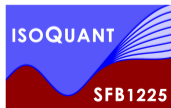


# Search for collective behaviour and multiparton interactions in $ep$ scattering at HERA

On behalf of the H1 and ZEUS collaborations

Dhevan Gangadharan

EPS, July 26th 2021



## Motivating questions

How small can a colliding system be while still exhibiting the collective features typically associated with the quark–gluon plasma in heavy-ion collisions?

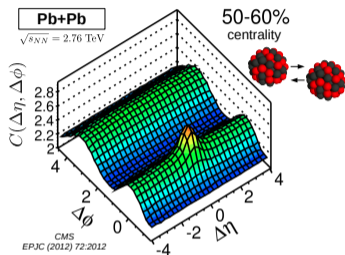
What kind of environment could collectivity evolve from?

Recent measurements using the H1 and ZEUS detectors will be presented in neutral current DIS and photoproduction.

New ZEUS publication: [arxiv:2106.12377](https://arxiv.org/abs/2106.12377) (submitted to JHEP)

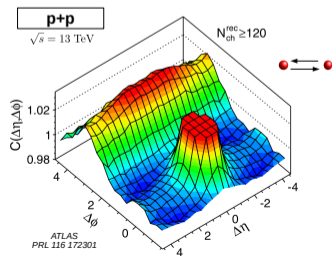
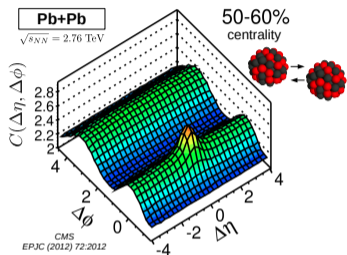
Recent H1 preliminaries: [Analysis note](#)

# Motivation for the analysis



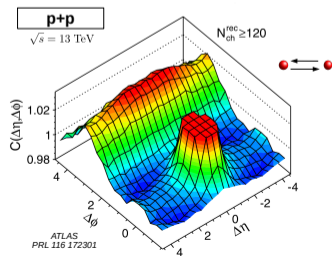
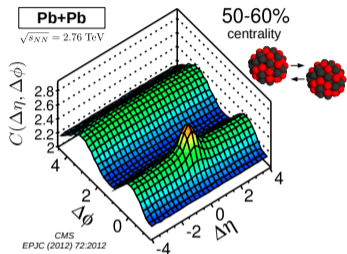
- Two-particle correlations in heavy-ion collisions show a clear **double ridge**, which is interpreted as a sign of fluid-like behaviour (QGP).
- $C(\Delta\eta, \Delta\varphi) = S(\Delta\eta, \Delta\varphi)/B(\Delta\eta, \Delta\varphi)$ ,  
 $S$  and  $B$  are formed from pairs from the same- and mixed-events, respectively.

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- Such collisions were thought to be too small to produce a thermally equilibrated QGP.

# Motivation for the analysis



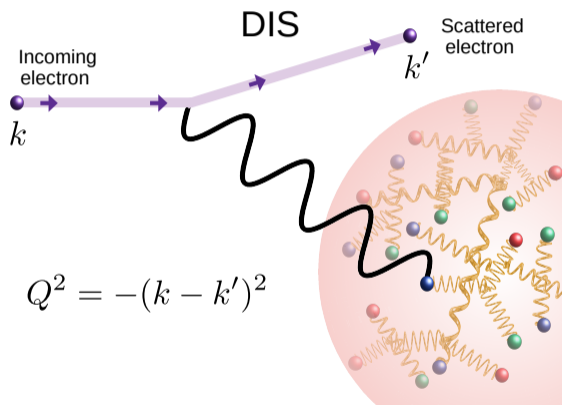
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- The start of the LHC revealed that high-multiplicity  $p + p$  collisions also have a double-ridge!
- Such collisions were thought to be too small to produce a thermally equilibrated QGP.
- **What about even more fundamental  $ep$  scattering at HERA??**

# The HERA collider and main experiments



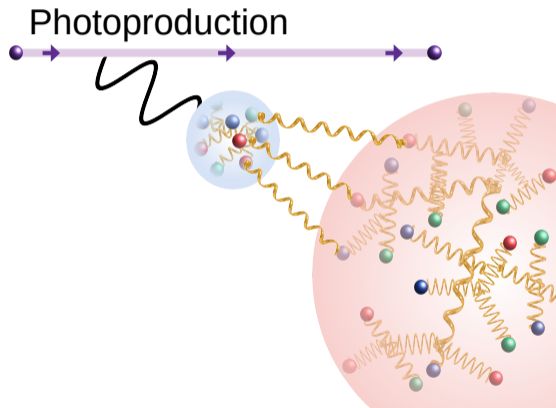
- Location: DESY, Hamburg, Germany
- Data taking: 1992 - 2007
- 27.5, 27.6 GeV electrons/positrons  
920 GeV protons  
→  $\sqrt{s} = 318, 319$  GeV
- HERA I+II:  
500  $\text{pb}^{-1}$  per experiment

