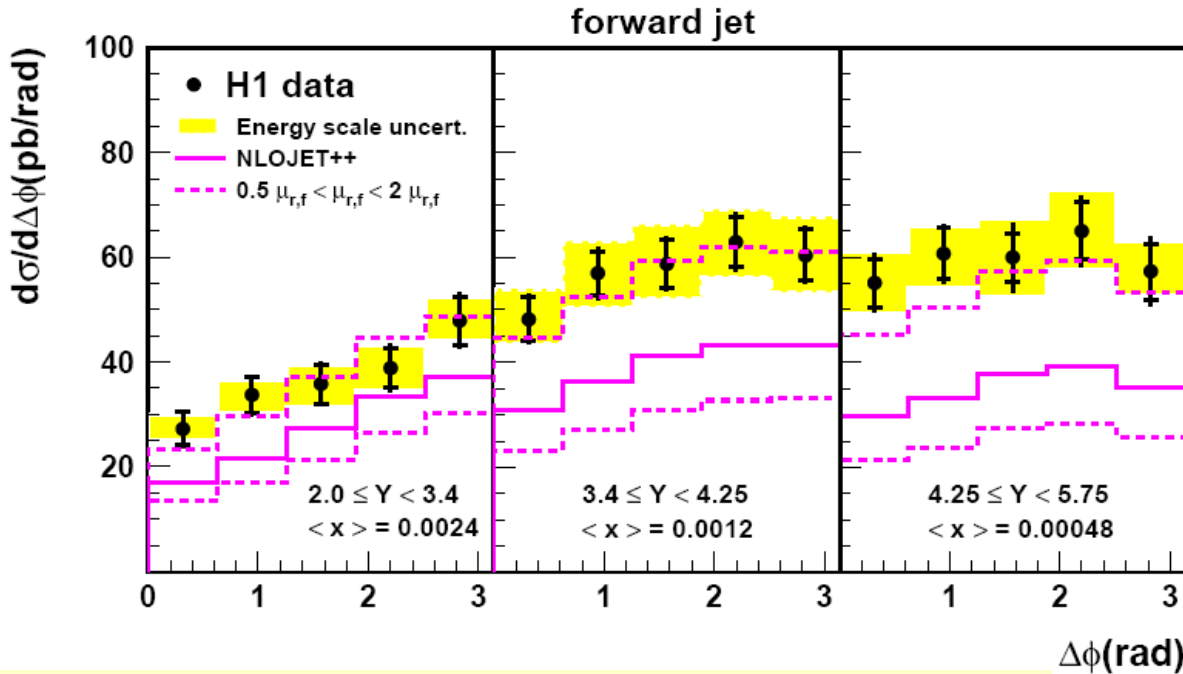
- **Cross sections**  
strongly depend on uPDF
- **Shape of  $\Delta\phi$  distributions**
  - at low Y shows sensitivity to uPDF
  - well described by the set A0

Predictions of the CCFM model depend on the choice of uPDF

# Forward Jet Azimuthal Correlations



## Comparison to NLO ( $O(\alpha_s^2)$ ) predictions



### NLO predictions

- shape of  $\Delta\phi$  distributions described, but central value too low
- large scale uncertainty ( of up to 50% ) indicates importance of higher orders

### NLOJET++

PDF : CTEQ6.6,  $\alpha_s(M_Z)=0.118$

Renormalisation and factorisation scales :

$$\mu_f = \mu_r = \sqrt{\frac{p_{T, fwdjet}^2 + Q^2}{2}}$$

Theoretical uncertainty :

factor 2 or  $\frac{1}{2}$  applied to  $\mu_r$  and  $\mu_f$  scales simultaneously

# Forward Jet Azimuthal Correlations (+central jet)



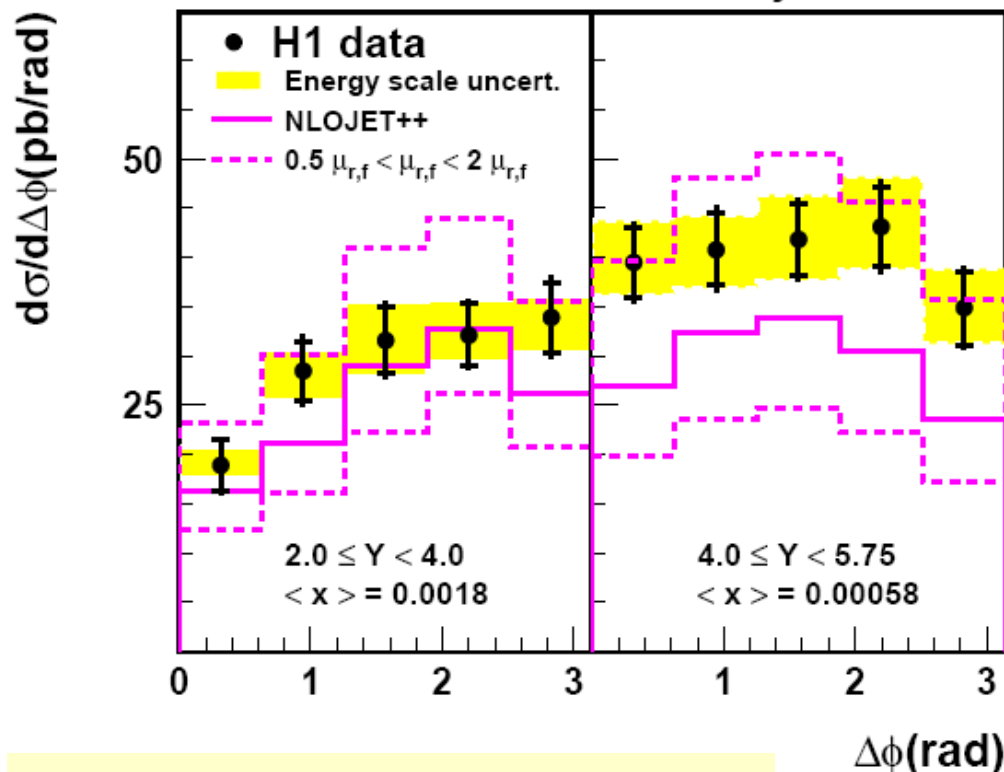
- **Subsample of events with forward jet + additional central jet**