# Deep inelastic scattering (DIS)



- DIS is defined by large virtualities:  
 $Q^2 \gg \Lambda_{\text{QCD}}^2$ .
- Transverse radius ( $R_t$ ) and longitudinal length ( $L$ ) of the probed region are given by:  
 $R_t \sim \frac{1}{Q}$   
 $L \sim \frac{1}{m_{\text{proton}} x}$  PRD 95 114008
- Neutral current (NC) DIS involves the exchange of photon or Z boson.

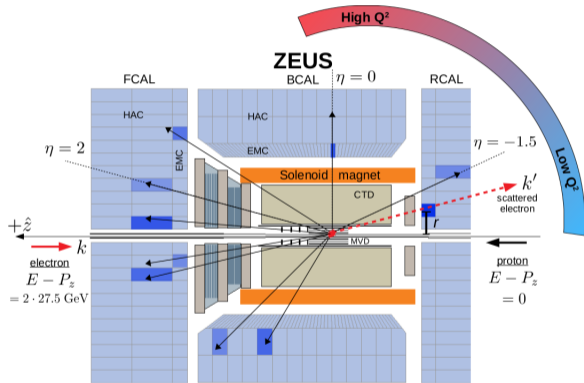
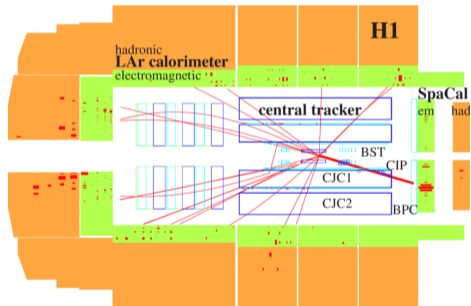
# Photoproduction (PhP)



- Photoproduction ( $\gamma p$ ) is defined by small virtualities:  $Q^2 \ll \Lambda_{\text{QCD}}^2$ .
- Exchanged photon may fluctuate into quarks and gluons.
- Larger interaction regions are probed.
- **Multiparton Interactions are possible.**
- Scattering is hadron-like.



# H1 and ZEUS detectors



## Event and track selection (main cuts only)

### DIS event selection

	scattered electron	$Q^2$	$\sum(E_i - P_{z,i})$	$N_{\text{ch}}$	$V_Z$
H1	in SpalCal	5 to 100 GeV <sup>2</sup>	35 to 75 GeV	$\geq 2$	-35 to +35 cm
ZEUS	in CAL	$\geq 20$ GeV <sup>2</sup>	47 to 69 GeV	$\geq 20$	-30 to +30 cm

### Photoproduction event selection

	scattered electron	$\sum(E_i - P_{z,i})$	$N_{\text{ch}}$	$V_Z$
H1	in tagger	NA	$\geq 2$	-30 to +30 cm
ZEUS	absent	$\leq 55$ GeV	$\geq 20$	-30 to +30 cm

### Track selection

	$p_T$	$\eta$	DCA
H1	0.3 to 3 GeV	-1.6 to 1.6 ( 0 to 5 for DIS in HCM)	< 5 cm in XY
ZEUS	0.1 to 5 GeV	-1.5 to 2.0	< 2 cm in XY and Z

There are several differences between the H1 and ZEUS analyses but compatible results are obtained nevertheless.

# Two- and four-particle correlation functions

Two-particle azimuthal correlations are measured:

$$c_n\{2\} = \langle\langle \cos n(\phi_i - \phi_j) \rangle\rangle.$$

$\phi_i$  is the azimuthal angle of particle  $i$ .

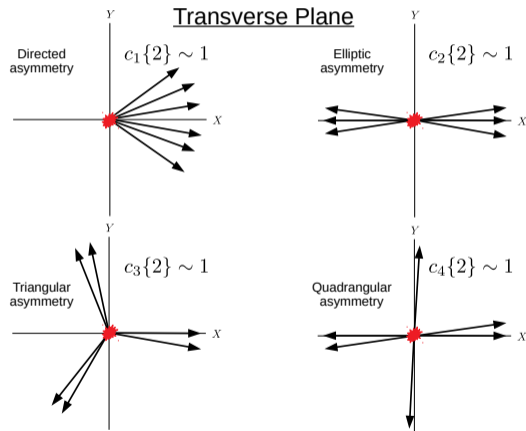
$n$  is the harmonic.

Four-particle cumulant correlations are also measured:

$$C_n\{4\} = \langle\langle \cos n(\phi_i + \phi_j - \phi_k - \phi_l) \rangle\rangle$$

$$c_n\{4\}(p_{T,1}) = C_n\{4\}(p_{T,1}) - 2 c_n\{2\}(p_{T,1}) c_n\{2\}$$

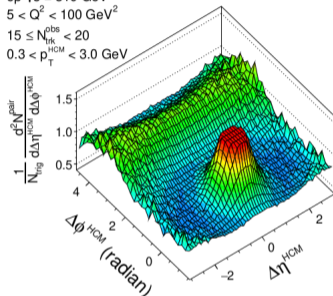
where  $p_{T,1}$  is the transverse momentum of particle  $i$ .



# Results: H1 & ZEUS ridge plots in DIS

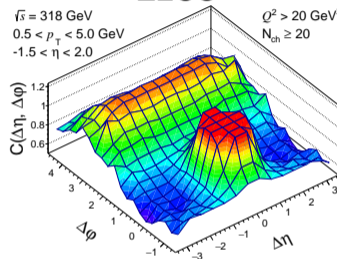
H1 Preliminary

ep  $\sqrt{s} = 319$  GeV  
 $5 < Q^2 < 100$  GeV<sup>2</sup>  
 $15 \leq N_{\text{trk}}^{\text{obs}} < 20$   
 $0.3 < p_{\text{T}}^{\text{HCM}} < 3.0$  GeV



ZEUS

$\sqrt{s} = 318$  GeV  
 $0.5 < p_{\text{T}} < 5.0$  GeV  
 $-1.5 < \eta < 2.0$   
 $Q^2 > 20$  GeV<sup>2</sup>  
 $N_{\text{ch}} \geq 20$



A near-side peak and away-side ridge are clearly visible.

**No visible double-ridge.**

Note: Kinematic selection differs between H1 and ZEUS.

# Results: H1 & ZEUS ridge plots in photoproduction

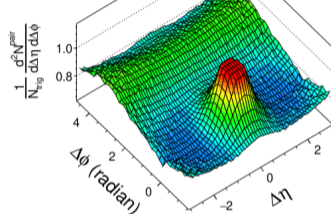
H1 Preliminary

ep photoproduction

$\langle W_{\text{yp}} \rangle = 270 \text{ GeV}$

$15 \leq N_{\text{trk}}^{\text{obs}} < 20$

$0.3 < p_{\text{T}} < 3.0 \text{ GeV}$



ZEUS

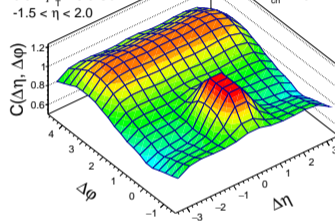
$\sqrt{s} = 318 \text{ GeV}$

$0.5 < p_{\text{T}} < 5.0 \text{ GeV}$

$-1.5 < \eta < 2.0$

$Q^2 < 1 \text{ GeV}^2$

$N_{\text{ch}} \geq 20$



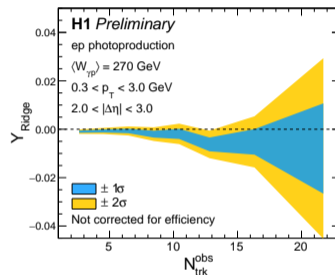
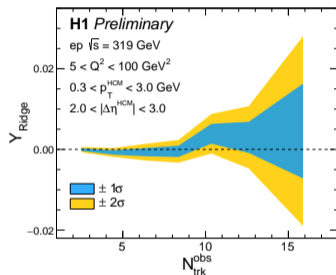
A near-side peak and away-side ridge are clearly visible.

**No visible double-ridge.**

Correlation strengths are significantly smaller than those in DIS.

Note: Kinematic selection differs between H1 and ZEUS.

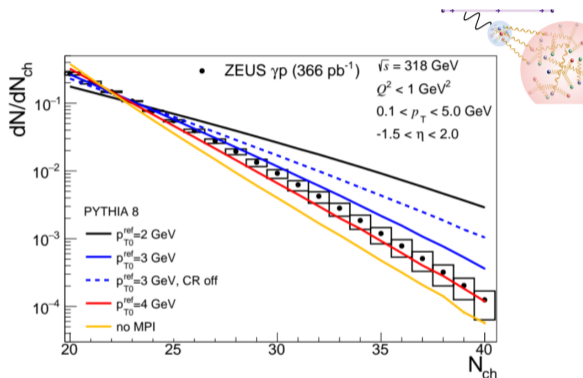
# Ridge yields in H1



Using a Zero-Yield-At-Minimum assumption, the ridge yields are extracted.