$$p_{T,centet} > 4 \text{ GeV}, \quad -1 < \eta_{centet} < 1$$

$$\Delta\eta = \eta_{fwdjet} - \eta_{centjet} > 2 \text{ ( enhance radiation between the forward and central jet )}$$

- $\Delta\phi$  still between the forward jet and the scattered positron

forward and central jet



## NLO ( $O(\alpha_s^2)$ ) predictions

- at low Y reasonable description of the data
- at high Y, central value too small but still within theory uncertainty
- large scale uncertainty ( of up to 40% ) indicates importance of higher order contributions

## NLOJET++

PDF : CTEQ6.6,  $\alpha_s(M_Z)=0.118$

$$\mu_f = \mu_r = \sqrt{\frac{\langle p_T \rangle^2 + Q^2}{2}}, \quad \langle p_T \rangle = \frac{(p_{T, fwdjet} + p_{T, centjet})}{2}$$

# Summary

## Jets & $\alpha_s$

HERA jet data among the most precise data for precision test of QCD  
Perturbative QCD NLO calculations in general describe the data

### **Precision Measurement of Jet Production in DIS**

- inclusive jets, dijets and trijets measurements
- absolute and normalised single and double differential cross sections
- multi-dimensional unfolding of various measurements simultaneously

### **Precision Measurement of Inclusive Jet Production in Photoproduction**

- single and double differential cross sections  
based on the three jet algorithms ( $k_T$ , anti- $k_T$ , SIScone)
- the three jet algorithms give very similar results

**Extracted values of  $\alpha_s(M_Z)$**  from jet production in different regimes competitive with other measurements, precision dominated by theoretical uncertainties

**Running of  $\alpha_s$**  determination over a wide range of scale

**Theory:** Missing higher orders calculation (NNLO) often is dominated source of uncertainty

## Azimuthal correlation of forward jets in DIS

Cross sections as a function of  $\Delta\phi$  and rapidity separation between the forward jet and the scattered positron are best described by the BFKL – like model CDM

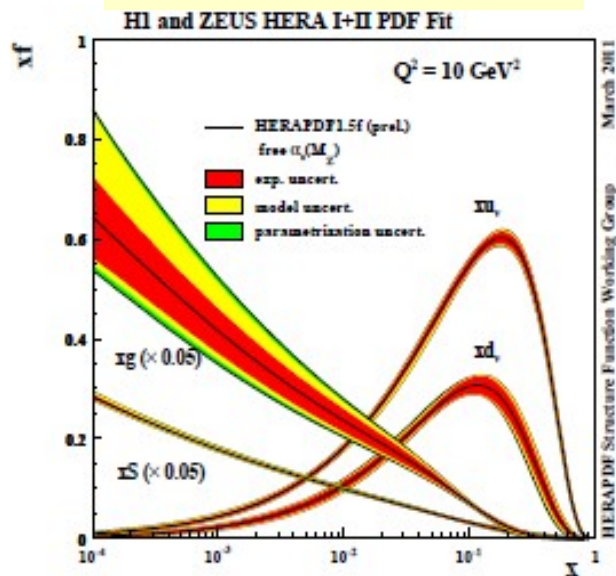
The shape of  $\Delta\phi$  distributions are well described by LO MC models with different QCD evolution schemes

NLO DGLAP predictions are in general below the data, but still in agreement within the large theoretical uncertainties

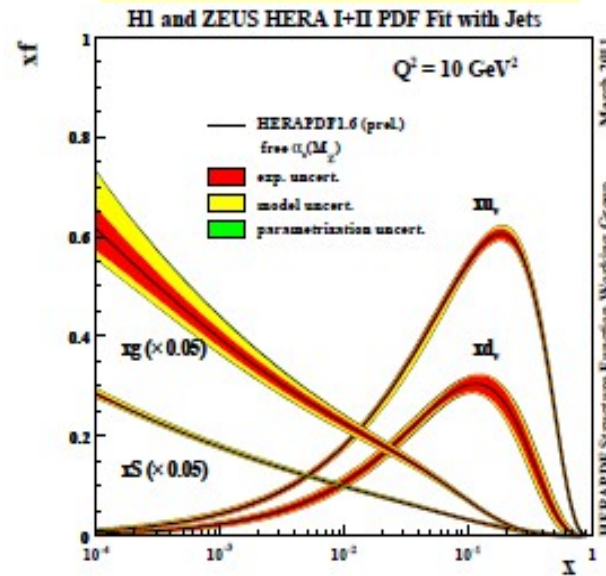


# Snapshot: Jets in PDFs fits

No Jets,  $\alpha_s(M_Z)$  free



+ Jets,  $\alpha_s(M_Z)$  free



Adding jets

Adding jet data dramatically decreases the low- $x$  gluon uncertainty, not only the experimental but also the model and parametrization uncertainties

See Achim Geiser's Wednesday talk