**Ridge yields in both DIS and Photoproduction are consistent with zero.**

# Results: $dN/dN_{ch}$



The level of MPI and IR divergencies are controlled by the  $p_{T0}$  parameter in PYTHIA.

It is used to regularize the interaction cross section in PYTHIA.

$$\frac{d\sigma}{dp_T^2} \propto \frac{\alpha_s^2(p_{T0}^2 + p_T^2)}{(p_{T0}^2 + p_T^2)^2}$$

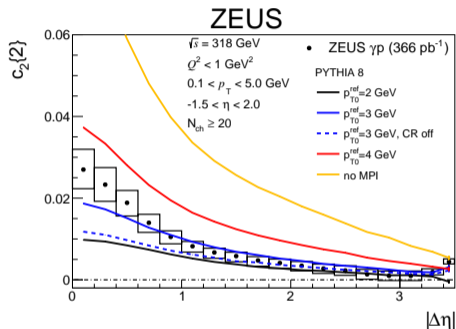
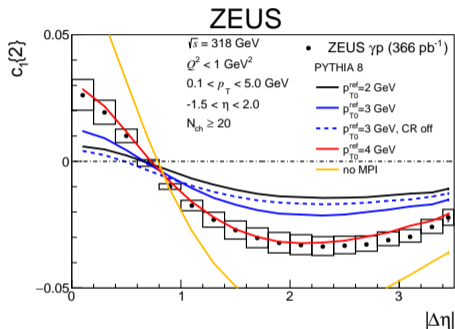
The energy dependence of this parameter is given by  $p_{T0} = p_{T0}^{\text{ref}} (W/7 \text{ TeV})^{0.215}$ , where  $W$  is the  $\gamma p \sqrt{s}$ .

**More MPI**  $\rightarrow$  **lower**  $p_{T0}^{\text{ref}}$

Colour Reconnection (CR) is PYTHIA's modeling of rescattering between partons from different MPIs

	$p_{T0}^{\text{ref}} = 2$	$p_{T0}^{\text{ref}} = 3$	$p_{T0}^{\text{ref}} = 4$
$\langle nMPI \rangle$	8.3	3.8	2.2

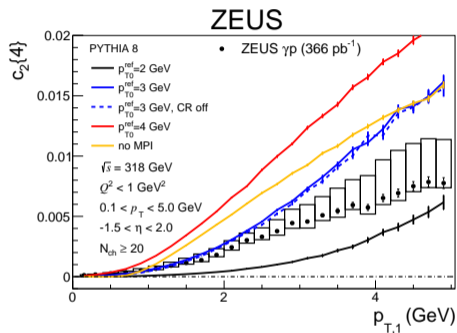
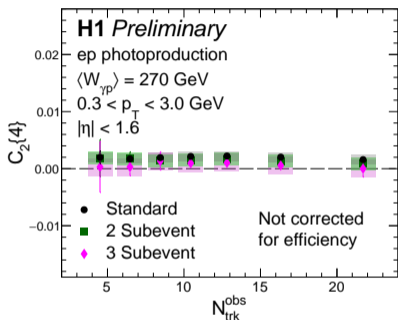
# Results: $c_1\{2\}$ and $c_2\{2\}$ versus $|\Delta\eta|$



- Correlation strengths are diluted by MPI.
- The scenarios of no MPI and very many MPI are disfavored.



# Results: Four-particle cumulants in photoproduction



- Four-particle cumulant is positive, which is in contrast to the negative values seen in non-central heavy-ion collisions.
- The scenarios of no MPI and very many MPI are disfavored.

## Summary

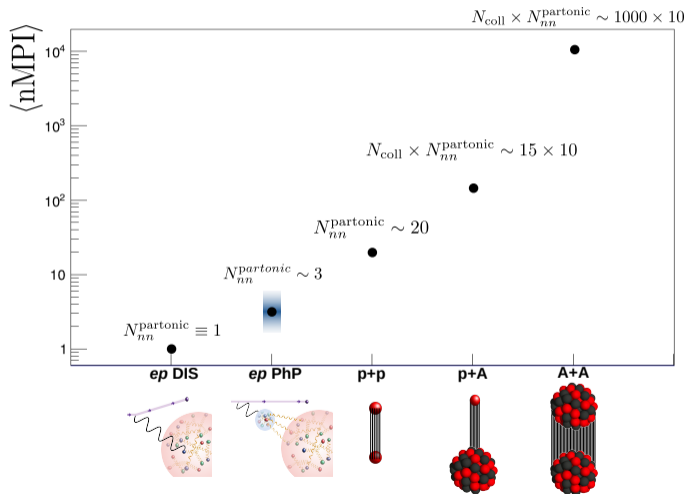
- Measurements of charged-particle azimuthal correlations have been presented using H1 and ZEUS data in  $ep$  photoproduction ( $\gamma p$ ) and NC DIS.
- There is no clear indication of a double ridge in either  $\gamma p$  or DIS. The observations **do not** reveal significant collective behaviour like that seen in heavy-ions or high-multiplicity hadronic collisions.
- The concept of multiparton interactions provides a useful tool to help understand the emergence of collective behaviour. It sets the stage for a potential rescattering phase.

	nMPI	Collectivity
$ep$ photoproduction	$\sim 3$	No
$pp$ high-multiplicity	$\sim 20$	Yes

The initial states in both systems may be similar in their spatial extent but completely different in the number of MPI.

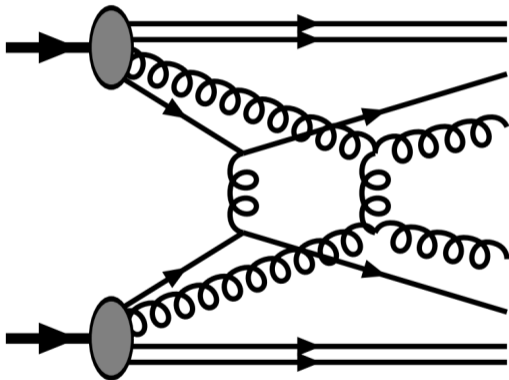
# Backup

# Illustration of MPI growth



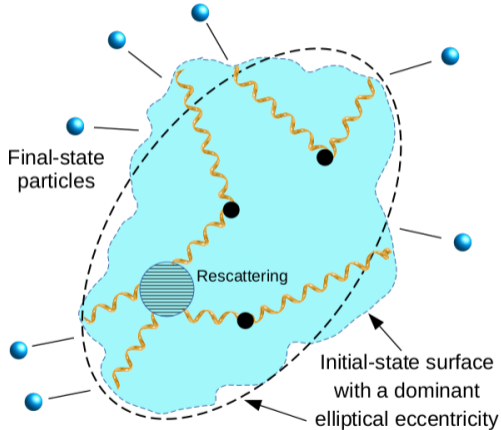
- Rough illustration of how nMPI grows from DIS to heavy-ions
- $N_{coll}$ : number of binary nucleon-nucleon collisions
- $N_{nn}^{partonic}$ : number of parton scatterings per binary nucleon-nucleon collision
- Estimates for  $N_{coll}$  taken from
  - Ann. Rev. Nucl. Part. Sci. 57, 205 (2007)
  - PRC 97 024905 (2018).
- Estimates for  $N_{nn}^{partonic}$  taken from PYTHIA

## Multiparton Interactions (MPI)



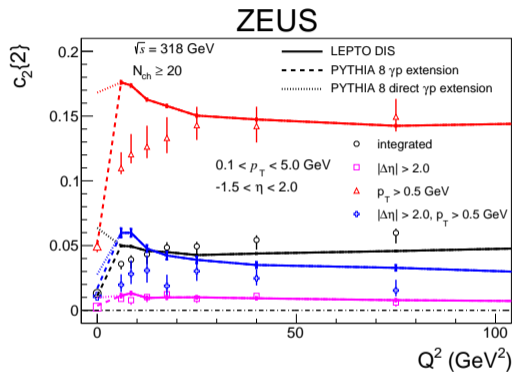
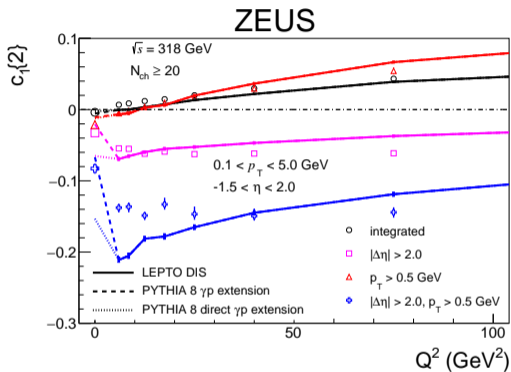
- MPI occur when there's more than one  $2 \rightarrow 2$  partonic scattering between the beam particles in a given event.
- If the scatterings are sufficiently hard ( $p_T \gtrsim 1$  GeV), they can be modeled in an event generator like PYTHIA.
- Established feature in high-multiplicity hadronic collisions. So far not conclusively observed in  $ep$  scattering.

## A subsequent rescattering phase is possible



- The initial scattering is shown here with 3 MPIs (black dots)
- Unlike in DIS, the spatial extent of this “initial state” is finite with an irregular shape in general.
- Subsequently, a phase of rescattering may occur, whereby a local thermal equilibrium might form.

# Results: $Q^2$ evolution of $c_1\{2\}$

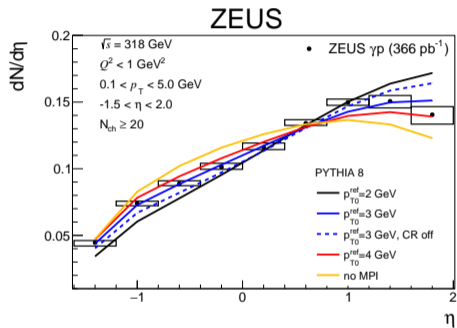
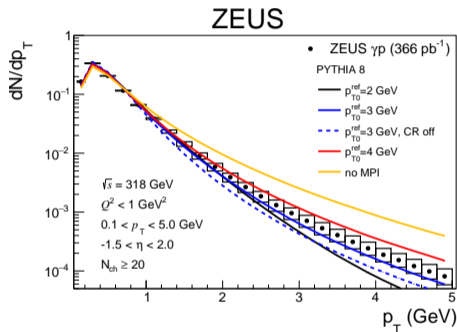


Photoproduction correlation strengths ( $Q^2 = 0$ ) are clearly diminished wrt those in DIS.

The LEPTO model of DIS gives a rough qualitative description of the data.

PYTHIA 8 with only the direct component of  $\gamma p$  predicts much stronger correlations than the full calculation (direct + resolved).

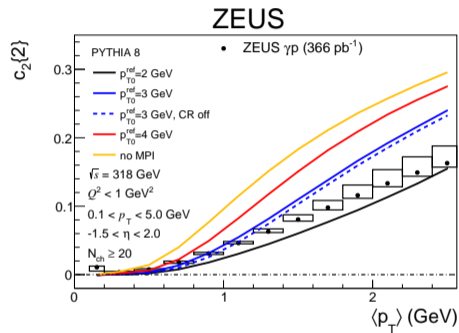
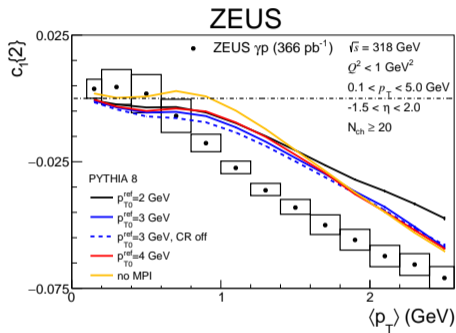
# Results: $dN/dp_T$ and $dN/d\eta$



- The scenarios of no MPI and very many MPI are disfavored.

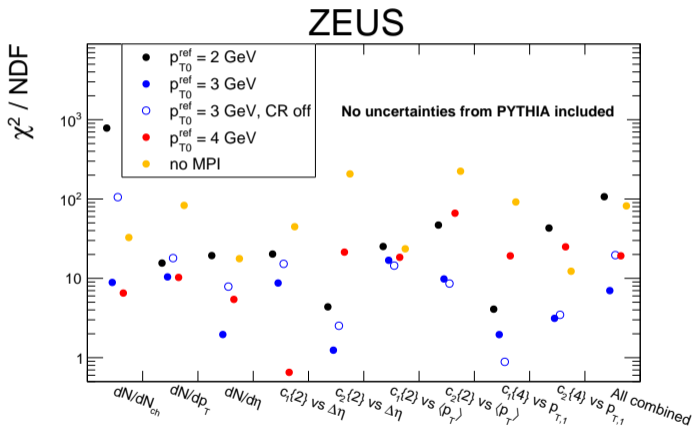


# Results: $c_1\{2\}$ and $c_2\{2\}$ versus $\langle p_T \rangle$

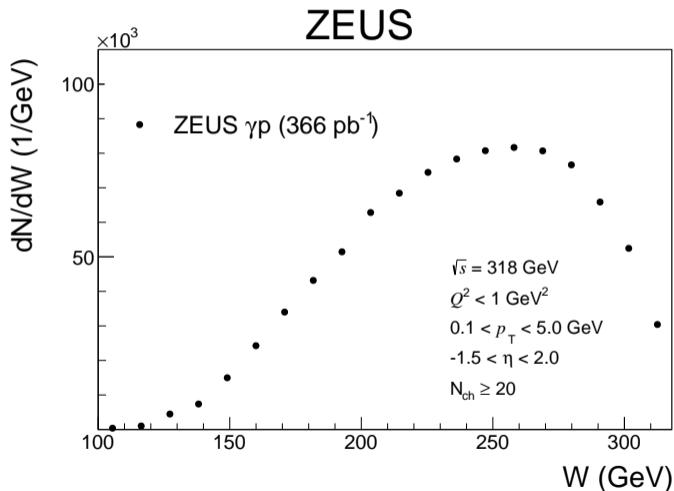


- $c_1\{2\}$  versus  $\langle p_T \rangle$  not sensitive to MPI and not described well by PYTHIA.
- More extreme levels of MPI are favored by  $c_2\{2\}$  versus  $\langle p_T \rangle$ .

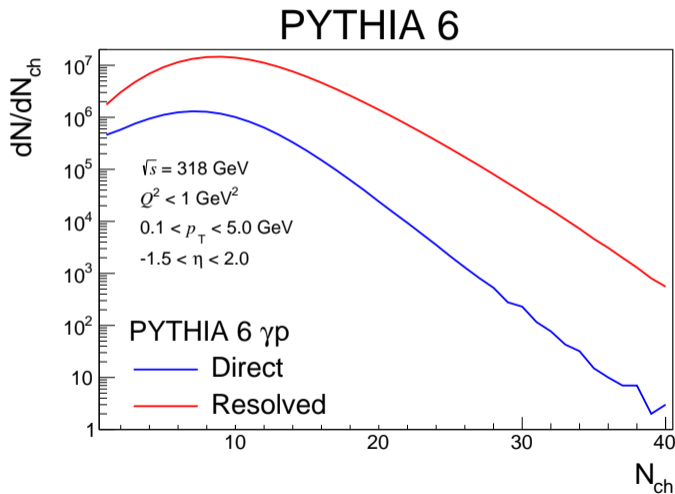
# Condensed view of PYTHIA 8 comparisons



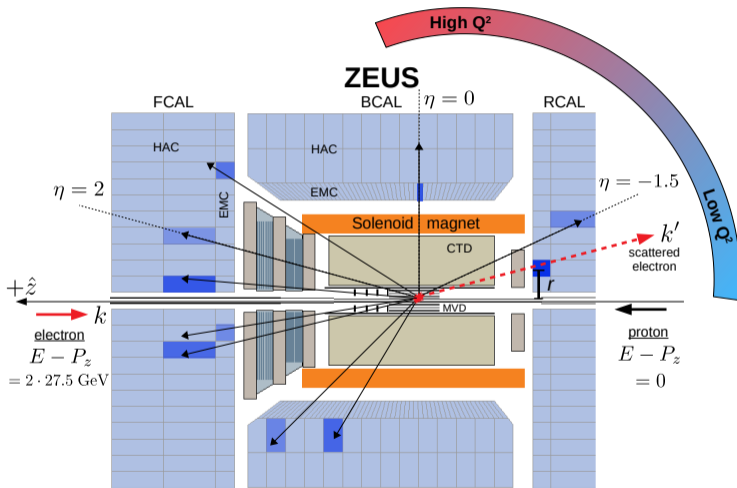
# W distribution



# Direct and Resolved event distributions



# ZEUS track selection

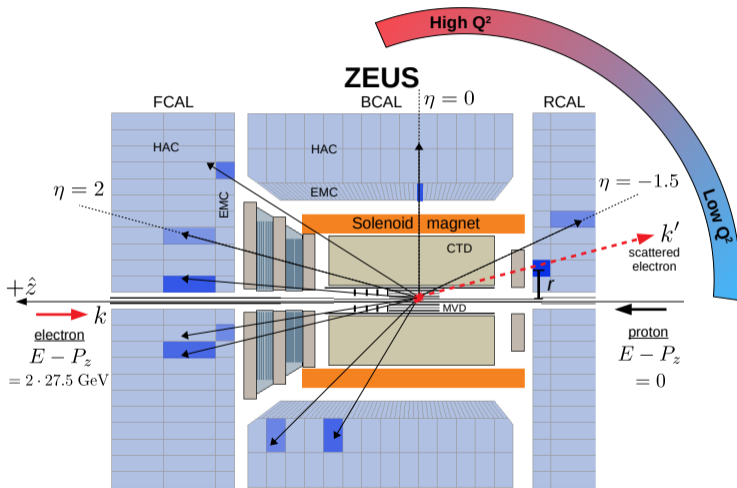


## Track selection for correlation analysis

- Reject scattered electron (if detected)
- $-1.5 < \eta < 2.0$
- $0.1 < p_T < 5.0$  GeV
- $\geq 1$  MVD hit
- $DCA_{XY,Z} < 2$  cm
- $\Delta R > 0.4$  (cone around scattered electron)

$$N_{\text{ch}} = \sum_i^{N_{\text{rec}}} w_i^{(1)}$$

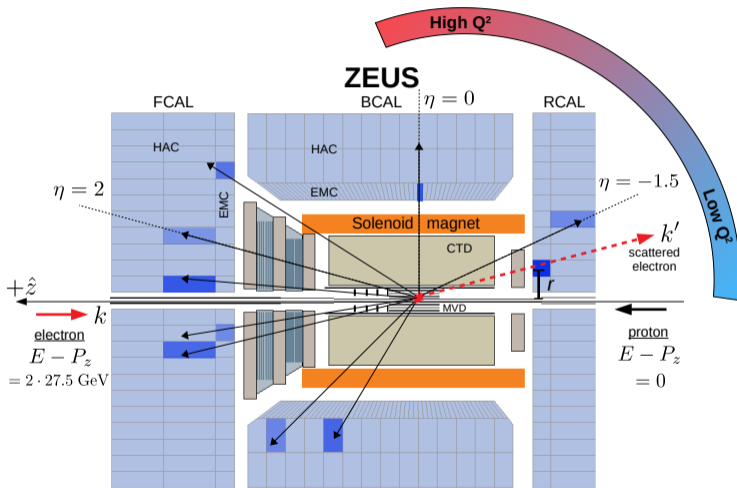
# ZEUS DIS event selection



## DIS Event selection (0.2 M)

- $N_{\text{ch}} \geq 20$
- DIS triggers
- electron probability  $> 90\%$
- $Q^2 = -(k - k')^2 > 5 \text{ GeV}^2$
- $k'_0 > 10 \text{ GeV}$
- $r > 15 \text{ cm}$
- $\theta_e > 1 \text{ rad}$
- $47 < \sum (E_i - P_{z,i}) < 69 \text{ GeV}$
- $|V_z| < 30 \text{ cm}$

# ZEUS photoproduction event selection



## Photoproduction event selection (5 M)

- $N_{\text{ch}} \geq 20$
- PhP oriented triggers
- electron probability  $< 90\%$
- $k'_0 < 15 \text{ GeV}$
- $\sum (E_i - P_{z,i}) < 55 \text{ GeV}$
- $|V_z| < 30 \text{ cm}$

## Tracking efficiency corrections

The efficiency correction weights for 1-, 2-, and 4-particle distributions are defined as:

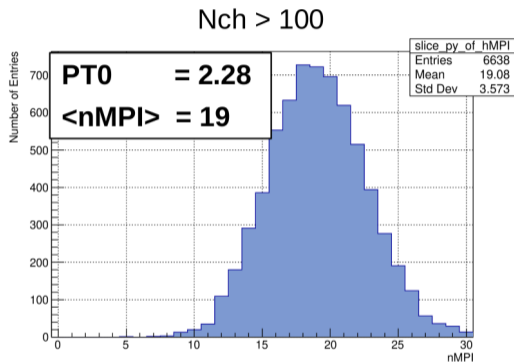
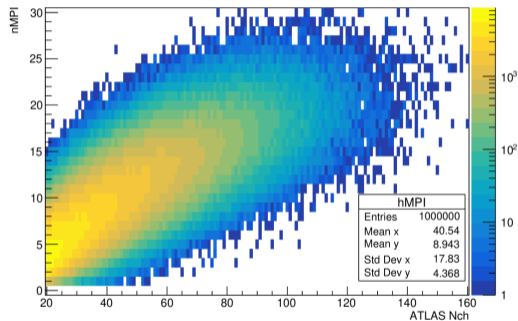
$$w^{(n)} = \frac{N_{gen}^n(\vec{x})}{N_{rec}^n(\vec{x})}$$

The are computed differentially in Monte Carlo simulations of the ZEUS detector:

dimension of $\vec{x}$	One-particle (n=1)	Two-particle (n=2)	Four-particle (n=4)
$x_1$	$\varphi$	$\varphi_1 - \varphi_2$	$\varphi_1 + \varphi_2 - \varphi_3 - \varphi_4$
$x_2$	$\eta$	$\langle \eta_i - \langle \eta \rangle \rangle$	$\langle \eta_i - \langle \eta \rangle \rangle$
$x_3$	$p_T$	$\langle p_{T,i} - \langle p_T \rangle \rangle$	$\langle p_{T,i} - \langle p_T \rangle \rangle$
$x_4$ (charge)	$q$	$ q_1 + q_2 $	$ q_1 + q_2 + q_3 + q_4 /2$
$x_5$	-	$N_{rec}$	$N_{rec}$



# nMPI in high-multiplicity $p + p$ PYTHIA at LHC energies



PYTHIA  $p + p$  events at  $\sqrt{s} = 13$  TeV were generated.

$N_{ch}$  was counted according to the ATLAS acceptance used in PRL 116 172301.

$-2.5 < \eta < 2.5$ ,  $0.4 < p_T < 50$  GeV